



## Research Paper

Volatile loss during heating of lunar mare simulants and related compositions<sup>☆</sup>Jennifer Sutherland<sup>a,b,\*</sup>, Kai-Uwe Hess<sup>a</sup>, Thilo Bissbort<sup>a,1</sup>, Simon Stapperfend<sup>c</sup>, Axel Müller<sup>d</sup>, Thomas C. Hansen<sup>b</sup>, Alexander Rack<sup>e</sup>, Donald B. Dingwell<sup>a,f</sup><sup>a</sup> Department of Earth and Environmental Sciences, Ludwig-Maximilians-Universität München, Theresienstraße 41, 80333 Munich, Germany<sup>b</sup> Institut Laue-Langevin, 71 Avenue des Martyrs, CS 20156, 38042 Grenoble Cedex 9, France<sup>c</sup> Chair of Space Technology, Technische Universität Berlin, Marchstraße 12-14, 10587 Berlin, Germany<sup>d</sup> OHB System AG, Manfred-Fuchs-Straße 1, 82234 Weßling, Germany<sup>e</sup> ESRF – European Synchrotron, 71 Avenue des Martyrs, CS40220, 38043 Grenoble Cedex 9, France<sup>f</sup> GEOLAB, Zhongfa Campus, Beihang University, Hangzhou 311115, China

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## ABSTRACT

Due to the scarcity of returned lunar regolith samples, simulants are currently widely used for testing and validation of technologies destined for resource extraction and manipulation on the Moon's surface. Deductions about thermal processing parameters (such as temperature ranges associated with sintering and melting) and material properties rely on a close fidelity of the terrestrially-sourced simulant materials. To evaluate the degassing behaviour of lunar regolith simulants during heating and thereby inform the validity of comparisons to the behaviour expected of actual lunar regolith, we compare profiles via simultaneous differential scanning calorimetry (DSC), thermogravimetric analysis (TGA), and evolved gas analysis (EGA) across a range of designated mare simulants, freshly recovered volcanic samples, and synthesised lunar-like glasses. Simulants generally exhibit higher mass loss (related to volatile release at sub-solidus temperatures) than fresh volcanic material, though overlap is observed with one simulant (JSC-1A). Of the fresh sources, Etna ash most closely limits the anticipated water loss difference to original lunar soils to within one order of magnitude. Results from complementary neutron powder diffraction (NPD) indicate that a lowering of the melting onset due to volatile fluxing persists ca. 200–300 °C beyond the major degassing range in samples of several grams. The possibility to reduce the simulant fidelity gap through critical source selection and their inherent limitations are discussed, and refined pre-treatment recommendations are made.

## 1. Introduction

The volatile (H, He, C, N, O, F, S, Cl) contents of lunar regolith have far-reaching implications for scientific and exploration endeavours. Volatile sources are highly sought, both to constrain the early history of the Earth-Moon system, and to reduce launch mass and cost for return missions by way of extracting elements for propellant and life support consumables (Hurley et al., 2023). Although the post-Apollo consensus perceived regolith as almost devoid of indigenous water (Epstein and Taylor, 1973, 1974; Fegley and Swindle, 1993), re-evaluation at

increased sensitivity and spectral reflectance data has overturned this notion. Within mare terrain, reported H<sub>2</sub>O/OH abundances average up to around 500 ppm in the upper mm of large pyroclastic deposits, and reach 83 ppm H<sub>2</sub> (742 ppm water-equivalent hydrogen) to depths of tens of cm in KREEP (potassium, rare-earth elements, and phosphorous)-rich regions (Li and Milliken, 2017; Lawrence et al., 2022). Direct measurements of hydrogen concentrations in returned regolith components, however, indicate the trend: agglutinates > bulk soil > volcanic glasses (Saal et al., 2008; Liu et al., 2012; Hauri et al., 2015). These small concentrations belie their significance on melt behaviour and properties

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(e.g. Burnham, 1979; Dingwell et al., 1996; Dingwell, 1998); nowhere is this more evident than in explosive volcanic eruptions (La Spina et al., 2022). For that reason, excess volatiles harboured in weathered products are highly undesirable in terrestrial material that acts to approximate the lunar surface for validating thermal processing strategies leveraged by in situ resource utilisation (ISRU).

A suite of lunar regolith simulants has been developed by space agencies, industry, and academia to support widespread technology testing in the context of crewed Artemis and Chang'e programmes, whilst preserving the carefully curated Apollo collection. Simulants should match the mineralogy, particle size and morphology, and bulk density of genuine lunar regolith, though fidelity varies in view of their intended purpose (Taylor et al., 2016; Martin and Wagoner, 2022; Slabic et al., 2024). Requirements are most stringent for construction (e.g. Taylor and Meek, 2005; Taylor et al., 2018; Fateri et al., 2019; Caprio et al., 2020; Farries et al., 2021; Han et al., 2024) and chemical extraction (e.g. Reiss et al., 2020; Schlüter and Cowley, 2020; Guerrero-Gonzalez and Zabel, 2023) efforts involving heating. Different approaches rely on manipulating regolith via sintering (grain bonding via atomic diffusion and neck formation at contact points below the melting point of any constituent, unless assisted by a minor liquid phase (e.g. Wadsworth et al., 2022)), through partial or complete melting.

Pre-treatment strategies for engineered simulants aim to limit degassing and thermal decomposition effects driven by non-lunar phases, whilst maintaining aspects of morphological and mineralogical similarity. Conversely, missing solar wind components ( $H_2$ , He) are not accounted for. The issue of volatile-driven vesiculation has been discussed in the context of mechanical properties (strength, conductivity) and density reduction (Song et al., 2019; Lim et al., 2023). Efforts to quantify (Petkov et al., 2024) and attribute gas species to their respective sources (Street Jr. et al., 2010; Petkov and Voecks, 2023) pave the way for targeted heating to remove secondary minerals in order to replicate lunar characteristics at elevated temperatures more closely. These primarily focus on optimising the most widely referenced simulant, JSC-1A, which has been exposed to atmospheric weathering for ca. 20,000 years (Duffield et al., 2006). Listed strategies vary from those that apply temperatures of up to 750 °C and exploit different sample environments (John and Rogers, 2021; Wilkerson et al., 2023) – recognising heating rate, vacuum, and atmosphere are known to affect volatile release (Kotra et al., 1982) – to somewhat arbitrarily-defined temperatures, or none at all.

Evaluation of those pre-treatment efforts is notably detached from the original material it seeks to imitate. This disconnect is likely due to blanket approaches to maximise volatile-derived mass loss, compounded by a paucity of pyrolysis results on lunar samples with which to compare. Updated measurements at improved sensitivities have been encouraged; however, recent studies continue to encounter issues of contamination (Mortimer et al., 2015; Verchovsky et al., 2020a). Whilst calorimetric analysis of samples with different mineralogy has been considered (Morrison et al., 2019; Whittington and Parsapoor, 2022; Whittington et al., 2023), and the fluxing effect of gases alluded to (Street Jr. et al., 2010), a quantitative perspective on volatile significance for sintering and melting applications with a view to informing pre-treatment is absent. Furthermore, while most attention has been placed on corrective measures, comparatively little discussion has been devoted to the potential for fresher source material in the ongoing development of simulants. Such pristine compositions align with mare, rather than highland, analogues, which pertain to sites outside the immediate Artemis exploration zone (Shearer et al., 2020).

To compare the spread in simulant degassing fidelity and evaluate the potential for terrestrially-sourced improvements, we measured several evolved gases and collected simultaneous thermogravimetric and heat flow data for 15 simulants and freshly collected volcanic samples. We catalogue these features through to 1300 °C in Section 3. Melt onset for larger volumes of two simulants is further probed using neutron powder diffraction. In Section 4, we sequentially attribute the

volatile loss origins and compare them against reported lunar abundances, before considering the impact of these differences on the as-yet undecided path(s) towards engineering of the Moon's surface.

## 2. Materials and method

### 2.1. Sample origins

#### 2.1.1. Mare simulants

Six simulants are included in this study: (1) ESA-supported EAC-1A, obtained from basanite in the Siebengebirge Volcanic Field, Germany (Engelschön et al., 2020). (2) NASA-developed JSC-1A, mined from ash deposit related to Merriam Crater in the San Francisco Volcano Field, Arizona, USA (Ray et al., 2010). (3) Space Research Technologies' (formerly Exolith Lab) LMS-1 (second manufacture batch), comprised of 32.0 wt% Butler Arts basalt of undisclosed provenance (procured instead from Merriam Crater as of September 2023), mixed with 19.8 wt% GreenSpar anorthosite from Kangerlussuaq, Greenland, 32.8 wt% pyroxene (Stillwater Mine, Montana, USA), 11.1 wt% olivine (Sisters Mining District, Washington, USA), and 4.3 wt% unspecified commercial ilmenite (Long-Fox et al., 2023). (4) Off Planet Research's OPRL2N, also matched (90.0 wt%) to the San Francisco Volcano Field, with a further 10.0 wt% originating from the Shawmere Anorthosite Complex, Ontario, Canada (Off Planet Research, 2025). (5) OPRL2NT, a high-titanium version with 14.4 wt% ilmenite, reducing the mare component by a corresponding fraction. (6) TU Braunschweig's TUBS-M, now continued at TU Berlin, sourced from a basalt complex west of Göttingen, Germany (Linke et al., 2020). Our supply of JSC-1A is missing the finest grain fraction due to prior mechanical testing. Bulk chemistry data, where available, is reported in Table S1.

#### 2.1.2. Volcanic deposits

Five of the six samples are recent volcanic products, collected shortly after eruption by members of the Department of Earth and Environment's volcanology group at LMU Munich. A lapillus sample from the August 2014 eruption of Mount Etna, Italy (approximately 20 mm in diameter), was collected by K.-U. Hess and subsequently crushed to powder. A reticulite sample from the 2018 lower East Rift Zone eruption of Kilauea, Hawaii (ca. 50 mm in diameter), was collected by D.B. Dingwell. From this sample, two sub-samples were extracted: the outer glassy rind ('Kilauea – outer') and an interior portion approximately 20 mm from the rim ('Kilauea – inner'). An ash sample from the 2021 Cumbre Vieja eruption, La Palma, was collected by U. Küppers. The Krafla sample was collected from a 1984 basaltic lava flow in Iceland by F. Wadsworth. A portion of this sample showing no visible signs of alteration was crushed to powder. All samples were stored under ambient conditions.

Meteorologically-exposed Hawaiian ash from an area to the west of the visitor centre for the Mauna Kea observatories (nearby the cinder cone Pu'u Haiwahini) was supplied by J. C. Hamilton of the University of Hawai'i at Hilo. The site was used for three multi-agency analogue field campaigns on ISRU in 2008, 2010, and 2012, including testing of the Regolith and Environment Science, and Oxygen and Lunar Volatile Extraction (RESOLVE), and the Volatile Analysis by Pyrolysis of Regolith (VAPoR) instruments (e.g. ten Kate et al., 2013). The sieved tephra was dried in a commercial bakery oven at 177 °C (350 °F) for two hours as feedstock for the Precursor ISRU Lunar Oxygen Testbed (PILOT) demonstration (Clark et al., 2009) in 2008, and has been in storage ever since.

#### 2.1.3. Processed analogue material

Synthesised 'green' and 'orange' glasses approximate the compositionally well-constrained pyroclastic beads first identified at Apollo 15 and 17 landing sites (Meyer, 2004, 2010). The glasses were prepared by mixing reagent-grade dried chemicals (Table S1) to match the corresponding lunar compositions. Oxide mixes totalling 100 g were loaded

into in Pt crucibles and melted at 1500 °C in a ZrO<sub>2</sub>-lined MoSi<sub>2</sub>-heated box furnace. The molten samples were then poured directly into Pt80Rh20 viscometry crucibles and subjected to high temperature viscometry to ensure homogeneity. Glasses were obtained post-viscometry by pouring the melts onto a water-cooled Cu plate. Reduced orange glass (Fig. S1) was generated by removing the viscometry crucible to a vertical gas-mixing tube furnace, where it was held and stirred for several days under a CO/CO<sub>2</sub> (18:2) gas mixture. The sample was manually removed from the base of the furnace and poured onto a steel plate.

A partially reduced version of the Mauna Kea sample was retained following the 2008 PILOT processing by Lockheed Martin (Clark et al., 2009). During that study, the ash was heated to ~900 °C for 90 min in a hydrogen reduction reactor to demonstrate low-yielding oxygen release from iron oxides.

## 2.2. Methods

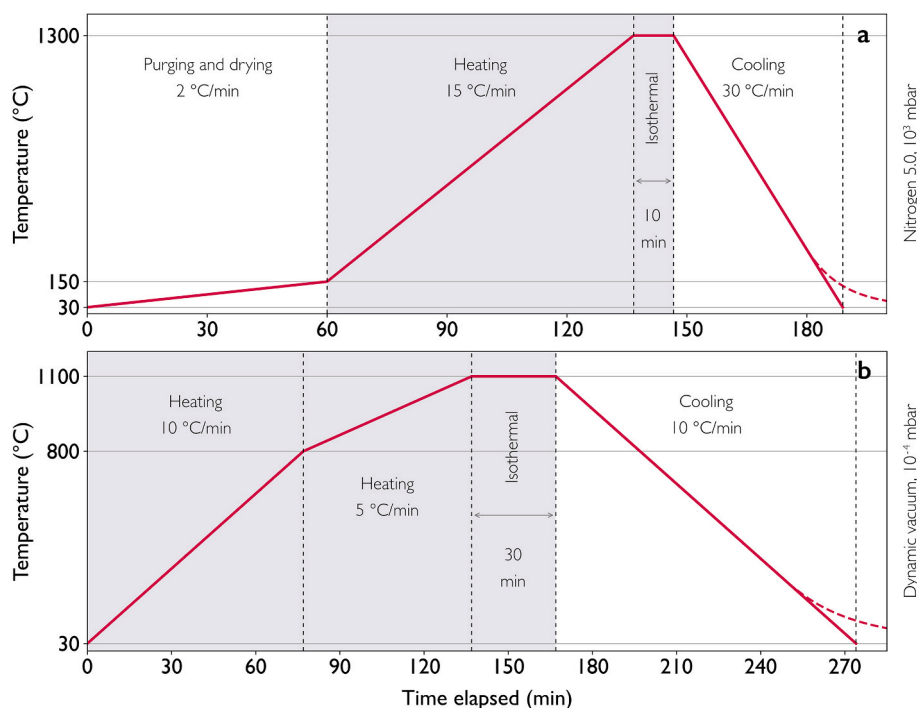
### 2.2.1. Differential scanning calorimetry (DSC), thermogravimetric analysis (TGA), and evolved gas analysis by mass spectrometry (EGA-MS)

A Mettler Toledo TGA/DSC 3+ with high temperature furnace was used for simultaneous thermal analysis (STA). For each run, 25–125 mg of as-received regolith simulant or crushed natural material was measured ( $\pm 0.1 \mu\text{g}$ ;  $\pm 5^\circ\text{C}$ ) in an alumina crucible and systematically corrected using reproducible baseline data conducted with empty crucibles. To synchronously identify released volatile species, the STA apparatus is linked to a ThermoStar Pfeiffer Vacuum GSD 320 EGA-MS via a transfer line heated at 150 °C to prevent gas condensation inside the capillary. An integrated vacuum pump generates a  $10^{-9}$  bar environment to retrieve the gases from the STA furnace, ensuring a near-concurrent response and analysis of the targeted evolved species (ca. 1 s delay between STA and EGA-MS signals). A quadrupole separates the gases, which have been ionised by a filament, according to their mass-to-charge ( $m/z$ ) ratio and trajectory stability within the applied oscillating electric fields. Arbitrary units are used for gas concentration as the signal pertains to ion current without direct physical interpretation. These

values represent the relative fraction of targeted species and are not normalised to sample mass.

The analytical protocol is depicted in Fig. 1a. A purging and drying segment is initiated from room temperature (ca. 30 °C) to 150 °C at 2 °C/min to purge the furnace, remove adsorbed water from the sample, and stabilise the baseline gas concentration, using a 100 mL/min nitrogen (N<sub>2</sub> 5.0, guaranteed purity of 99.999 %) flow. Next, the sample is heated from 150 °C to 1300 °C at 15 °C/min. This rate is a compromise based on returned volatile abundances: faster heating would increase the gas signal, improving detection sensitivity of low-abundance species, though dilute temperature resolution. At 1300 °C the sample undergoes a 10-min isothermal period, before cooling at 30 °C/min. Each analysis took ca. 200 min. Purging N<sub>2</sub> flow provided near-inert external conditions by displacing evolved gases and enabled the simultaneous detection of any H<sub>2</sub>O, CO<sub>2</sub>, SO<sub>2</sub>, O<sub>2</sub>, CH<sub>4</sub>, or H<sub>2</sub>S (extracted from unique  $m/z$  ratios of 18, 44, 64, 32, 16, and 34, respectively). Sensitivity is at 1 ppm level ( $10^{-14}$  A), except for H<sub>2</sub>O and O<sub>2</sub>, which is an order of magnitude greater owing to higher background in blank analyses. Empty baseline measurements indicate the supposedly inert N<sub>2</sub> atmosphere bears up to 200 ppm O<sub>2</sub>, likely contributed by diffusion from the Teflon connection tube and ineffective purging. Owing to mass balance sensitivity, higher fidelity vacuum conditions could not be recreated, although Wilkerson et al. (2023) demonstrate similar mass loss evolution under vacuum ( $< 10^{-4}$  mbar) and helium with JSC-1A. Furthermore, inert atmospheric conditions are more replicable in ongoing large-scale laboratory sintering and melting experiments.

Data are normalised to the mass at 150 °C to negate the effects of different storage conditions on adsorbed water quantity. Characteristic crystallisation ( $T_{\text{Cx}}$ , peak) and melting ( $T_{\text{Mx}}$ , onset) temperatures are identified, whereby the subscript x denotes distinct chronological events. The latter is constrained by the intersection of the inflectional tangent to the heating curve and the extrapolated slope of the endothermic melting baseline. Solidus ( $T_{\text{AS}} = T_{\text{M1}}$ ) and liquidus ( $T_{\text{AL}}$ ) temperatures are termed ‘apparent’, as at 15 °C/min they are expected to be an overestimation of the values at which samples attain thermal equilibrium. Furthermore, there may be residual crystalline phases that melt



**Fig. 1.** Thermal protocols. Results herein correspond to the signals recorded during the heating and isothermal segments (shaded grey). (a) Simultaneous TGA-DSC and EGA. (b) NPD: slower heating and increased isothermal period owing to longer signal integration time and lower maximum temperature. In both cases, true sample temperature begins to appreciably deviate from the idealised profile below approximately 300 °C on cooling (dashed curve).

above 1300 °C. We refer to melt onset rather than maximum as a closer approximation of the intrinsic temperature at which the transition starts, less influenced by variations in sample mass, shape, density, and crucible contact. Where sample temperature is marginally elevated above the applied experimental limit (up to 1316 °C was recorded), the data has been plotted sequentially within the isothermal section, which may be interpreted as a minimum of 1300 °C. An experimental artefact expressed as a jump in the heat flow signal marks the transition from dynamic to isothermal heating. Strong fluctuation in mass loss rate around 150 °C is also an artefact from the shift in heating programme.

### 2.2.2. Neutron powder diffraction (NPD)

To examine the replicability of the DSC-derived melting onset and potentially attribute decomposition reactions, two simulants (JSC-1A and TUBS-M) with contrasting degrees of weathering were additionally probed using the D20 instrument at the Institut Laue-Langevin (ILL). Datasets are listed under Sutherland et al. (2021, 2024). JSC-1A represents an extensively referenced simulant, with a source material that provides a common basis for comparison across multiple simulant series. The selection of TUBS-M reflects the authors' involvement in a broader campaign to characterise this European-origin modular simulant system. To promote earlier melting, further measurements were taken after mixing TUBS-M with 30 wt% of a distinct, well-characterised (Jagota et al., 1990; Vasseur et al., 2016; Wadsworth et al., 2016, 2017) soda-lime glass (Spherglass 2530, Potters Industries LLC, 45–90 µm).

A relatively long wavelength of 2.41 Å was prioritised over greater flux for high angular resolution ( $\Delta 2\theta \sim 0.1^\circ$ ) among densely clustered crystalline reflections. Individual patterns were integrated from the position sensitive detector ( $0 < 2\theta < 150^\circ$ ) over two minutes to improve counting statistics with a continuous temperature ramp. D20's dedicated dynamic vacuum furnace with a vanadium heating element imposed an upper limit of 1100 °C, although is sufficient to capture melt onset for basalt-like material (Chen et al., 2017; Giuliani et al., 2020). The nominal atmosphere at the pump decreases by three orders of magnitude to  $10^{-4}$  mbar at the sample position in the furnace vessel. The thermocouples ( $\pm 5$  °C) at the tip of the sample rod are ca. 3 cm from the neutron beam centre, and 1 cm above the sample surface. With one end at room temperature, the rod is a thermal bridge over ca. 60 cm, whose gradient likely lowers the recorded nominal temperature by an estimated 20 °C as compared to the temperature at the sample. This discrepancy has been accounted for in the display and interpretation of data.

The heating profile is highlighted in Fig. 1b. From room temperature the sample is raised to nominal 800 °C at 10 °C/min, which is then slowed to 5 °C/min from 800 until 1100 °C to better delimit phase changes approaching the melting interval. Each sample was held at 1100 °C for a further 30 min, before cooling at 10 °C/min. Vanadium cylinders (60 × 8 mm Ø) contained ca. 4 g of powder with 10 mm clearance and were loosely screwed with a vanadium cap to the sample rod. The attachment is not airtight, allowing evolved gases to be exchanged and removed via the furnace vessel. Neutrons permit larger sample volumes than X-ray alternatives, yielding a more representative bulk measurement for heterogeneous samples. Destructive testing of several grams of genuine lunar regolith, however, will require more compelling justification when several mg may suffice for synchrotron usage. Pertinently, neutrons have greater sensitivity to light elements such as hydrogen, carbon, and oxygen due to their relatively large scattering cross-sections, making them an ideal probe of crystalline structures incorporating volatile components.

## 3. Results

### 3.1. TGA-DSC-EGA

#### 3.1.1. Simulants

Characteristic temperatures and mass loss are presented in Table 1.

Mass loss within the tabulated ranges is summarised in Fig. 2. CH<sub>4</sub> and H<sub>2</sub>S were below detection limit for all runs. EAC-1A exhibits the highest mass loss of any sample (1.785 wt% total), with a prominent trimodal pattern peaking around 300, 550, and 750 °C (Fig. 3). Strong H<sub>2</sub>O signals are recorded during all three, with the last preceded by a major two-step CO<sub>2</sub> release between 600 and 850 °C. SO<sub>2</sub> emerges as early as 500 °C and is most significant between 800 and 900 °C, with a smaller contribution after melt onset at 1072 °C. The melting interval, between apparent solidus ( $T_{AS}$ ) and liquidus ( $T_{AL}$ ), spans a 115 °C range. Neither an endothermic peak that could be related to a glass transition nor an exothermic peak related to crystallisation of glass is present in the DSC curves, indicating its absence or sub-detection levels in the as-received simulant.

By contrast, the basaltic cinder comprising JSC-1A is glass-rich (~30 wt% glass (Patridge et al., 2025)), enhanced in the missing fine grain (50 wt% <20 µm (Collins et al., 2022)) – resulting in a crystallisation peak ( $T_{C1}$ ) at 824 °C. It demonstrates least alteration of all considered simulants at 0.161 wt% cumulative loss. Water content is about an order of magnitude less than EAC-1A, with two subtle contributions peaking around 300 and 650 °C, and almost exhausted by  $T_{C1}$ . The latter of these is associated with a pronounced CO<sub>2</sub> release. From 250 °C, very shallow oxidation is detectable through to ~1150 °C. The small upturn of the O<sub>2</sub> trace around 1300 °C is a persistent artefact accounted for by exact overlap of sample and background signals. The SO<sub>2</sub> signal following melt onset is comparable to EAC-1A; although spiked sources of SO<sub>2</sub> and CO<sub>2</sub> after  $T_{AL}$  are more prominent here, they remove little further mass. Among all simulants, no more than 9.3 % (JSC-1A) of the total cumulative mass loss is derived from the 1000–1300 °C range.

LMS-1 sustains consistently elevated water release up until approximately 1000 °C. No release of latent heat of crystallisation is observed, though glass is reported in the basaltic component (Long-Fox et al., 2023). Total mass loss is relatively high at 1.102 %. Three distinct contributions, from CO<sub>2</sub> then H<sub>2</sub>O, are clustered between 650 and 900 °C. The middle of these is correlated with a subtle sharp oxidation at 770 °C; an earlier event is also detected centred on 460 °C. CO<sub>2</sub> release remains relatively far above the baseline through to the second melt onset ( $T_{M2}$ ). Due to the considerable anorthosite fraction (19.8 %) and much later arrival of the apparent liquidus, the melting interval is comparatively drawn out over 210 °C.

OPRL2N crystallises expectedly close to JSC-1A (829 °C), given the basaltic cinder is derived from the same source. O<sub>2</sub> is, however, sharply consumed ~460 °C, in parallel with an additional release of CO<sub>2</sub> and H<sub>2</sub>O. There is a further small inflection in the water profile between 900 and 1000 °C, presumably owing to the 10 % anorthosite substitution. For these two simulants only, the mass loss rate is not led by water release above 200 °C. SO<sub>2</sub> is more conspicuous following the onset of melting at 1049 °C (within 2 °C of JSC-1A); this is the latest emergence of SO<sub>2</sub> among all simulants, with a negligible share prior to  $T_{AS}$ . Like LMS-1, the anorthosite fraction pushes  $T_{AL}$  considerably higher (1225 °C, 63 °C above JSC-1A).

In the higher-Ti version OPRL2NT, where ilmenite partially substitutes basalt, mass loss is approximately doubled. This difference is primarily expressed by increased H<sub>2</sub>O release under 600 °C, as well as a stronger CO<sub>2</sub> contribution in the 300–600 °C range. An additional decomposition around 940 °C simultaneously releases H<sub>2</sub>O, CO<sub>2</sub>, and SO<sub>2</sub>. There is no distinct oxidation event as in OPRL2N, but a baseline discrepancy persists from ~260–1200 °C.  $T_{C1}$  is identical to the former. Post-liquidus CO<sub>2</sub> bubbling is the most pronounced of any simulant. Although the anorthosite content remains consistent with OPRL2N,  $T_{AL}$  is lower by 55 °C.

TUBS-M displays a prominent H<sub>2</sub>O peak at 240 °C, such that the degassing profile is heavily skewed towards lower temperatures (81 % of the total mass loss is under 500 °C). Nearly all 0.797 wt% lost in this range is attributable to water. The relative CO<sub>2</sub> contribution, meanwhile, is the lowest of all simulants considered. From 600 °C evolved H<sub>2</sub>O remains minor, aside from a small event at 750 °C overlapping a



**Table 1**

Temperature estimations (all in °C to the nearest integer) of key thermal events ( $\pm 5$  °C) identified by DSC and mass loss contributions (3 d.p.) normalised to the dried sample mass at 150 °C. Crystallisation peaks ( $T_{\text{Ck}}$ ), apparent solidus ( $T_{\text{AS}} = T_{\text{M1}}$ ), distinct melting onsets ( $T_{\text{MX}}$ ), and apparent liquidus ( $T_{\text{AL}}$ ) correspond to labels in Figs. 3–5. Mass loss ranges refer to sample temperatures, except for the applied 1300 °C upper limit, where the maximum sample temperature varies.

	Crystallisation $T_{\text{C1}}$	Crystallisation $T_{\text{C2}}$	Apparent solidus $T_{\text{AS}}$	Melting onset $T_{\text{M2}}$	Melting onset $T_{\text{M3}}$	Apparent liquidus $T_{\text{AL}}$	-Δwt% 150–500 °C	-Δwt% 500–1000 °C	-Δwt% 1000–1300 °C 10' (% of total -Δwt%)	-Δwt% 150 °C - $T_{\text{AS}}$	-Δwt% $T_{\text{AS}}$ - 1300 °C 10'	-Δwt% 150–1300 °C 10'
<i>Simulants</i>												
EAC-1A	–	–	1072	1157	1168	1187	0.681	1.042	0.0611 (3.42 %)	1.743	0.042	1.785
JSC-1A	824	–	1051	1138	–	1162	0.072	0.074	0.015 (9.32 %)	0.150	0.011	0.161
LMS-1	–	–	1035	1193	1223	1245	0.315	0.751	0.036 (3.27 %)	1.074	0.028	1.102
OPRL2N	829	–	1049	1139	1175	1225	0.103	0.212	0.017 (5.14 %)	0.320	0.011	0.331
OPRL2NT	829	–	1066	1138	1155	1170	0.373	0.271	0.026 (3.88 %)	0.652	0.018	0.670
TUBS-M	–	–	1161	1181	–	1196	0.797	0.169	0.012 (1.23 %)	0.978	0.001	0.979
<i>Volcanic deposits</i>												
Etna	795	1122	1135	–	–	1185	–0.002	0.024	0.068 (75.56 %)	0.062	0.028	0.090
Kilauea (inner)	896	1069	1139	–	–	1162	0.050	0.079	0.012 (8.15 %)	0.152	–0.011	0.141
Kilauea (outer)	891	–	1119	–	–	1159	0.054	0.063	0.042 (26.42 %)	0.151	0.008	0.159
Krafla	829	1010	1099	1122	1136	1154	0.088	0.046	0.065 (32.83 %)	0.137	0.061	0.198
La Palma	803	–	1096	–	–	1135	–0.007	0.060	0.093 (63.70 %)	0.085	0.061	0.146
Mauna Kea (exposed)	–	–	1052	1172	–	1205	0.546	0.275	0.044 (5.09 %)	0.830	0.034	0.864
<i>Processed analogue material</i>												
Green glass (air)	786	848	1171	1282	–	>1300	–0.007	–0.014	0.123 (120.59 %)	–0.017	0.119	0.102
Orange glass (reduced)	723	812	1131	1201	–	1221	–0.007	–0.020	–0.021 (43.75 %)	–0.034	–0.014	–0.048
Mauna Kea (reduced)	–	–	1033	1122	–	1171	0.151	0.088	0.074 (23.64 %)	0.253	0.060	0.313

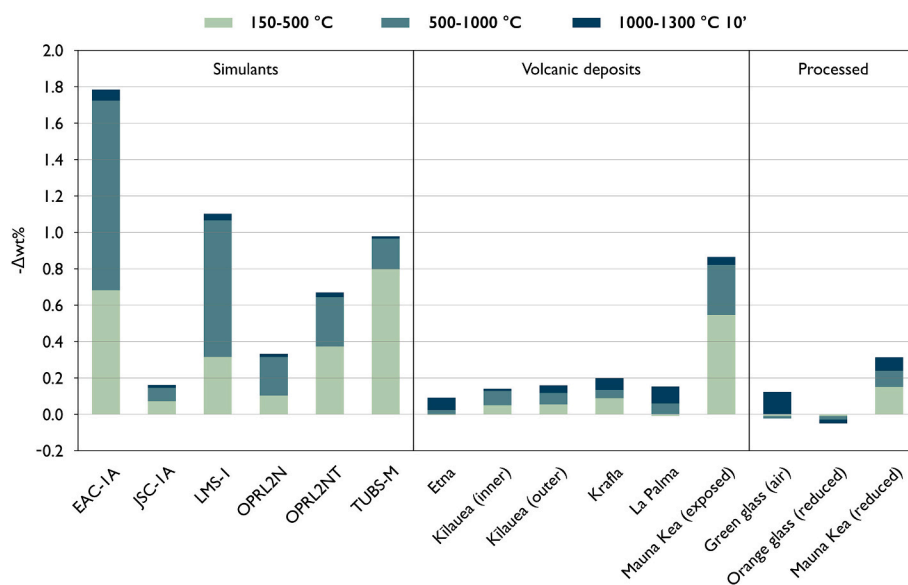


Fig. 2. Summarised cumulative mass loss over sample temperature ranges defined in Table 1.

secondary CO<sub>2</sub> peak. O<sub>2</sub> deviation is visually negligible. T<sub>AS</sub> is nearly 90 °C above the next closest simulant, giving a melt interval of only 36 °C, with consequently little associated degassing above its solidus (0.001 wt %).

### 3.1.2. Volcanic deposits

The volcanic samples (Fig. 4), with exception of the non-fresh Hawaii ash, undergo comparable or weaker de-volatilisation to the simulant exhibiting least mass loss (JSC-1A). Etna ash displays a double CO<sub>2</sub> peak below 500 °C, which then tapers until about 700 °C. The effect on weight loss is marginal. H<sub>2</sub>O release coincides with the larger of these two peaks between 400 and 500 °C, although oxidation offsets the mass loss resulting in slight net gain up to 500 °C (+0.002 wt%). Mass loss only perceptibly commences above 900 °C from H<sub>2</sub>O and SO<sub>2</sub>, and persists throughout the isothermal segment. The cumulative value of 0.090 wt% is the lowest of any natural sample considered. In stark contrast to the simulants, 76 % of Etna's weight loss occurs above 1000 °C; the absolute wt% in this range (0.068) is also larger than for any simulant. Major crystallisation occurs at 795 °C, with a second minor event at 1122 °C, just prior to T<sub>AS</sub> at 1135 °C. There is no discernible second melting event before T<sub>AL</sub> at 1185 °C, which lies within the range of simulants without an anorthosite constituent.

Crushed inner reticulite from Kilauea features a relatively high first T<sub>C</sub> at 896 °C, and a second at 1069 °C. Like Etna, water loss is consistently extremely small, though persists throughout the entire run, including at 1300 °C. Evolved CO<sub>2</sub> is the lowest of any fresh sample. Two isolated SO<sub>2</sub> releases around 500 and 1100 °C are observed. The melt interval is condensed over 23 °C, in which overlapping melt peaks are present. From 1300 °C CO<sub>2</sub> is consumed by the sample, contributing to an upturn in the cumulative mass loss curve; however, the net loss above 1000 °C remains positive.

The outer portion of the reticulite is slightly more volatile-enriched. The CO<sub>2</sub> signal is higher throughout, with a main degassing plateau between ~375 °C (coinciding with stronger oxidation) and 875 °C, and an additional small H<sub>2</sub>O release between 1000 and 1100 °C. A marginally wider melt interval has only one clear melting event. Like the inner sample, mass gain is seen at the end of the isothermal segment. The capacity for the barely perceptible O<sub>2</sub> baseline deviation to account for this increase is more clearly demonstrated with the orange glass. However, an assessment of oxidation trends at these low concentrations is complicated by the aforementioned background impurity.

Krafla ash has a markedly higher SO<sub>2</sub> contribution before the onset of

melting. Its melt interval is more staggered: three distinct events are observed during this 55 °C period. Like the outer reticulite, an early CO<sub>2</sub> peak under 500 °C is accompanied by more pronounced oxidation. Of the fresh samples, Krafla has the greatest water release in this first interval; more so than JSC-1A.

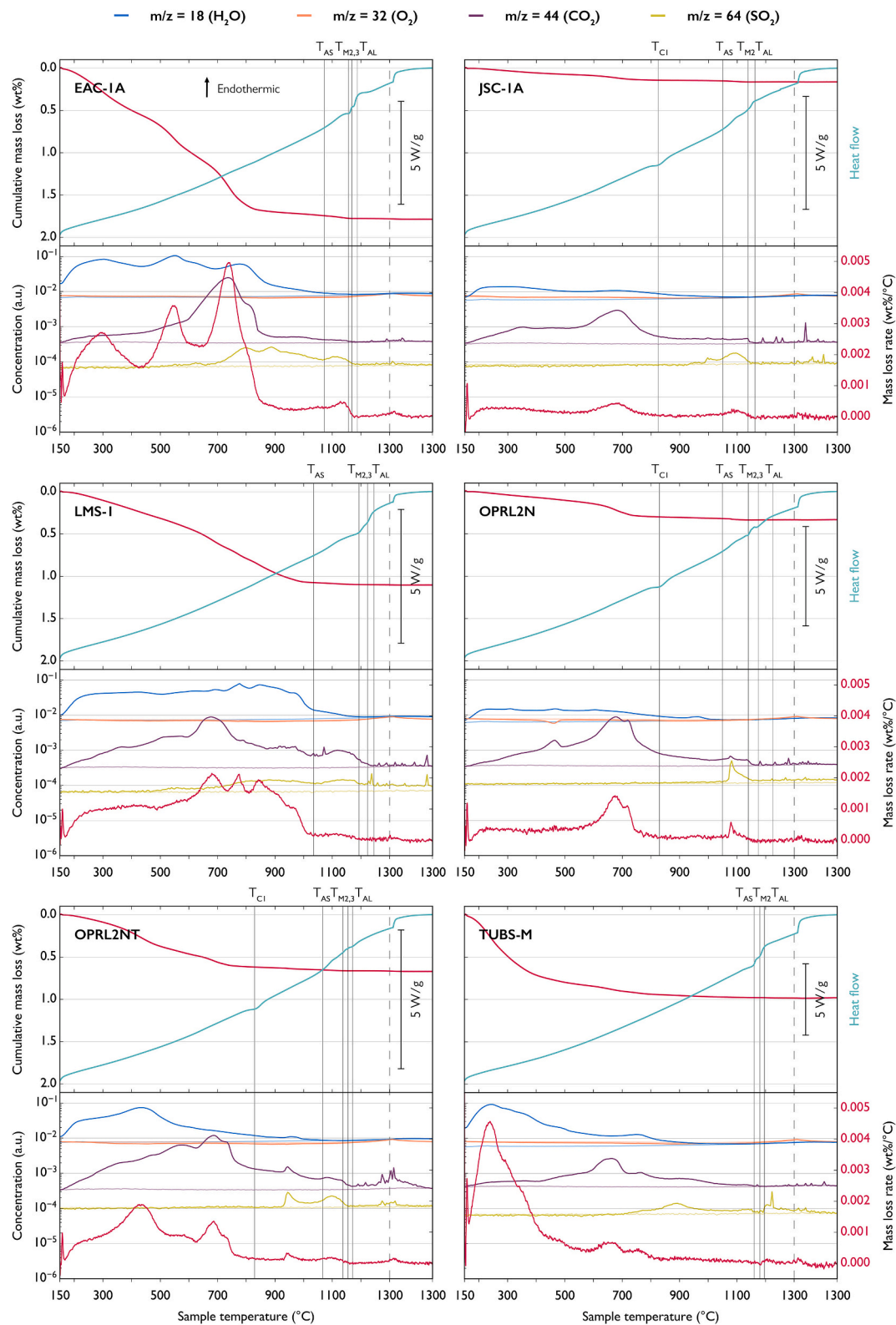
La Palma ash exhibits even greater weight gain than Etna below 500 °C (+0.007 %), attributed to oxygen uptake coinciding with very low initial H<sub>2</sub>O. Evolved CO<sub>2</sub> remains small. SO<sub>2</sub> is a prominent contributor to mass loss, such that the sample has the highest absolute fractional mass loss (0.093 wt%) above 1000 °C of all the unprocessed material. Unlike Krafla, the major release is delayed until after T<sub>AL</sub>. Common to the fresh deposits, the melting interval is on the order of only 40 °C.

Exposed Mauna Kea ash has a tapering H<sub>2</sub>O trace from 200 °C until T<sub>AS</sub>, and a broad CO<sub>2</sub> profile that extends to ~1150 °C. Oxygen is more heavily consumed from 250 °C until liquidus, peaking around 370 °C. Total mass loss is nonetheless comparatively high at 0.546 wt% up to 500 °C. SO<sub>2</sub> is released continuously from 500 °C in minor quantities, with a barely detectible contribution to mass loss at 1300 °C, alongside CO<sub>2</sub>. No crystallisation is observed. An additional subtle endotherm before T<sub>M2</sub> is unlabelled, as quantifying an onset from the shallow amplitude would be overly subjective. T<sub>AL</sub> is the highest of the six volcanic deposits, at 1205 °C.

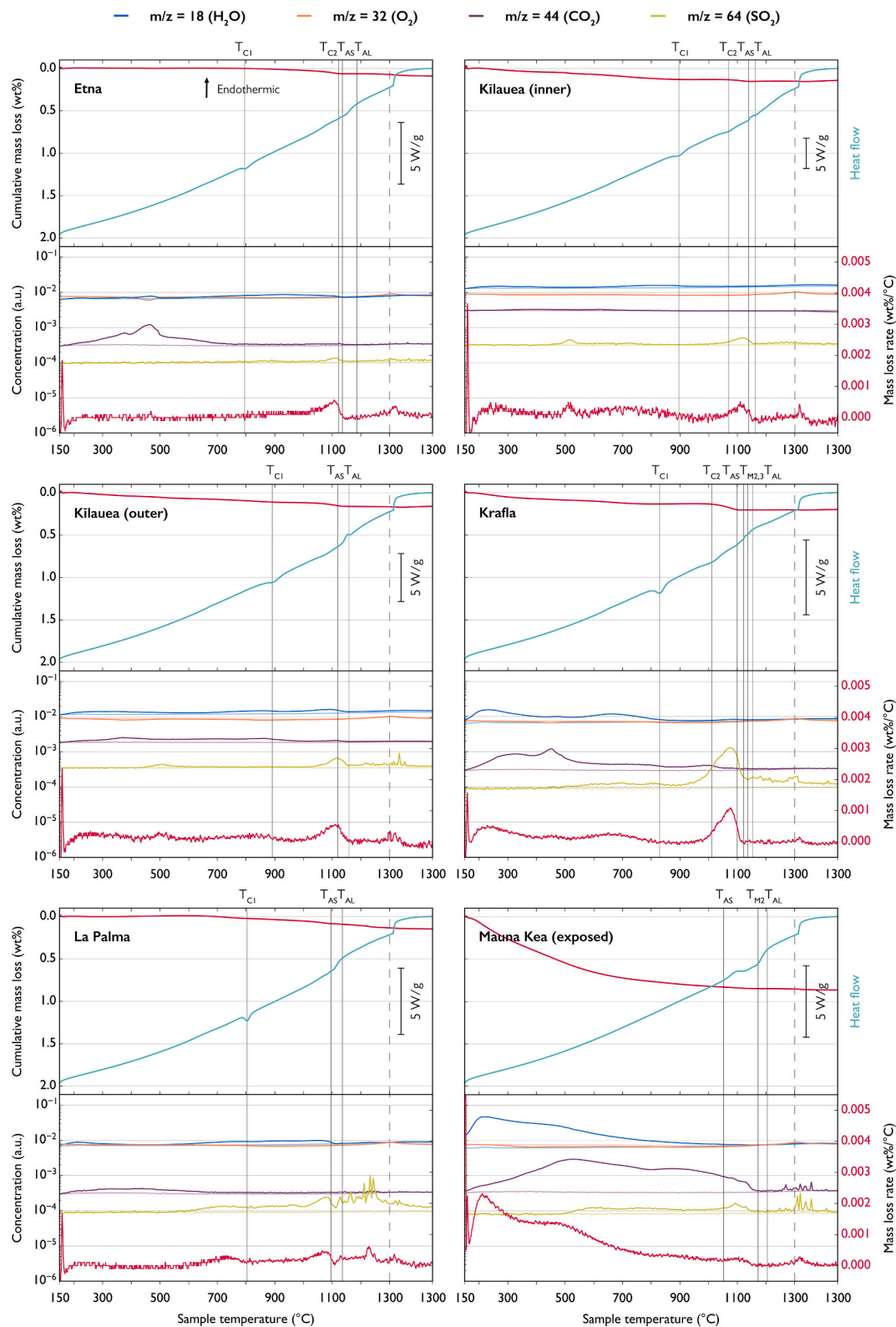
### 3.1.3. Processed analogue material

Two distinct crystallisation troughs are visible in the green glass synthesised under air (Fig. 5), dominated by the second at 848 °C. Although CO<sub>2</sub> is released below 800 °C, net mass loss is not seen until 1200 °C, after the onset of melting, in which O<sub>2</sub> is uniquely detected. This oxygen profile directly correlates with the independently measured mass loss rate curve. Water loss is negligible throughout (a slight baseline deviation is likely fed by an impurity from the purging gas itself), though a tiny expulsion is associated with crystallisation at T<sub>C2</sub>. Total mass loss stands at 0.102 %. By 1300 °C, melting appears incomplete.

Orange glass produced under reduced conditions meanwhile returns an overall net mass gain of 0.048 %, due to the steady slight oxidation of the sample from ~250 °C. The H<sub>2</sub>O profile matches that of the green glass. This higher-Ti composition is less stable after glass transition and crystallises earlier at 812 °C. A weak CO<sub>2</sub> signal extends until 1200 °C; the upturn ~1170 °C could provide an additional source of oxygen to drive a potential exothermic peak superimposed on the larger melting event, or alternatively indicate a double melting event in quick succession. T<sub>AL</sub> is reached within experimental range, by 1221 °C.



**Fig. 3.** Simultaneous DSC-TGA (upper subplot) and synchronous EGA (lower subplot) via mass spectroscopy performed on lunar mare simulants. TGA raw data is presented as cumulative mass loss normalised to the dried sample mass at 150 °C. DSC raw data is expressed as sample heat flow, with variable vertical scale. Mass loss rate is depicted on the lower plot for closer alignment with contributing gas species. EGA raw data collected as ion current is a proxy for relative gas concentration (arbitrary units) using the given mass-to-charge ( $m/z$ ) ratios. The fainter curves are baselines for their corresponding colour.  $T_{C1}$  and  $T_{Mx}$  represent characteristic temperatures associated with crystallisation peaks (if detectable) and onset of distinct melting events, respectively, whilst  $T_{AS}$  ( $= T_{M1}$ , i.e. first observable onset of melting) and  $T_{AL}$  are estimations of apparent solidus and liquidus temperatures as reported in Table 1. The subscript number 'x' denotes chronological occurrence.



**Fig. 4.** Simultaneous DSC-TGA (upper subplot) and synchronous EGA (lower subplot) via mass spectroscopy performed on mostly fresh volcanic deposits. Labels are as described for Fig. 3.

Partial oxygen reduction of the Mauna Kea ash lowers overall mass loss almost three-fold, predictably coinciding with the sub-1000 °C interval. However, from 800 °C the CO<sub>2</sub> signature is comparatively elevated and increased oxidation is observed in association, peaking just

after T<sub>AL</sub> and steadily growing throughout the isothermal segment. The prior reduction procedure lowers the onset of melting by 19 °C and apparent liquidus by 50 °C.



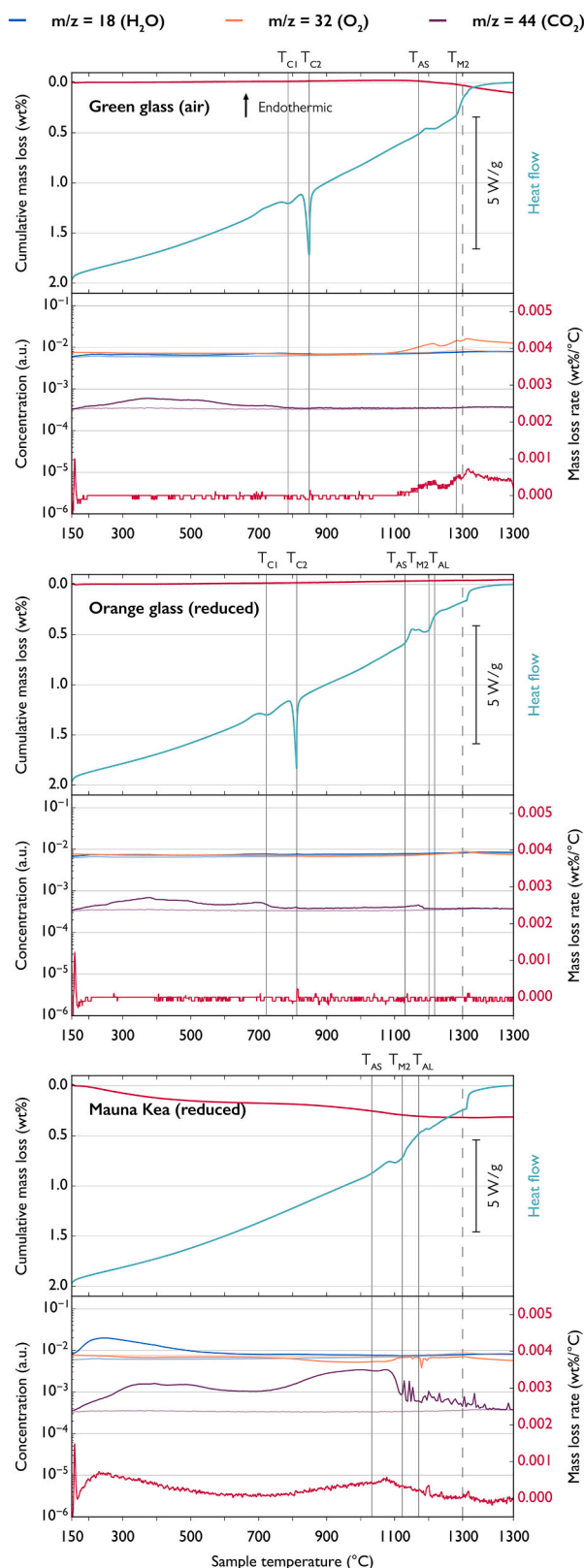


Fig. 5. Simultaneous DSC-TGA (upper subplot) and synchronous EGA (lower subplot) via mass spectroscopy performed on processed analogue material. Labels are as described for Fig. 3.

### 3.2. NPD

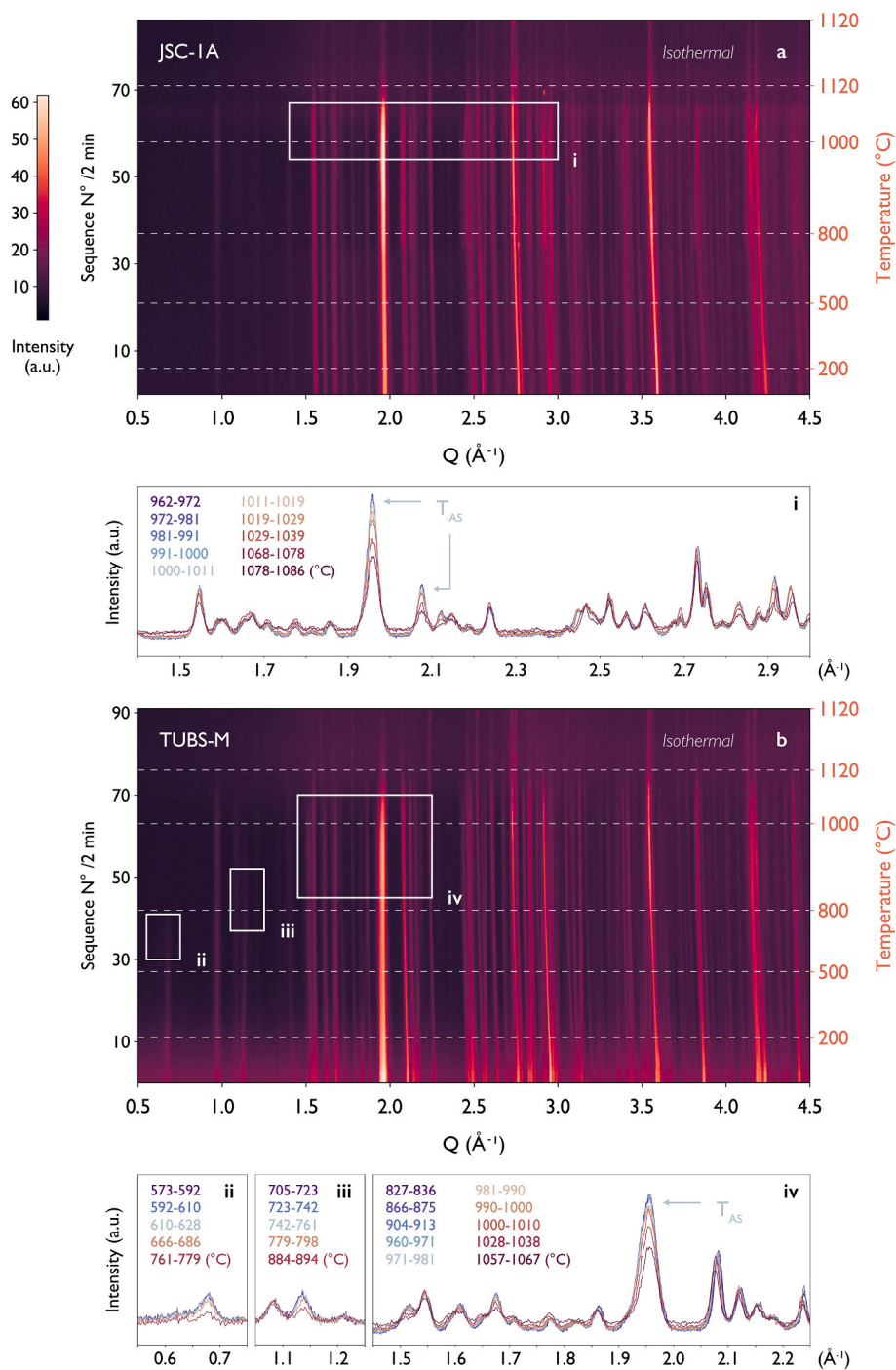
#### 3.2.1. JSC-1A

Vertically stacked sequential diffraction patterns are shown in Fig. 6a. Phase identification is assisted by reported components (e.g. NASA, 2007) and the International Centre for Diffraction Data (ICDD)'s search indexing software Sleave+, drawing upon the Powder Diffraction File database PDF-5+. Scarcity of low angle reflections point to the absence of layered structures with large unit cells that incorporate water molecules (e.g. zeolites and phyllosilicates including clays). Accordingly, the general leftward slant of peak traces with temperature indicates thermal lattice expansion. The faint reflection at  $1.30 \text{ \AA}^{-1}$  (plausibly trace alunite,  $\text{KAl}_3(\text{SO}_4)_2(\text{OH})_6$ ) is the sole observation of sub-solidus crystalline decomposition, becoming undetectable after around  $325^\circ\text{C}$  as dehydration in Fig. 3 begins to taper. Branching reflections (e.g.  $2.57$  and  $2.95 \text{ \AA}^{-1}$ ) from around  $300^\circ\text{C}$  are related to lattice distortion due to slight oxidation of  $\text{Fe}^{2+}$ , partially transforming ulvöspinel to ilmenite without detectable impact on the DSC profile. Drop off in intensity at  $3.19 \text{ \AA}^{-1}$  after  $600^\circ\text{C}$  could be attributed to loss of the long-range magnetic contribution from magnetite above its Curie temperature. Peak strengthening approaching  $800^\circ\text{C}$  is through crystallisation of the initial glass component. No evidence of any crystalline phase contributing to the elevated  $\text{CO}_2$  signal around  $600\text{--}750^\circ\text{C}$  is observed; given the cumulative mass loss of this reaction is less than  $0.1 \text{ wt\%}$ , water-bearing phases encompassed, the associated mineral breakdown is unsurprisingly below neutron detection limit ( $\sim 1 \text{ wt\%}$ ). The recognition of alunite, though in proportions beyond general detection, is possible with these counting statistics only through enhanced H sensitivity. Melt onset is first observed in the strongest plagioclase (bytownite) reflection centred on  $1.96 \text{ \AA}^{-1}$ , as well as the prominent pyroxene (augite/diopside) peak at  $2.08 \text{ \AA}^{-1}$  during the count ending  $1011^\circ\text{C}$  (inset i). The progressive decrease in crystalline intensity is accompanied by a subtle rise in the broad, diffusely scattered baseline, indicating loss of long-range order and the development of an amorphous component. Having adjusted for the systematic underestimation of the thermocouple, the onset of melting is lower than the DSC-derived value by ca.  $40^\circ\text{C}$ .

#### 3.2.2. TUBS-M

The release of adsorbed and interstitial water largely below  $400^\circ\text{C}$  is marked by a broad decrease in the diffuse baseline intensity (Fig. 6b). Subsequently, two distinct structural decompositions are identified (inset ii and iii). In the count ending  $628^\circ\text{C}$  through to  $779^\circ\text{C}$ , dehydrated chabazite ( $(\text{Ca},\text{Na})\text{Al}_2\text{Si}_4\text{O}_{12}\cdot 6\text{H}_2\text{O}$ ) is a candidate for the decaying low-angle peak at  $0.68 \text{ \AA}^{-1}$ , with an additional zeolite, analcime ( $\text{NaAlSi}_2\text{O}_6\cdot \text{H}_2\text{O}$ ), at  $1.13 \text{ \AA}^{-1}$  from  $761^\circ\text{C}$  until loss of detection at  $942^\circ\text{C}$ . Contrary to the generally observed lattice expansion, partial collapse of structural voids following zeolitic water loss leads to a reduction in unit cell volume. This contraction is more pronounced for analcime owing to constrained water positions and a rigid framework that cannot compensate as flexibly. With relatively open channels, chabazite releases water more readily. No potential carbonate phases nor sulfur-bearing candidates are observed within detection limits. Melt onset is detected as early as  $981^\circ\text{C}$ , again first in the  $1.96 \text{ \AA}^{-1}$  plagioclase peak (inset iv), well below the  $1161^\circ\text{C}$  predicted by DSC (reproducible to  $\pm 5^\circ\text{C}$ ).

Suspecting that the increased water content of TUBS-M relative to JSC-1A is responsible for the more pronounced solidus reduction, the protocol was repeated with TUBS-M dehydrated at  $500^\circ\text{C}$  in air for 24 h prior. Fig. 7a illustrates the differences at select temperature increments. Under ambient conditions, the diffuse water contribution is evident from the vertical separation of the two curves. Peaks at  $0.68$  and  $1.45 \text{ \AA}^{-1}$  (chabazite) are barely visible in the partially dehydrated sample – prolonged exposure at  $500^\circ\text{C}$  is sufficient to approach thermodynamic equilibrium, contrasted with inset ii in Fig. 6. The statistics of the repeat pattern appear weaker, although the trace is entirely reproducible. At

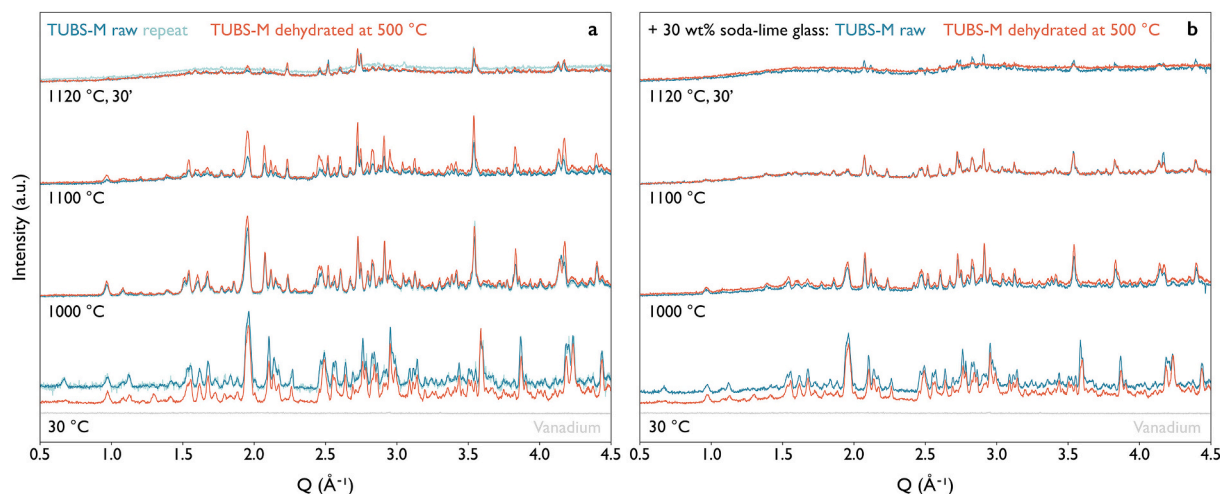


**Fig. 6.** In situ NPD of (a) JSC-1A and (b) TUBS-M collected at D20 (ILL). Successive patterns integrated over two-minute intervals with continuous temperature ramping (10 °C/min < 820 °C; 5 °C/min 820–1120 °C) are vertically stacked. Crystalline decomposition and melt onset ( $T_{AS}$ ) are highlighted in inset i-iv; not all temperature intervals contained in the outlined range are displayed and vertical scale varies between plots.

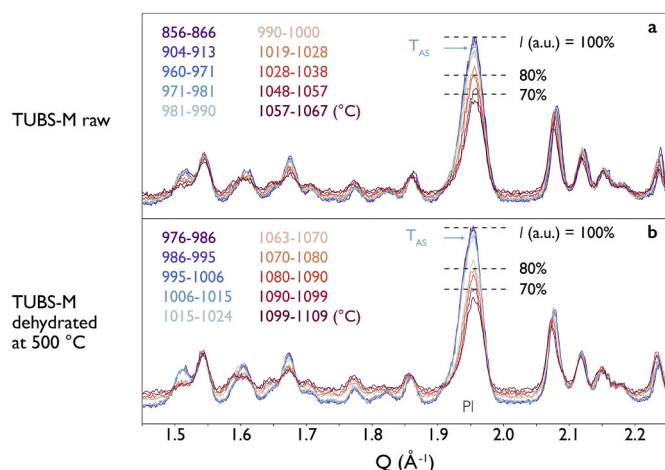
1000 °C both samples are totally dehydrated and raw TUBS-M has just begun to melt. As the incoherent baseline contributions converge, no significant difference in hydrogen content remains between the two, i.e. excess trapped intergranular water vapour is not discernibly present in the raw sample. There is a marked increase in melt extent in the raw sample by 1100 °C that therefore cannot be explained by ongoing water presence. However, after holding for a further 30 min the pre-treated sample is largely able to catch up. When melting is kinetically constrained, even beyond the water evacuation range, the dehydration history appears to influence melt onset and development. Furthermore, with a residual 0.169 wt% loss between 500 and 1000 °C (more than

double JSC-1A), the apparent solidus of the partially-dehydrated sample (1015 °C) remains almost 150 °C below the DSC-determined onset. To indicate potential differences in the progression of partial melting, Fig. 8 additionally compares the temperatures at which the first-melting plagioclase reflection surpasses 80 % and 70 % of its maximum intensity. The temperature discrepancy between samples in both cases is consistent at ca. 42 °C, marginally higher than for  $T_{AS}$  (34 °C).

This picture is complicated when increasing the starting amorphous fraction by 30 wt% through incorporation of soda-lime glass to promote earlier melting (Fig. 7b). Absolute peak intensities at ambient temperature are diminished by ~30 % as expected. Both samples undergo



**Fig. 7.** Select in situ NPD patterns contrasting the evolution of as-received (raw) simulant and after pre-treatment at 500 °C in air for 24 h. (a) Fully crystalline starting samples. (b) Additionally incorporating 30 wt% soda-lime glass (45–90  $\mu\text{m}$  bead size), following any pre-treatment.



**Fig. 8.** In situ NPD of (a) TUBS-M and (b) TUBS-M pre-dehydrated at 500 °C in air for 24 h. Dashed lines highlight melting onset and progression (exceeding 80 % and 70 % maximum intensity) in the first-melting peak at 1.96  $\text{\AA}^{-1}$  (plagioclase).

increased though comparable melting through 1000 and 1100 °C. Unlike the initial fully crystalline samples, patterns diverge after holding at 1120 °C; the raw sample evolves little over the 30-min increment, whereas the partially dehydrated version is nearly completely molten. Thus, the addition of soda-lime glass beads appears to preferentially encourage melting of the partially dehydrated crystalline component, whereas the raw pattern remains relatively unchanged in melt extent at 1120 °C aside from the proportional reduction in crystalline intensity.

## 4. Discussion

### 4.1. Sample trends

Simulants exhibit higher mass loss and volatile release at sub-solidus temperatures than fresh volcanic materials and synthesised glasses, though overlap exists between JSC-1A and unexposed ash. Among the simulants, over an order of magnitude difference in cumulative mass loss is seen (Fig. 2) in both 150–500 °C (between JSC-1A and TUBS-M) and 500–1000 °C (between JSC-1A and EAC-1A/LMS-1) ranges. Conversely, they exhibit comparable or in many cases slightly lower mass loss from 1000 to 1300 °C than the fresh material. For all 15

samples, degassing persists from room temperature until at least the apparent liquidus, although mass loss is not perceptible throughout.

The relative magnitude of species contributions is largely consistent among simulants, however, the loss rate is not always dominated by water (notably with material from the San Francisco Volcano Field). This distribution is more in line with the fresh samples, whose water contribution is small – or negligible in the case of Etna and La Palma ash. For these five rapidly collected deposits, mass loss rate is shaped by  $\text{SO}_2$ . The profile for synthesised glasses is driven by uptake or release of oxygen, depending on the fugacity ( $f_{\text{O}_2}$ ) during formation.

### 4.2. Mass loss origins

Here we combine insights from STA and NPD to briefly attribute mass loss in the samples. Sequential detection of  $\text{H}_2\text{O}$ ,  $\text{CO}_2$ , and  $\text{SO}_2$  is geologically consistent with dehydration, decarbonation, and desulfation (in their several forms) during decomposition reactions, diffusion from volcanic glass, exsolution from the melt phase, as well as the removal of fluid and melt inclusions. Below  $\sim 300$  °C, adsorbed, inter-layer, and caged water is liberated from secondary minerals, while crystalline ( $\text{H}_2\text{O}$ ) water may persist to  $\sim 600$  °C depending on the mineral phase. These develop in low-temperature environments through processes associated with weathering (e.g. clays), hydrothermal activity (e.g. zeolites), diagenesis (e.g. caliche), or direct precipitation from aqueous solutions (e.g. gypsum, hydromagnesite). Water evolution is sustained to 900 °C from the breakdown of more tightly-bound hydroxyl groups (e.g. amphiboles, micas). Early release of carbon-containing species is found in all cases except Kilauea (inner) from adsorbed  $\text{CO}_2$  or organic residues. Common carbonate weathering products can dissociate into metal oxides and  $\text{CO}_2$  from  $\sim 400$ –900 °C. Low temperature isolated  $\text{SO}_2$  detection (e.g. around 500 °C in the Kilauean reticulites) may be ascribed to fluid inclusions, which can be released over a broad temperature range, depending on the thermal stability of the host mineral. If volcanic glass is present, the majority of dissolved hydrogen, carbon, and sulfur species are released following the glass transition (typically 600–800 °C), as its structure relaxes and permits increased diffusion rates. Under inert atmosphere,  $\text{SO}_2$  in the range  $\sim 550$ –750 °C is a by-product of two-step sulfide decomposition, primarily pyrite ( $\text{FeS}_2$ ), provided oxygen is available from the surrounding phases or purging gas impurity (Hu et al., 2006). Non-lunar sulfates such as anhydrite ( $\text{CaSO}_4$ , the dehydrated form of gypsum) contribute to the signal up to  $\sim 1000$  °C. The reduced sulfur species  $\text{H}_2\text{S}$  was here below detection limit.

Approaching solidus temperatures, the (re-)emergence of  $\text{SO}_2$  loss in



all simulants (Fig. 3) and natural material (Fig. 4) can be attributed to either melt inclusions in primary crystals, or lattice-bound (structural or impurity) volatiles from softening and then partial melting. Stepwise heat flow indicates incongruent melting, particularly for simulants of more hybrid sources. Beyond the liquidus, punctuated signals from SO<sub>2</sub> and CO<sub>2</sub> containing-bubbles are indicative of low viscosity melt (on the order of 10<sup>1.12</sup> Pa·s, TUBS-M at 1300 °C (Stapperfend et al., 2024)). Their reception is not instantaneous (e.g. a lag in the material of Hawaiian origin), suggesting some time is required to allow volatile products to reach the melt surface for a given temperature. Exsolution of residual water is also observed therein (both reticulites), as its solubility in the complete melt steadily decreases. Where the gas traces of SO<sub>2</sub> and CO<sub>2</sub> align, for example between 900 and 1150 °C in OPRL2NT, the decomposition reactions of sulfates can be carbon-assisted, as highlighted by Petkov and Voecks (2023). While the major degassing period for simulants ends between 700 and 1000 °C (Fig. 3), for Etna and La Palma samples devolatilisation comparatively accelerates above 1000 °C, as these volatile sources dominate (Fig. 4).

Oxidation, when occurring, appears from 250 °C onwards. The signal is particularly enhanced during an isolated CO<sub>2</sub> event around 460 °C with OPRL2N, Etna, and Krafla ash. Reactions involving transient oxygen and granular surfaces are typically Fe<sup>2+</sup> to Fe<sup>3+</sup>, as seen with the ulvöspinel unmixing in JSC-1A, but can apply to glass as well. Besides iron, oxidising reactions may involve carbon (i.e. combustion), sulfur (e.g. pyrite), titanium (e.g. titanomagnetite, Fe<sup>2+</sup>(Fe<sup>3+</sup>,Ti)<sub>2</sub>O<sub>4</sub>), and chromium (e.g. chromite, FeCr<sub>2</sub>O<sub>4</sub>), all of which have been reported among the samples (e.g. Schrader et al., 2008). For strongly reduced samples (Kilauea, Krafla, orange glass, and processed Mauna Kea) degassing has tapered to the extent that oxidation drives net mass gain above 1300 °C. In the absence of sufficient O<sub>2</sub>, the molten inner reticulite also consumes CO<sub>2</sub>.

#### 4.3. Analogues in the lunar context

Direct measurements from Chang'e-5 (11–87 ppm H<sub>2</sub>O/OH particle size-dependent average (Lin et al., 2022, 2024)) and revisited Apollo samples (Liu et al., 2012) constrain typical bulk mare soil to sub-100 ppm H<sub>2</sub>O/OH, consistent with the background level from near-infrared spectra within ±35° latitude (Milliken and Li, 2017). Attributing the majority of mass loss under 1000 °C to water, simulants release two to three orders of magnitude (0.15–1.72 wt% from JSC-1A to EAC-1A) more, reduced by an order in the unexposed material (0.02–0.13 wt% from Etna to Krafla). This simplistic assumption discounts varying mass gain from oxygen uptake. Exceptions exist for areas of localised enhancement in various pyroclastic deposits and KREEP-rich material up to several hundred ppm water equivalent hydrogen; with a perspective towards attractive landing sites for ISRU, it is notable that the water mass of the Aristarchus Plateau pyroclastic deposit, for instance, has been calculated as a significant fraction of that estimated for the lunar poles (Lawrence et al., 2022). In contrast to the initial Apollo vacuum pyrolysis studies (Table S2), this study underestimates absolute volatile availability in the analogues as melting under atmospheric pressure prolongs solubility, particularly of water (Holtz et al., 1995).

Overlap in the carbon content range of volcanic melt from terrestrial and lunar sources is reported. Measured abundance in lunar mare soils and regolith breccias has been found to reach 216 ppm C (average 109–154 ppm among Apollo mare locations, concentrated in the finest fractions: up to 500 ppm < 53 µm (Moore et al., 1970; Gibson, 1977; Haskin and Warren, 1991)), though more than half (51–98 %, *n* = 6 basalts) of which is likely to be associated with terrestrial contaminants below 600 °C (Mortimer et al., 2015). If speciated entirely as CO<sub>2</sub>, this upper (contaminated) bulk value is equivalent to 792 ppm. However, CO is the primary exsolved C species on the Moon (Rutherford and Papale, 2009); across the realistic range of lunar *f*O<sub>2</sub>, the degassed CO<sub>2</sub> proportion is estimated between 4 and 29 % (Varnam et al., 2024). Water-

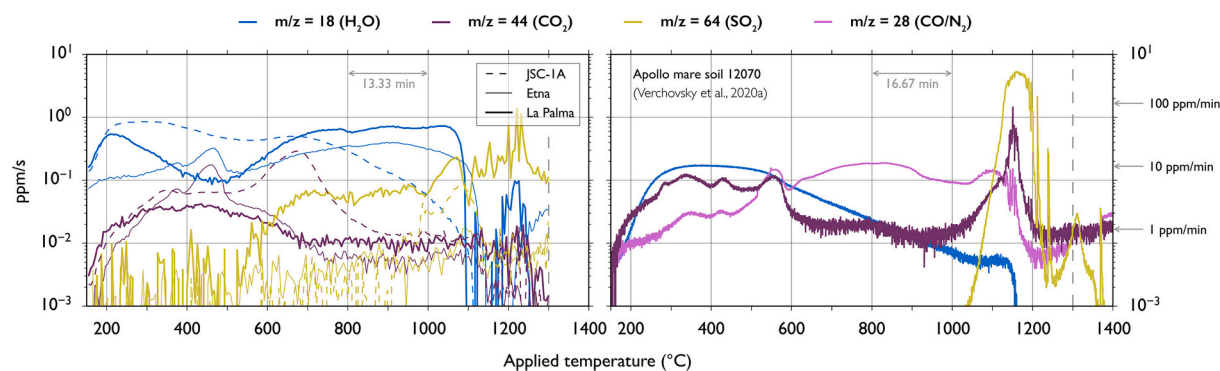
poor subaerial ocean island basalts (OIB) like Kilauea typically retain up to 350 ppm CO<sub>2</sub> (95 ppm C, excluding CO contribution) in quenched glass, though highly degassed products in low-viscosity systems characteristic of OIB may contain as little as <10 ppm CO<sub>2</sub> (Anderson and Brown, 1993; Johnson et al., 1994). Under standard oxidising conditions during eruption, most carbon in the system preferentially exists as CO<sub>2</sub> (Ni and Keppler, 2013), which remains more soluble than CO at high temperature by integrating into the melt via carbonate complexes. Supposing equal C abundance, these diverging redox and solubility behaviours differentiate lunar and terrestrial degassing profiles. Furthermore, whereas CO readily diffuses to escape from silicate melts, lower-mobility CO<sub>2</sub> tends to accumulate (Koch and Schmidt, 2023) until saturation is reached, so that larger gas voids may be expected.

Sulfur is locally more heterogeneous than hydrogen or carbon on the Moon. It is comparable or even enriched in lunar basalts relative to their terrestrial counterpart (up to 0.30 wt% versus less than 0.15 wt% (Taylor et al., 1991)), with 440–1400 ppm S (mission-dependent average 748–1280 ppm) identified in mare or mixed terrain Apollo soil (Gibson, 1977), primarily as troilite (FeS). In Earth-based basaltic melts, the two dominant forms are sulfide (S<sup>2-</sup>)- and sulfate (SO<sub>4</sub><sup>2-</sup>). Their ratio is strongly dependent on *f*O<sub>2</sub> over an interval (Δnickel-nickel oxide (NNO) buffer = −0.8 to +1.2, Wallace and Carmichael, 1992) encompassing redox conditions of potential analogue sources, where preferred sulfur speciation will vary accordingly. Whilst mare simulants tend to be derived from intraplate continental settings, fresher samples additionally include more oxidised environments i.e. OIBs (e.g. La Palma, Burton et al., 2023) and subduction-influenced intraplate basalts (e.g. Etna, Metrich and Clocchiatti, 1996). Sulfides are typically less thermally stable and degas at lower temperatures than sulfates (Gazulla et al., 2009) although oxygen fugacity influences the precise release pathways. Under lunar *f*O<sub>2</sub>, sulfur only occurs in its reduced state. When saturation is reached, reduced species including S<sub>2</sub>, H<sub>2</sub>S, COS, and CS<sub>2</sub> (ten Kate et al., 2010; Varnam et al., 2024) tend to escape more readily owing to their higher diffusivities, likely leading to more efficient release than in sulfate-rich melts. Outside of the volatiles considered in this study, N, F, and Cl species are additionally present in lunar soils up to a few hundred ppm (McCubbin et al., 2015).

The observed release patterns herein are broadly consistent (Fig. 9) with previous analyses on (to some extent contaminated) regolith (Gibson and Johnson, 1971; ten Kate et al., 2010; Mortimer et al., 2015; Verchovsky et al., 2020a), even if the origin and concentrations of volatiles diverges. Sustained H<sub>2</sub> or water release is derived from solar wind implantation, impactors, or indigenous OH in trace minerals like apatite, with a similar drop off around solidus. Agglutinitic and impact glasses, rather than minerals, have been shown to host the majority of H (Liu et al., 2012; Zhou et al., 2024). At latitudes below 30° (where surface temperatures exceed 400 K), the water content correlation with optical maturity diminishes (Li and Milliken, 2017). The integrated water abundance in Apollo soil 12070 (Table 2) is inflated above known values as the gas transfer was not optimised for water analysis. Below 600 °C, the carbon contribution is dominated by terrestrial sources. Sulfur is absent below 1000 °C. Exponential growth of carbon and sulfur species above ca. 1000 °C highlights a consistent degassing between simulants, fresh material, and original regolith: that of endogenous gases dissolved in the magma. Later-stage vesiculation arising from S<sub>x</sub>O<sub>y</sub>/CO<sub>x</sub>/N<sub>2</sub>/H<sub>2</sub>S exsolution therefore is also to be expected during processing of lunar regolith, a conclusion similarly drawn by Wilkerson et al. (2023). To our knowledge, there is no recorded lunar benchmark specifically for post-solidus volatile abundances in bulk regolith. Sulfur and carbon released above 1000 °C mainly originates from troilite (magmatic, or a product of vapour or fluid interaction)/carbide decomposition, volcanic glass, or redistributed sulfur/trapped exogenic organic material in vapour-condensed impact glass coatings or in mineral defects (McCubbin et al., 2015).

Although it is the most abundant volatile in lunar soils and volcanic glasses (Gibson, 1977; Saal et al., 2008), studies of lunar S have been





**Fig. 9.** Rate of release profiles for  $\text{H}_2\text{O}$ ,  $\text{CO}_2$ , and  $\text{SO}_2$  of select analogues (left) and Apollo mare soil 12070 (right). MS of the lunar sample (Verchovsky et al., 2020a; A. Verchovsky, personal communication, March 2025) was performed under a baseline pressure of  $10^{-3}$  mbar according to the setup in Verchovsky et al. (2020b). An additional reference curve at  $m/z = 28$  is therefore included; CO expectedly dominates this signal with a small contribution from  $\text{N}_2$ . Linear heating rates over the range are  $15^\circ\text{C}/\text{min}$  for the analogues, and  $12^\circ\text{C}/\text{min}$  for the Apollo sample. The release rate (ppm/s) is normalised to values in Table 2: for  $\text{H}_2\text{O}$  and  $\text{CO}_2$  total ppm corresponds to that stated for  $150\text{--}1300^\circ\text{C}$ , whereas  $\text{SO}_2$  uses the elevated estimation from  $T_{\text{AS}}\text{--}1300^\circ\text{C}$ .

limited compared to H, C, and N. Previous work indicates positive correlation of sulfur with  $\text{TiO}_2$  concentration, and negatively with  $\text{Al}_2\text{O}_3$ , in multiple sample types, including soils (e.g. Gibson et al., 1977; as compiled in Renggli et al., 2024). A positive correlation between sulfur content and  $\delta^{34}\text{S}$  exists convincingly in pyroclastic glasses, derived from magmatic processes and variations in source composition (Saal and Hauri, 2021). In lunar soils, a similar correlation is only loosely observed within individual source missions, linked with increasing maturity (Renggli et al., 2024). In contrast, impact glasses returned by Chang'e 5 show decreasing S abundance with increasing  $\delta^{34}\text{S}$  (Wang et al., 2024), attributed to isotropic fractionation during impact degassing, some of which may then condense on regolith grains (Dottin et al., 2022). Sulfur-rich material is preferentially concentrated in the fine grain, though with a tendency for S concentration to initially fall, before rising with decreasing grain size (Rees and Thode, 1974). Integrated  $\text{SO}_2$  abundance in soil 12070 exceeds the approximated concentrations in the La Palma ash sample and JSC-1A by one to two orders of magnitude, respectively. The difference in total S budget is significantly underestimated without the unmeasured reduced contribution (e.g., as presented in ten Kate et al., 2010). Terrestrial analogues are then more likely to resemble the post-solidus volatile budgets of lunar material sourced from lower-Ti, more felsic compositions, which, if targeted, would limit bubbling after melt onset. As temperatures exceed  $1200^\circ\text{C}$  in vacuum, degassing in both terrestrial and lunar volcanic material will encompass vaporisation of the constituent materials, fundamentally altering the composition of the remaining melt (Šeško et al., 2024).

Fresh volcanic material bridges the magnitude and distribution of documented mass loss in lunar soils, with as little as a few hundred ppm of water. 76 % of the total 0.090 wt% loss between  $150$  and  $1300^\circ\text{C}$  in Etna ash is above  $1000^\circ\text{C}$ , compared with ca. 91 % in the  $1000\text{--}1400^\circ\text{C}$  range for Apollo 11 and 12 soils (Gibson and Johnson, 1971). These fresh samples are highly glassy which is not only relatively more stable against chemical weathering (not applicable here due to rapid collection), but may initiate volatile release more readily above the glass transition temperature through enhanced mobility in the supercooled liquid state. By contrast, flash heating preferentially sequesters volatiles in quenched lunar impact melt (Liu et al., 2012). The high glass content and arid origin of JSC-1A source material is a strong determiner of its low mass loss, although minor substitution of basalt with anorthosite can dramatically raise this, i.e. 10 wt% in OPRL2N doubles mass loss below the solidus. Increasing volcanic glass fraction is not a realistic prospect; while lunar pyroclastic glasses have local pockets of abundance, these are a minor (typically  $<5$  vol% (McKay et al., 1991)) component of mare regolith. Elevated analogue glass content is not problematic per se as a sintering agent in processing applications, however, as largely unimitated impact glass (ca. 3–5 vol%), glass-

containing agglutinates (5–65 vol%, average 25–30 vol%), roapy glasses, and shocked minerals (maskelynite) cumulatively make up a comparable fraction of the lunar soil (McKay et al., 1991). Although JSC-1A exceeds the Etna ash water content almost twofold, its more tapered release above  $800^\circ\text{C}$  (Fig. 9) is potentially less influential for melt onset, provided the earlier evolved gases can be evacuated. The  $\text{H}_2\text{O}$  degassing rate nonetheless remains an order of magnitude higher than soil 12070 around  $T_{\text{AS}}$  ( $1051^\circ\text{C}$ ).

#### 4.4. Significance of volatiles for sintering and melting applications

We refer to approximations of the solidus and liquidus for specific conditions of pressure, rate, and  $f\text{O}_2$ , but they are not thermodynamically-determined temperatures, which would neither be applicable to ultimate lunar processing conditions. Partial and total gas pressures influence both reaction kinetics and the efficiency of volatile degassing itself (Zhang and Ni, 2010). Water vapour is known to enhance catalytic decomposition of carbonates (Giammaria and Lefferts, 2019) and sulfates (Gallagher et al., 1970). By way of example, caliche, a multi-constituent sediment primarily composed of calcium carbonate, has been proposed as the main  $\text{CO}_2$  source in JSC-1A (Street Jr. et al., 2010). Its  $\text{CO}_2$  trace peaks earlier than for isolated calcium carbonate decomposition, demonstrating the effectiveness of water (captured in the caliche matrix and alunite) to aid the breakdown of the  $\text{CaCO}_3$  structure via formation of intermediate species (e.g.  $\text{Ca}(\text{HCO}_3)_2$ ) by over  $100^\circ\text{C}$ . Even though water is no longer detectable by  $1000^\circ\text{C}$  in Fig. 7a, prior water loss appears to have facilitated earlier melting of the raw TUBS-M sample. Hydration-weakening of grain boundaries, e.g. via water-mediated diffusion or by enhancing reactivity, can lower the activation barrier to incipient melting, potentially initiating the first reduction in crystalline coherence. Once underway, the rate of intensity decrease initially appears consistent (Fig. 8), suggesting that the melting process itself is proceeding at a similar kinetic pace. Prolonged heating at  $1120^\circ\text{C}$  (Fig. 7a) allows the samples to approach equal melt fractions once water-related kinetic advantages are no longer decisive and residual resistances (e.g. structural rigidity, dry grain boundaries) are overcome. The much larger partial pressure within TUBS-M compared to JSC-1A overwhelms the contribution towards lowering solidus from a  $10^7$ -fold reduction in external atmospheric pressure (recognising trapped inter-granular air during capillary loading will locally reduce this figure), at moderately slower heating. A pressure gradient is expected towards the container opening due to gas migration (Figs. 10 and S2). In more hydrated systems, rapid degassing can steepen chemical potential gradients. If gases cannot escape efficiently, pressure builds up, and the resulting increase in volatile solubility subsequently suppresses further bulk degassing. In contrast, drier samples exhibit steadier degassing and

**Table 2**

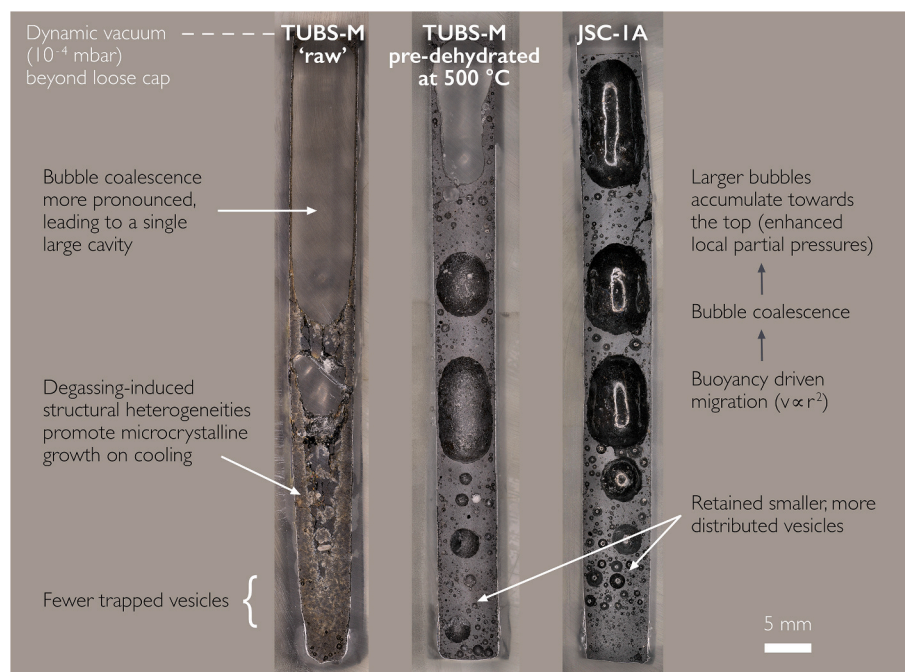
Integrated concentrations of H<sub>2</sub>O, CO<sub>2</sub>, and SO<sub>2</sub> in Apollo mare soil 12070 and select analogues. Apollo data are derived from Verchovsky et al., 2020a. In parenthesis are the contributions of CO<sub>2</sub> and SO<sub>2</sub> towards the overall C and S budgets. Ranges refer to applied (furnace) temperature, except T<sub>AS</sub>, which pertains specifically to sample temperature, consistent with values in Table 1. Integrated ppm is calculated by normalising the sum of the relative gas fractions (H<sub>2</sub>O, CO<sub>2</sub>, SO<sub>2</sub>, and O<sub>2</sub>) against the absolute measured mass loss fraction. For asterisked values, oxygen is neglected in the calculation as apparatus diffusion significantly contributes to the total. The increase in SO<sub>2</sub> and to a lesser extent CO<sub>2</sub> ppm over narrower intervals is a consequence of the calculation a) assuming all gases behave similarly across the full range, and b) normalising a greater relative share of gas concentration over a considerably smaller denominator at high temperature in water-dominated systems.

	Integrated ppm											
	Apollo mare soil 12070 (Verchovsky et al., 2020a)				JSC-1A simulant				Etna volcanic ash			
	150–1400 °C		1000–1300 °C		150–1300 °C		1000–1300 °C		150–1300 °C		1000–1300 °C	
m/z = 18, H <sub>2</sub> O	316 (0.03 wt%)	4	1564 (0.16 wt%)	47	2	906 (0.09 wt%)	474	78	1598 (0.16 wt%)	597	94	
m/z = 44, CO <sub>2</sub> (C)	324 (89)	124 (34)	244 (67)	59 (16)	67 (18)	111 (30)	18 (5)	12 (3)	80 (22)	28 (8)	49 (13)	
m/z = 64, SO <sub>2</sub> (S)	1533 (766)	1529 (765)	5 (3)	29 (15)	36 (18)	7 (4)	15 (8)	19 (10)	88 (44)	165 (83)	354 (177)	

maintain lower internal pressures, potentially enabling more efficient volatile escape.

Above the solidus, any dissolved water retained in the melt helps depolymerise the silicate network and enhance cation mobility, as well as accelerating diffusion of CO<sub>2</sub> (Koch and Schmidt, 2023). The effective reduction in viscosity of the melt phase encourages bubble migration and expansion, which are in turn pathways for further degassing. However, water-rich melts can dissolve and retain greater amounts of other volatiles before exsolution begins (Mysen, 1976). Formation of gas bubbles can also temporarily raise the effective local viscosity as the melt must deform around many small inclusions. Relatively wet melts, e. g. late release simulants like LMS-1 or EAC-1A, or Kilauea, may then follow less predictable viscosity evolution during melting due to the fluxing effect of volatiles destabilising the Si-O-Si framework. Meanwhile, drier systems depend on more gradual temperature-driven depolymerisation, with implications for material workability. Faster heating leads to oversaturation and rapid bubble nucleation and growth (e.g. Lim et al., 2023), as well as considerable broadening of evolved gas peaks (ten Kate et al., 2010), whereas slower rates permit steady-state conditions to be approached. Although the hydrogen and amorphous contents of both raw and pre-dehydrated TUBS-M converge upon holding at 1120 °C (Fig. 7a), the diverging textures in Fig. 10 indicate that degassing history still influences crystallisation kinetics. The raw sample sustained greater volatile loss during the heating programme, possibly creating chemical and structural heterogeneities that acted as nucleation triggers during cooling. These effects were most pronounced near the capillary edge, leading to the observed microcrystalline textures. In contrast, the pre-treated sample had already lost most of its water and reached structural equilibrium prior to high-temperature treatment, resulting in a more uniform melt structure and correspondingly homogenous texture on cooling. Further ex situ investigation (e.g. via SEM-EDS/EBSD or Raman) could offer more detailed insights into these cooling textures.

The mechanism by which vacuum conditions induce greater mass loss has been explained in terms of extended sulfur outgassing towards maximum temperatures due to earlier melt onset (Wilkerson et al., 2023), and more effective outward gas transfer (Lim et al., 2023) from heightened chemical potential gradient between the melt and its surroundings, thereby increasing diffusive flux. Even under so-called vacuum, an inevitable gradient of volatile vapour will locally emerge. These volatile by-products can transiently convect, aiding the heat transfer mechanism for furnace sintering and melting more significantly than for laser irradiation, where such contributions are proportionally less impactful under higher imposed temperatures. With complete melting, deviations in the target temperature are more forgiving. It remains to be determined at what level nominal vacuum becomes material-limited, and where negligible returns of effort are seen. Petkov and Voecks (2023) recommend no perceivable advantage beyond 10<sup>-6</sup> mbar with JSC-1A, to which we suggest better than 10<sup>-4</sup> mbar has diminished benefit for more heavily weathered untreated simulants like TUBS-M. Non-lunar oxidation is observed in a majority of cases from ca. 200 ppm apparatus contamination; although such levels would likely be absorbed into the already-saturated melt as part of the silicate network, higher concentrations (including via dissociation of volatiles) could affect liquidus temperatures if more refractory metal oxide crystals are stabilised under higher oxidation states. In an important distinction for volatile extraction efforts, O<sub>2</sub> release does not occur when the synthesised glass is formed under reducing conditions. We speculate this is not down to the difference in Ti content between the orange and green glasses. So far, no lunar-like analogues have been produced starting from oxides using a gas mixing furnace on account of practicality and volumes involved; gas release from samples produced under different redox states should be further studied, particularly in view of imitating accurate lunar C and S speciation. Simulating iron in the predominantly reduced state Fe<sup>2+</sup> would better replicate its role as a network modifier in lunar melts, leading to lower viscosity (Dingwell, 1991; Di Genova



**Fig. 10.** Sample sections (set in resin) following NPD to 1120 °C, with cooling at 10 °C/min. Fewer trapped vesicles towards the base of the relatively volatile-enriched TUBS-M is potentially due to more efficient local bubble migration at lower initial viscosity, with melt compaction from the expanding upper void displacing deeper bubbles upwards. Apparent porosities may be exaggerated following densification upon partial melting, in tandem with sample migration up container walls (Dominguez and Whitlow, 2019), creating a natural central cavity (Fig. S2).

et al., 2017; Casas et al., 2023), whilst enhancing the solubility (Brooker et al., 2001) and diffusivity (Persikov et al., 2010) of volatiles contained therein.

Given the considerable advancement in melt progression within the larger TUBS-M volumes ca. 200–300 °C beyond the bulk of water release, overhead percolation distance will affect the sintering window for a certain composition and heating rate. Agitation suggestions (Lim et al., 2023) may not be practical for uncontained structures. If partial melting is engaged to enhance consolidation, the temperature window will coincide with the primary – lunar-like – mass loss for fresh samples around 1100 °C, just before or after the apparent solidus. The post-solidus contribution is independent of material weathering and limited to <0.06 wt% (up to 1300 °C) across the investigated samples, which falls short of actual lunar concentrations, notably in S. Sulfur in lunar regolith has a relatively minor effect on melt onset as the reduced forms do not strongly interact with the silicate network. An exception may be near Fe-FeS eutectic temperatures in Fe-rich regions, where troilite and other sulfide phases could promote localised early melting. The primary relevance of additional sulfur in lunar melting processes is rather in degassing and redox balance. The preference for less mobile carbon and sulfur species in terrestrial material might therefore inadvertently replicate aspects of lunar volatile behaviour, such as increased bubble density, despite the inherently more efficient degassing style associated with lunar species like CO, H<sub>2</sub>S, or S<sub>2</sub>.

#### 4.5. Focusing pre-treatment

Efforts to reduce the bulk (up to 700–1000 °C) of volatiles in simulants are particularly encouraged for larger-scale studies due to primarily water-driven reaction influences, recognising freshness cannot be easily ‘corrected’ (i.e. processed Mauna Kea ash versus Kilauea) and introduces new deviations, e.g. grain consolidation (Wilkerson et al., 2023) and residual oxides. Moreover, it is suspected that the surprising result in Fig. 7b arises because of structural defects or reactive sites created during partial dehydration that interact more effectively with the subsequently added soda-lime glass, leading to more extensive

melting than when combined with the raw ‘wetter’ simulant. This outcome may be an unavoidable consequence of pre-treatment of partially glassy analogue material. Whilst 500 °C is sufficient to drive off the majority (81 %) of volatiles in TUBS-M without discernibly affecting particle size distribution, it would liberate less than half of that (38 %) in EAC-1A. The translational fluxing effect of this remnant water likely obscured the influence of glass content in the comparison of sintering temperatures between EAC-1A and JSC-1A in Ginés-Palomares et al. (2023).

A conscious pre-treatment of any designated lunar analogue should consider the envisaged thermal strategy, as well as sample individuality. For furnace heating where particular viscosities may be targeted, long and elevated pre-heating is advised. Occasional sample turnover would help liberate exsolved volatiles at depths beyond a few cm. Volume-appropriate time scales should be further investigated. Exceeding  $T_C$  is not critical as temperatures are too low for stable consolidation and would be experimentally surpassed. Sample loading would benefit from being performed within a vacuum chamber, if practical, to limit trapped intergranular air. Similarly, if the aim is to evaluate the cooled product of rapid (laser, microwave) beam heating, this approach would reduce porosity from terrestrial contamination at the expense of particle agglomeration. However, pre-treatment temperatures should be more conservative (below sintering) under this setup when investigating melting parameter refinement, whereupon strict grain distribution is sought. Therefore, although naturally-sourced terrestrial material cannot be an effective proxy in replicating complete thermal treatment of lunar regolith, simulants and fresher samples are best applied to isolate the effects of particular processing stages.

## 5. Conclusions

Regolith from the ancient mare basins is better approximated by younger volcanic material uncompromised by terrestrial alteration. With the exception of JSC-1A, the simulants catalogued here exhibit higher volatile degassing than freshly collected samples below solidus temperatures. Differences with the lunar regolith can be divided by this



threshold. Excess terrestrial H<sub>2</sub>O, and to a lesser extent CO<sub>2</sub>, act as depressants of melt onset. The fluxing effect of these partial pressures and degassing (in)efficiency is exacerbated at scale from impeded gas percolation; heating timescales are relevant for all pre-treatment and melting strategies and require further exploration at different sample volumes. When water and carbonate content is minimised, analogue fidelity is limited by the volatiles inherent to the melt. Above the solidus, elevated lunar sulfur and carbon contents primarily promote more abundant and efficient bubble formation. Unless fresher sources become widely distributed, sintering and melting experiments with different simulants compositions will rely on more deliberate pre-treatment strategies at typically higher temperatures than previously conducted. Temperature ranges and features reported herein to assist this targeting relate to a dry nitrogen atmosphere at a given rate. Using analogue or even returned samples to predict degassing behaviour under lunar surface conditions is challenged by the reproducibility of lunar vacuum, low gravity, and strongly reducing conditions. These environmental factors could further influence degassing kinetics, melt viscosity, bubble coalescence/migration, and crystallisation pathways. Going forward, it is not unreasonable that smaller quantities of oxides could be manipulated under controlled redox states to synthesise more realistic-performing crystalline material for validating melt processing strategies.

### CRedit authorship contribution statement

**Jennifer Sutherland:** Writing – original draft, Visualization, Validation, Software, Methodology, Investigation, Formal analysis, Conceptualization. **Kai-Uwe Hess:** Writing – review & editing, Validation, Supervision, Resources, Methodology, Investigation, Formal analysis, Conceptualization. **Thilo Bissbort:** Writing – review & editing, Formal analysis. **Simon Stapperfend:** Writing – review & editing, Resources, Investigation. **Axel Müller:** Supervision, Project administration, Funding acquisition. **Thomas C. Hansen:** Writing – review & editing, Supervision, Methodology, Investigation, Data curation. **Alexander Rack:** Investigation, Data curation. **Donald B. Dingwell:** Writing – review & editing, Supervision, Resources, Project administration, Funding acquisition, Conceptualization.

### Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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### Appendix A. Supplementary data

Supplementary data to this article can be found online at <https://doi.org/10.1016/j.chemgeo.2025.123115>.

### Data availability

Data will be made available on request.

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