



4-D imaging of sub-second dynamics in pore-scale processes using real-time synchrotron X-ray tomography

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Received: 1 March 2016 – Published in Solid Earth Discuss.: 21 March 2016

Revised: 1 June 2016 – Accepted: 7 June 2016 – Published: 15 July 2016

Abstract. A variable volume flow cell has been integrated with state-of-the-art ultra-high-speed synchrotron X-ray tomography imaging. The combination allows the first real-time (sub-second) capture of dynamic pore (micron)-scale fluid transport processes in 4-D (3-D + time). With 3-D data volumes acquired at up to 20 Hz, we perform in situ experiments that capture high-frequency pore-scale dynamics in 5–25 mm diameter samples with voxel (3-D equivalent of a pixel) resolutions of 2.5 to 3.8 μm . The data are free from motion artefacts and can be spatially registered or collected in the same orientation, making them suitable for detailed quantitative analysis of the dynamic fluid distribution pathways and processes. The methods presented here are capable of capturing a wide range of high-frequency nonequilibrium pore-scale processes including wetting, dilution, mixing, and reaction phenomena, without sacrificing significant spatial resolution. As well as fast streaming (continuous acquisition) at 20 Hz, they also allow larger-scale and longer-term experimental runs to be sampled intermittently at lower frequency (time-lapse imaging), benefiting from fast image acquisition rates to prevent motion blur in highly dynamic systems. This marks a major technical breakthrough for quantification of high-frequency pore-scale processes: processes that are critical for developing and validating more accurate multiscale flow models through spatially and temporally heterogeneous pore networks.

1 Introduction

Porosity, permeability, and flow in geological systems are all highly dynamic. Changes in the confining conditions (e.g. pressure, temperature), flow volume, fluid chemistry/viscosity, or suspension composition can drive mass transport (through processes such as precipitation, dissolution, deposition, or erosion) and change the connectivity and tortuosity of the pore network. The evolving porosity and permeability then cause further changes in both micro- and macro-scale flow. These mass transport processes and the passage of fluid–fluid and fluid–rock reaction fronts through heterogeneous geological systems occur at the pore scale. Developing a detailed appreciation of macroscopic processes as diverse as pollutant transport, hydrocarbon recovery, CO₂ sequestration, storage of nuclear waste, aquifer management, nutrient bio-accessibility, building stone preservation, and hydrothermal deposit formation requires a quantitative understanding of the multiscale effect of pore-scale processes.

Synchrotron imaging for dynamic geoscience applications

Laboratory and synchrotron X-ray computed micro tomography (usually XCT and sXCT respectively) are ideal approaches for imaging these processes as they allow in situ observation in a high-resolution, non-destructive way. The general principles of X-ray tomography, reconstruction, and data processing as related to geoscience applications are covered

in a series of recent reviews (Bultreys et al., 2016; Cnudde and Boone, 2013; Hess et al., 2011; Maire and Withers, 2014; Wildenschild and Sheppard, 2013). For the specific application of understanding pore-scale processes, XCT and sXCT are now becoming widely used for both qualitative and quantitative imaging of complex natural pore networks and the distribution of liquid(s) within them (Al-Raoush et al., 2011; Al-Raoush and Willson, 2005; Berg et al., 2013; Bhreasail et al., 2012; Boone et al., 2014; Bultreys et al., 2015a; Cnudde and Boone, 2013; Dewanckele et al., 2012; Geraud et al., 2003; Herring et al., 2013; Iglauer et al., 2011; Katuwal et al., 2015; Ma et al., 2016; Naveed et al., 2013b; Olafuyi et al., 2010; Sakellariou et al., 2003; Sok et al., 2010; Wildenschild et al., 2002; Wildenschild and Sheppard, 2013). The data are also being used as a basis of, and validation, for numerical simulations (Al-Raoush and Papadopoulos, 2010; Alhashmi et al., 2015; Bultreys et al., 2015b, 2016; Degruyter et al., 2010; Menke et al., 2015; Naveed et al., 2013a; Raeini et al., 2014, 2015; Fourie et al., 2007).

Standard operation is to collect a set of 2-D “projections” or “radiographs” at constant angular spacing (hereafter angular density) while the sample is rotated through 180° or 360° . Standard 3-D tomographic image datasets require seconds (synchrotron), minutes (synchrotron and laboratory), or even hours (laboratory) to acquire (Maire and Withers, 2014). Until recently this has limited 3-D experimental investigations of dynamic processes because the critical pore-scale processes occur over much shorter durations. Many key fluid–rock and fluid–fluid interactions therefore remain poorly constrained. Furthermore, while numerical simulations can now consider multiple mass and thermal transport processes simultaneously and can incorporate realistic pore geometries (see Bultreys et al., 2016, for in depth review), many generally still lack validation from experimental data from natural systems where observations are needed with a range of temporal and spatial resolutions.

In recent years, careful alignment of 3-D datasets collected at fixed time points over moderate duration experiments has enabled “time-integrated” tomography as a tool for quantification of fluid dynamics and porosity evolution (Andrew et al., 2014, 2015; Armstrong et al., 2014a; Berg et al., 2013; Blunt et al., 2013; Herring et al., 2013, 2014; Lin et al., 2016; Menke et al., 2015; Wildenschild and Sheppard, 2013). The increased image acquisition rates now available at third-generation synchrotron facilities have driven development of faster “continuous” imaging. In this mode, sample rotation is not stopped for each projection; instead data are acquired over a narrow arc (typically 0.1 – 0.5°). For most geological materials, collection of the projection data can be just a few seconds (Andrew et al., 2015; Berg et al., 2013; Bhreasail et al., 2012; Pistone et al., 2013, 2015; Youssef et al., 2014). However, 2-D imaging (the same projection acquisition rates but without rotation) is still needed to observe processes occurring on a timescales of milliseconds to a few seconds, such as individual Haynes jumps and the sub-

sequent relaxation dynamics (Armstrong et al., 2014b; Berg et al., 2013). Two-dimensional imaging cannot capture pore and pore–fluid interface morphology, pore throat orientation and size, or the location of the events within the fluid volume.

Advances in temporal resolution through camera and beamline technology have been followed by highly parallelised iterative reconstruction methods which require far fewer projections yet still achieve adequate reconstructed image quality (Batenburg and Sijbers, 2011; Brabant et al., 2014; Kaestner et al., 2015; Kazantsev et al., 2015a, b; Van Eyndhoven et al., 2015). However, direct in situ observation of sub-second events and processes, as well as those that cause rapid change over longer durations, still remain challenging. Pushing time-integrated tomography towards true long duration 4-D imaging would therefore have high impact on our understanding of pore-scale processes.

Here we present a state-of-the-art 4-D (3-D + real time) imaging methodology that enables visualisation and quantitative assessment of dynamic pore-scale processes in real-time (3-D acquisition rates of up to 20 Hz) over variable experimental durations. We present two key methodology developments: sub-second real-time imaging of fluid transport up to an order of magnitude faster than previously presented and the extension of that imaging protocol to allow slower or variable flow dynamics to be quantitatively assessed over longer time periods. We illustrate the potential of these developments for providing insight into pore-scale processes using three case studies showing preliminary data from experiments that quantify initial wetting, the evolution from dynamic to steady state flow, and the evolution of the reaction front during chemical dilution or fluid mixing.

2 Experimental set-up

Disaggregated sandstone gravels (1–2 or 2–4 mm size fraction) were loaded into a reusable cylindrical gravity-fed flow cell and mounted on the JEEP i12 beamline (Diamond Light Source, Harwell, UK) (Fig. 1). Full beam specification can be found in Drakopoulos et al. (2015). The cell diameter for this system can be varied (up to 50 mm diameter) according to experimental need and X-ray transmission through the sample. Perspex sample chambers of 6 and 25 mm internal diameter were used here. A circular piece of fine mesh was placed at the bottom of the cell above the outlet feed to prevent fine material being washed through to the outflow reservoir (maximum volume 125 mL in the configuration used). Particles were added to a minimum depth of 15 mm to ensure no effects from basal flow were observed in the imaging field of view. Particle diameters were sufficiently small that the influence of the cell geometry on particle packing was minimal. The random nature of the packing structure was checked before each experiment through the collection of a single tomography dataset (Fig. 2). In the well-cemented host sandstones (prior to disaggregation) grain sizes ranged from 200

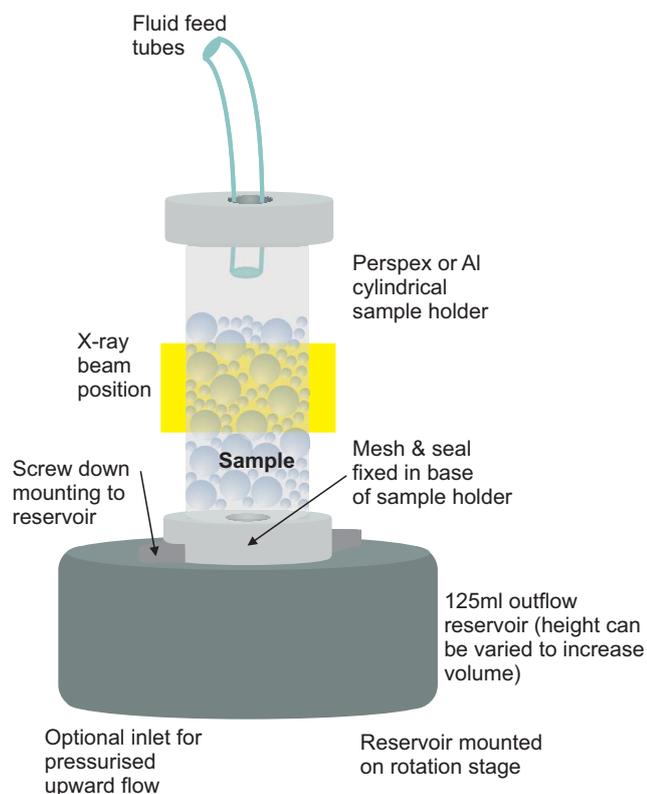


Figure 1. A schematic showing the simple peristaltic pump drip-fed gravity-driven flow cell. The system was designed to be modifiable to incorporate laminar and multiphase flow and to be suitable for use with confining pressure and variable temperatures. All key dimension (sample cell diameter, height, and material; reservoir volume; feed) can be adjusted to fit experimental parameters and beamline imaging conditions.

to 500 μm , with pore diameters of 10–150 μm . Mercury and helium porosimetry and water absorption measurements on the disaggregated samples were 13–18, 15–23, and $\sim 19\%$ respectively, and the measured permeabilities ranged from 250 to 600 md. In the configuration used here, fluid was supplied at a constant flow rate by a 12-channel peristaltic pump operated from the beamline control room (outside the X-ray enclosure or “hutch”) so that flow rates and compositions could be adjusted in real time, i.e. while 3-D data are being acquired. The inlet tube was attached centrally to the top of the cell using a slip ring to ensure a fixed drip point (approx. 5 mm above the centre of the rotating sample). Flow rates were varied from 40 to 200 $\mu\text{L min}^{-1}$, which corresponded to a drip every 2 and 10 s. The majority of experimental runs were performed using a 6 M KI solution to ensure high contrast between air, fluid, and rock particles. Pilot scans showed un-doped H_2O and some oils could be observed under the hard X-rays of the JEEP white beam (50–150 kV), although low contrast would make post-processing challenging.

The flow cell was designed to be easily modified and can accommodate base fed laminar flows, simultaneous flow

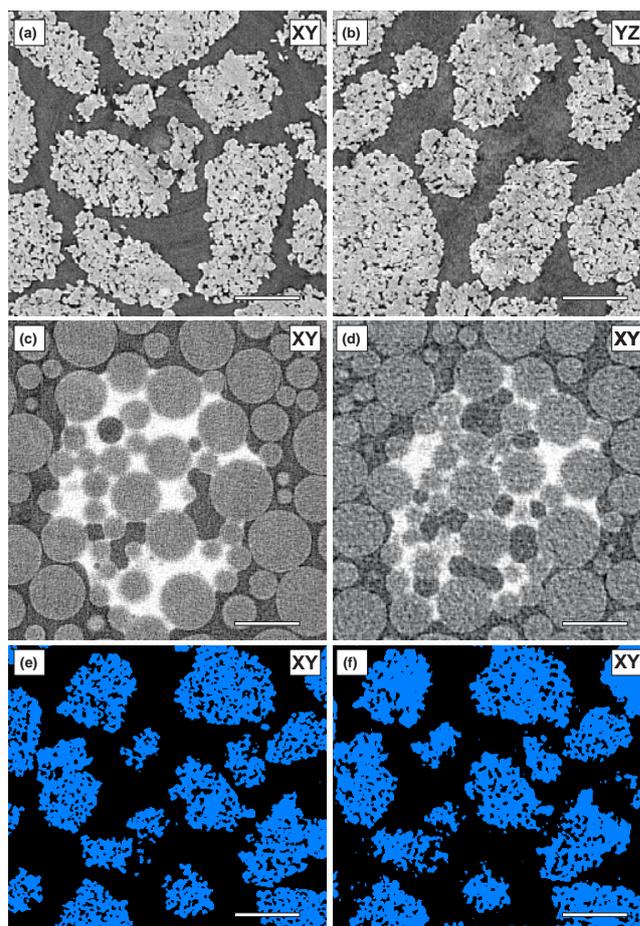


Figure 2. Two-dimensional reconstructed slices through a typical gravel sample. High-resolution (1800 projection data) perpendicular (a; XY) and parallel to flow direction (b; YZ) reconstructed using filtered back projection. (c) A 2-D slice (XY) through a simple spherical bead pack containing a small volume for KI solution reconstructed from 360 projections (2°) using filtered back projection; (d) the same 2-D slice reconstructed from 180 projections (1°) using filtered back projection; (e, f) comparison of the segmentation of the rock possible from the 1800 and 180 projection data. All scale bars are 200 μm .

of multiple fluids (variable composition or temperature), or multiphase suspensions. Thin-walled (1 mm) aluminium sample chambers can be used for both variable thermal and pressure experiments. The cell can be loaded with disaggregated material (e.g. soil, sand, gravel, glass beads) or with more coherent cored soil/rock samples of any diameter. However, when designing ultra-high-speed experiments or processes that cause subtle changes in attenuation/composition, some higher density or large volume solid samples will increase the minimum exposure time and may also cause artefacts that reduce the ability to resolve features of interest.

3 Ultra-fast image acquisition

The maximum rotation speed of the JEEP tomography stage is 10 Hz (continuous rotation). As synchrotron tomography datasets are acquired over 180°, 3-D volumes can be acquired at up to 20 Hz. The actual operational rotation speed for the given experiment will be controlled by the exposure time for each projection and the total number of projections collected per 3-D dataset. Exposure time is controlled by X-ray flux and sample density. Imaging at the JEEP beam-line can be performed using a polychromatic (white, higher flux) or monochromatic X-ray beam (selectable energy range 53–150 kV). Images were acquired using a Vision Research Phantom Miro™ 310 M camera (up to 5000 images per second). In monochromatic beam mode, the exposure time for the large diameter cell was 200 μ s (3.8 μ m pixel resolution), which dropped to 90 μ s for the smaller diameter system (2.5 μ m pixel resolution). The system can be in continuous rotation at this speed, in one direction, for as long as required. Beam hardening can occur under white beam in response to the energy dependence of the attenuation, but none was seen in our test data. While white beam would allow shorter exposure times (between 30 and 50 % reduction in the test scans), there would be lower contrast between the low density phases and more X-ray scattering and associated noise in the images. Imaging in monochromatic beam gave increased contrast between air and water in a three phase system, without requiring long exposures. The diameter of the samples used here exceeds the field of view of the highest magnification module available at JEEP (1.25 μ m pixel resolution), but region-of-interest images could be obtained with a 200 μ s exposure. Qualitative analysis of the images at different rotation speeds showed no visible dependence of liquid distribution on acquisition (rotation) speed, suggesting the centrifugal effect on this sample and fluid was negligible.

In most “traditional” tomography experiments, over 1000 projections are collected for each 3-D volume. We collected high-resolution data (1800 projections, 0.2 and 0.5 s per 3-D data acquisition) for each dry sample prior to flow initiation to characterise the inter- and intra-grain porosity of each sample (Fig. 2a, b). This dataset can also be used as a structural prior for iterative reconstruction methods (e.g.; Kazantsev et al., 2015a, b; Van Eyndhoven et al., 2015). Reducing the angular density of the projections will shorten overall scan time. The reconstructed image quality does reduce as the projection density drops, but a significant increase in temporal resolution can be achieved. To achieve the maximum 20 Hz 3-D image acquisition rate, it is necessary to reduce the number of projections to 250 (200 μ s exposure) and 550 (90 μ s exposure) per scan.

Under-sampling (using fewer projections) to this degree does cause image quality to degrade, making differences between phases of similar attenuation difficult to observe and reducing the precision with which phase boundaries and volumes can be defined. When features of interest are relatively

small, or have complex geometries, then under-sampling can prevent quantitative analysis. However, here the contrast between the saline solution, rock, and air is high, and the primary target is the sub-second distribution of the fluid phase. Using the routine filtered back-projection reconstruction algorithms, testing with as few as 360 (Fig. 2c) and 180 projections (Fig. 2d) gave adequate images for basic quantification (Fig. 2e). Comparison of the two under-sampled datasets shows little increase in quality and little overall change between the 360 and 180 projection scans. For the smaller cells, it is possible to collect 360 projections at maximum image acquisition frequency.

The Miro camera has fixed on-board data storage, which allowed a maximum of approximately 20 000 projections to be collected for the field of view needed in this experiment (in this set-up, the number of projections is dependent on the field of view needed). The under-sampling allowed between 53 and 141 3-D datasets to be acquired in a single experimental run (2.5–7 s of acquisition at maximum frequency). Capturing a specific event with this short a duration of collection is challenging, especially as the high rotation speeds prevent observation of small volume changes in the live view radiographs. During collection the camera is armed and then triggered. Arming the camera starts projection acquisition, but those data are automatically overwritten until the trigger signal is received. The user can define the trigger to mark (a) the last image (all existing projections on the card are saved), (b) the first image (projections are collected until the card volume is full), or (c) an event (a user-defined number of projections before and after the trigger is saved). The latter option allows a small amount of buffer time for the user to capture a specific event, such as a drip entering the sample or a flow volume/composition change. Automatic triggering mechanisms such as light gates can be incorporated although in this experiment a webcam was used to observe the delivery of each drip.

To clarify the real-time nature of the data, the following terminology is used in all subsequent discussion: one or more scans are collected during each experiment, and each scan contains a number of 3-D “frames” which were acquired at a known speed and with a given frequency (defined by the time between the first projections of sequential frames). Data are offloaded from the camera between scans (download takes between 5 and 7 min). Several scans can be collected during an experiment. To highlight specific processes and features, most figures presented below show only 2-D slices (either perpendicular or parallel to flow) or a sub-volume of the reconstructed dataset. The post-processing and image segmentation are the same for all examples and were performed in Avizo™ and MATLAB. The methods are outlined in Appendix A.

Case study: observing sub-second dynamics using ultra-fast imaging

Wetting processes and the formation of rivulet networks are poorly understood in heterogeneous media. The processes operate on short timescales and are strongly influenced by sample morphology. Non-invasive 4-D tomography is ideal for capturing the dynamics of these behaviours. Figure 3 shows the evolution of the fluid distribution in a sub-volume of the sample immediately below the drip point. The data were collected under continuous acquisition mode, with a 3-D frame acquisition frequency of 0.07 s (system operating at 7 Hz rotation speed).

The first drip of KI solution (shown in blue) enters the dry sandstone bed (rendered in translucent grey) between frames F0 and F1 (Fig. 3). Initially ($t = 0.07$ s) the KI solution fills the inter-grain pore space to the front of the sub-volume (A) and then immediately begins lateral transport into the intra-grain porosity in the neighbouring grains (Fig. 3b). By the second frame ($t = 0.14$ s) the inter-grain volume has largely emptied, and the majority of the fluid is now within the grains. Intra-grain transport has also occurred with capillary-driven wetting of grains not in direct contact with inter-grain fluid (e.g. C), and drainage of previously saturated intra-grain pores can also be observed from other viewing angles. Between F1 and F2 most grains in contact with the inter-grain fluid have become partially or fully saturated. The remaining intra-grain fluid is generally limited to a few small pores and surface films (D). The figure only shows the relaxation after the first drip. As the experiment continues (F3 to F5), the fluid moves away from the initial location, with grain saturation increasing by radial and vertical redistribution, with oblique upward propagation of the saturation front in some instances (E). Most grains show an expansion and coalescence of grain-surface fluid films after initial intra-grain saturation (F), although grain surface wetting does occur with no, or prior to, intra-grain saturation (G). Capillary forces continue to redistribute the intra-grain pore fluid over the remainder of the experiment, with saturation of many grains decreasing as the system moves towards equilibrium. Grains at both the upper and lower extent of the wetted region (H) are almost completely saturated and then drained again within 1.54 s.

Vertical slices taken through the same sub-volume of the sample in the second scan of the experiment show the arrival of the third drip (Fig. 4). In this experiment the fluid feed was stopped during data download. When the pump was restarted, the first drip was used as the trigger for the scan acquisition. The initial frames of the second scan were before the arrival of the drip and act as a wet reference. Differences in the fluid volume across the download period result from the fall of a drip at the very end of the experiment, immediately after fluid feed was stopped. Figure 4 captures the major changes in the fluid distribution in the inter-grain porosity (blue). At this time, the intra-grain fluid distribution (grey)

has become largely stable, although a few grains still show some variability (Fig. 4b).

Passage of the drip causes transient fluid bridging (Fig. 4a). Most bridging connections are maintained throughout the drip cycle, but some transient connections are formed with every drip, and others form only occasionally. Temporary isolation of fluid (enhanced by capillary-driven draw in of the surface film during equilibration as observed at A; Fig. 4) could control chemical reaction rates in reactive systems. Upon reconnection the relatively small volume of this fluid would be rapidly diluted, but rates of dissolution would be locally reduced by limited fluid supply. A later part of the experiment is shown in Movie 1 (see Supplement).

Although the focus of this contribution is on method development, we include a limited selection of the quantitative analysis performed on these samples, as an illustration of the adequacy of the data quality. Quantitative analysis of inter- and intra-grain fluid volume shows the strong fluctuation over the course of two drips (Fig. 5) and the general instability of the fluid volume as it is redistributed. After the first drip, the saturation of the intra-grain porosity becomes more gradual, with significant increases on the arrival of subsequent drips. The inter-grain distribution equilibrates more slowly.

4 Ultra-high-speed imaging for longer-duration experiments

The experiments presented above allow capture of the high-speed dynamics but have very short overall durations (< 10 s). Capturing longer duration processes with the same 3-D frame acquisition rates requires lower frame acquisition frequencies. The Miro camera allows frames to be acquired either continuously (as above) or at user-defined intervals (hereafter called gapped acquisition). Gapped acquisition decreases frame acquisition frequency, while maintaining frame acquisition speeds and the resultant data therefore benefit from a lack of motion induced blurring in the reconstructed data.

Although any time or angular gap can be defined, in these experiments they were limited to odd multiples of 180° . This ensures that the initial projection of every frame was collected at the same angle. Registration of the data in this manner saves significant time aligning datasets during post-processing. A gapped acquisition with a spacing of 1 (G1) collects projections over the 0 – 180° sector of every rotation (one full rotation between acquisition start points); a gapped acquisition with spacing of 5 (G5) acquires projections from 0 to 180° on every third rotation (three full rotations between acquisition start points).

Although not used in this experiment, the camera memory can also be partitioned, allowing for a fixed and constant number of projections/frames in each partition. This could enable more efficient imaging across different tempo-

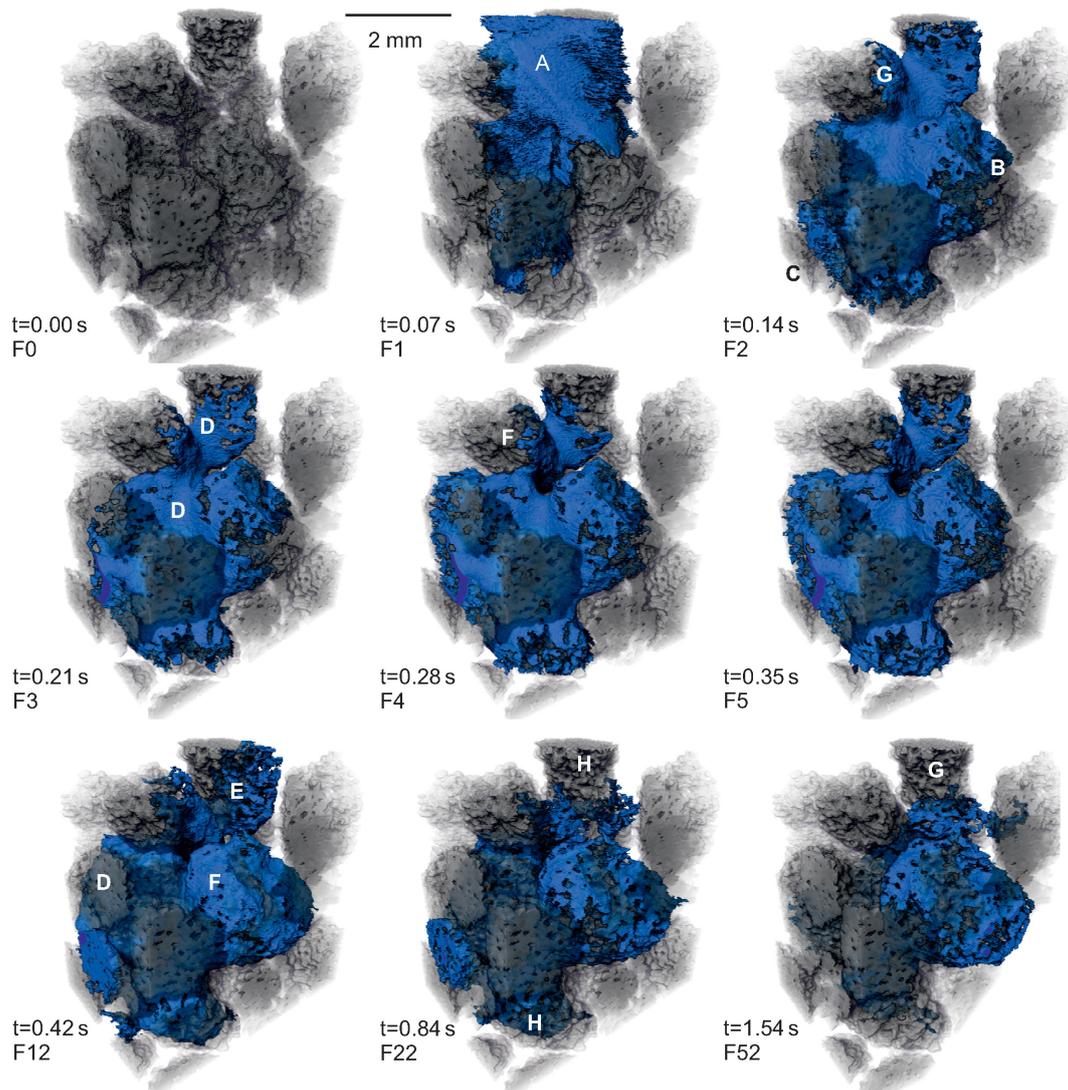


Figure 3. Time series showing the wetting of the sandstone gravel under the first drip. Three-dimensional frame acquisition frequency = 0.07/3-D volume (14 Hz). The first frame shows the dry sample with the sandstone grains rendered in semi-transparent grey. From the second frame the KI solution is shown in blue. Note the rapid changes in both the inter- and intra-grain fluid distribution (see text for discussion). The rendered volume is a small sub-volume of the larger dataset and shows a volume approximately 3.5 mm × 3.5 mm × 2 mm. Data are from sample LH_9, scan 1.

ral scales in one scan or the inclusion of much longer delays between high speed frame acquisitions.

4.1 Case study: capturing multiscale dynamics in pore-scale processes

To show the suitability of this method for capturing the variability of pore-scale flow dynamics, wetting experiments were scanned repeatedly at different gap lengths (continuous through to G25). Between 53 and 106 frames were collected in each scan. High-frequency data acquisition (continuous or G1) was used to capture the first few drips, with longer acquisition intervals (G3–G25) used to capture the slower

processes as the flow network evolved towards steady state through time (Fig. 6). In some experimental runs, variability in the high-speed processes through time was also assessed by repeating continuous G1 or G3 acquisitions at intervals among G5 or G7 acquisitions.

The larger fluid volume in these images for longer durations experiments makes visualisation of the 3-D fluid distribution challenging to show on the page (as in Fig. 3), and so a semi-transparent render of a distance map is used to show the variability in the local fluid volume (Fig. 6). Areas of inter-grain fluid appear as red regions, while the majority of the intra-grain fluid is shown in blue and green. In all scans at all frame acquisition frequencies, we see significant vari-

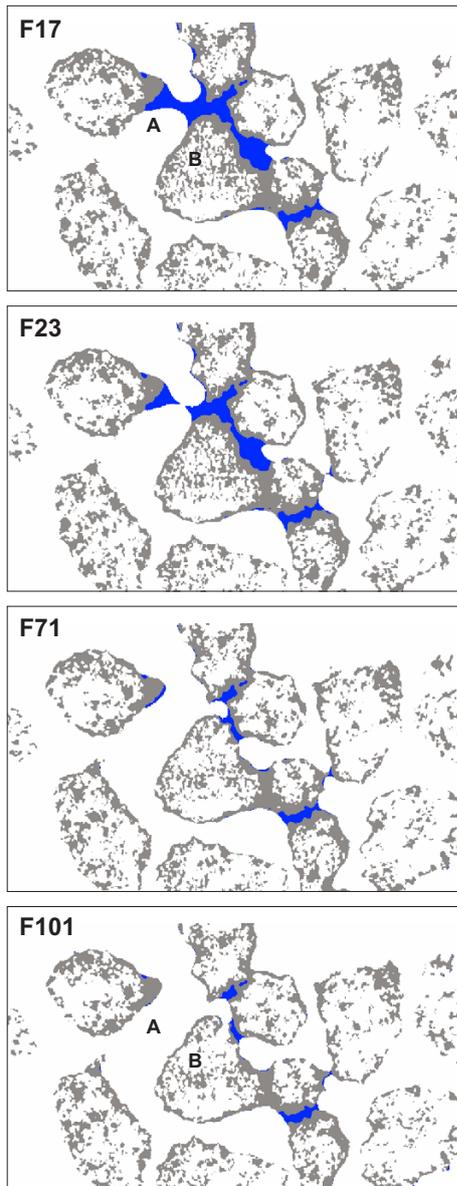


Figure 4. A single vertical 2-D slice through the second frame. The frames shown cover the passage of the through the sample. The inter-grain fluid is shown in blue and the intra-grain fluid in grey. Wet rocks are grey; KI solution outside the rocks is blue; air and dry rocks are in the background. Field of view is $3.5 \text{ mm} \times 2 \text{ mm}$. Data are from sample LH_9, scan 3.

ability in the distribution of the KI solution (Fig. 6) over the 35–60 min experiments.

From the same data it is possible to track the local changes in saturation at the grain scale. Well-established network analysis and morphological quantification algorithms can be applied to extract details of saturation on the pore scale (Fig. 7) or on the geometry of individual phase interfaces and contact angles (Andrew et al., 2014, 2015). For under-sampled data, higher levels of image noise in filtered back-

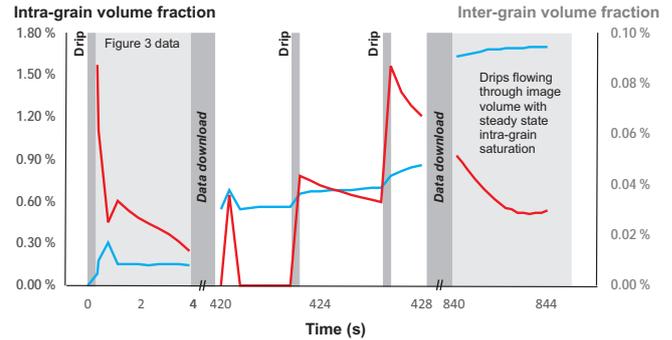


Figure 5. The change in fluid volume within the inter- (red) and intra-grain (blue) porosity over the experiment showing the arrival and subsequent equilibration of two drips. After the first drip the system equilibrates to very low inter-grain saturation. Data are from sample LH_9, scans 1 to 3.

projection reconstructions may limit quantitative surface curvature and contact angle analysis, but optimised iterative reconstruction methods may give more suitable data (see Sect. 5 below).

4.2 Case Study: in situ observation of mixing dynamics

One key application for ultra-high-speed in situ imaging is to track the interaction between multiple phases: capturing both spatial and temporal heterogeneity of replacement, dilution, and chemical reactions. After 30 min under a 6M saline feed the system was flushed with H_2O . Scanning was performed at G1 (0.07 s scan acquisition time, 7 Hz acquisition frequency); the scan started on arrival of the first H_2O drip. Other experimental runs tested replacing H_2O with KI solution and replacing oil with KI solution. The need for ultra-high-speed imaging to capture the complex and spatially variable mixing and dilution processes can be seen in Fig. 8. Dilution of the uppermost part of the local inter-grain network is instantaneous (Fig. 8a), but percolation of the dilution/replacement front through the image volume takes approximately 5 s (Fig. 8a–f). After this time there is little major change in the local greyscale values (Fig. 8g, h), implying a shift in mixing mechanisms. The 2-D slices through the same sub-volume perpendicular to the flow direction (Fig. 8) show the dilution is extremely heterogeneous and can occur on a range of timescales in neighbouring regions. The central pore volume (orange box, F30) is diluted over 10–15 scans, whereas a neighbouring region (connected in 3-D) is replaced in under five frames (yellow box F35). Some grains show variability in the KI concentration within the porosity, implying that, although the bulk porosity of this sandstone is approx. 20% and the connectivity is high, individual grains have discrete pore networks within them (blue boxes F45). Using slower data acquisition methods would not capture the dynamics of replacement.

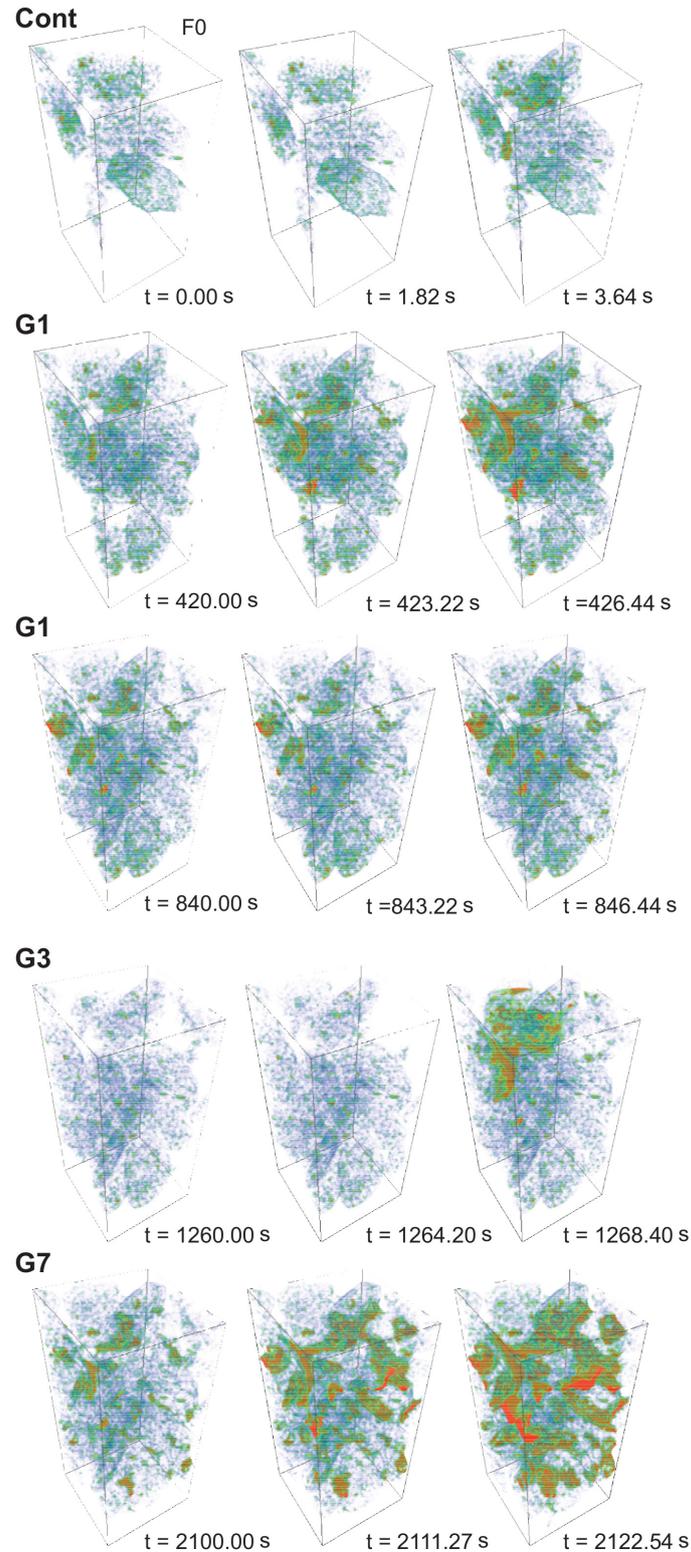


Figure 6. Three-dimensional renders of a gapped experimental run. For each scan series the first, middle, and last frame of each acquisition period is shown. Only the fluid volume has been rendered, with the colour representing the distance to the fluid interface. Cool colours represent small fluid volumes; warm colours represent fluid more than $20\ \mu\text{m}$ from the interface. The field of view is $1.5\ \text{mm} \times 1.5\ \text{mm} \times 3\ \text{mm}$. Data from sample LD_1, scans 1 to 5.

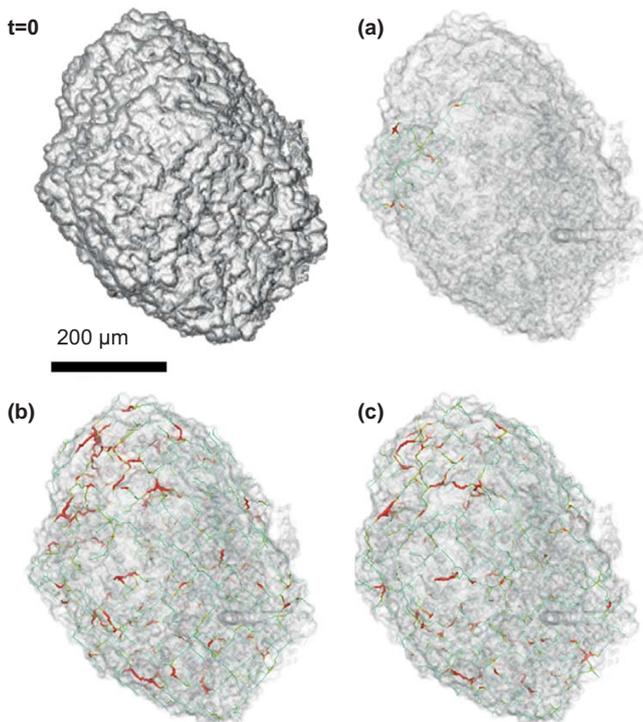


Figure 7. A 3-D render of a single grain showing the progression of saturation. The images show the evolution from first contact with the fluid to a steady saturated state over three scans. The colours and the size of tubes representing the pore network correspond to the local volume of the fluid within the pores: thick red tubes represent the higher local volume; thin blue tubes represent lower local saturation. Data are from sample LD_1, showing the dry scan ($t = 0$) and the final frames of (a) scan 1 (continuous), (b) scan 2 (G1), and (c) scan 3 (G1).

5 Moving beyond the state of the art

The ultra-high-speed imaging presented here allows capture of wetting and transport behaviour at the moderate-to-high spatial resolution necessary for quantitative understanding of the processes involved. The 3-D frame acquisition speed and frequency are substantially faster than has been previously achieved. Events such as Hayne jumps (millisecond timescale) and the subsequent relaxation (a few seconds) have been observed in 2-D imaging (acquisition of projections without rotation) (Armstrong et al., 2014b). The method presented here can provide data on the 3-D nature of those changes, at comparable timescales. Such quantitative data can be used to develop or validate larger-scale flow models and to better incorporate pore-scale processes.

The data presented here are reconstructed using standard filtered back-projection reconstruction algorithms. However, these methods are not optimal for under-sampled data. Recent advances in iterative reconstruction algorithms use information about the different phases contained within the data to enhance the quality of the reconstruction, with great

effect (Fig. 9). Visual comparison of this type of method (Fig. 9b, d) with reconstructions using filtered back projection (Fig. 9a, c) shows the improvement in phase resolution that can be achieved on fewer projection data. This example applies an iterative optimisation solver, and along total variation penalisation is incorporated as a regularisation step (Rudin et al., 1992; Little and Jones, 2010) that takes the sample porosity into account (i.e. uses the dataset themselves as prior information), with refinement at each iteration. The iterative data have fewer line artefacts and lower noise and permit accurate phase segmentation from fewer than 45 projections.

Iterative methods have significant potential to develop imaging dynamic pore-scale processes further still. Of specific interest are those methods that use high-quality dry scan data as a structural prior (Kazantsev et al., 2014, 2015a, b; Van Eyndhoven et al., 2015) and those that divide the data volume into static and dynamic regions using the static information from sequential scans to better define the reconstruction volume (Kazantsev et al., 2014, 2015a, b; Van Eyndhoven et al., 2015). Both allow reconstruction of better-quality images from under-sampled data (data with fewer projections) and also allow data of comparable or better quality to those displayed here to be produced from datasets with as few as 18 projections (Van Eyndhoven et al., 2015). These methods can improve the quality of the quantitative data, especially on the key phase interfaces; however, the key advance is likely to come from working with still higher degrees of under-sampling. However, iterative methods are computationally expensive and can take an order of magnitude longer to reconstruct. The methods applied here took between 15 (Fig. 9d) and 45 (Fig. 9b) min per volume, depending on the number of projections in being used. As such, their use in processing every frame of long-duration high-frame-rate data for near-real-time assessment remains unrealistic at present. This may change as computational efficiency increases, but we envisage that iterative methods are more practically applied retrospectively to achieve higher spatial resolution over short periods of critical interest.

At the 3-D volume acquisition rates achieved here, the projection exposure times and mechanical limit of the rotation stage means that imaging faster is not currently possible for the current samples. For higher-density materials, or higher resolution cameras where longer exposure times mean slower overall volume acquisition, iterative methods may improve temporal resolution. However, on the JEEP experimental setup, the main advantage of iterative methods is the ability to acquire fewer projections per scan and so collect more scans per experiment without sacrificing image quality. This could enable experimental run times to be extended by up to an order of magnitude.

The methods presented here provide a workflow for visualisation and quantification of sub-second dynamics in porous media (Fig. 10). We used the disaggregated sample to capture two spatial scales of fluid–pore interaction, but

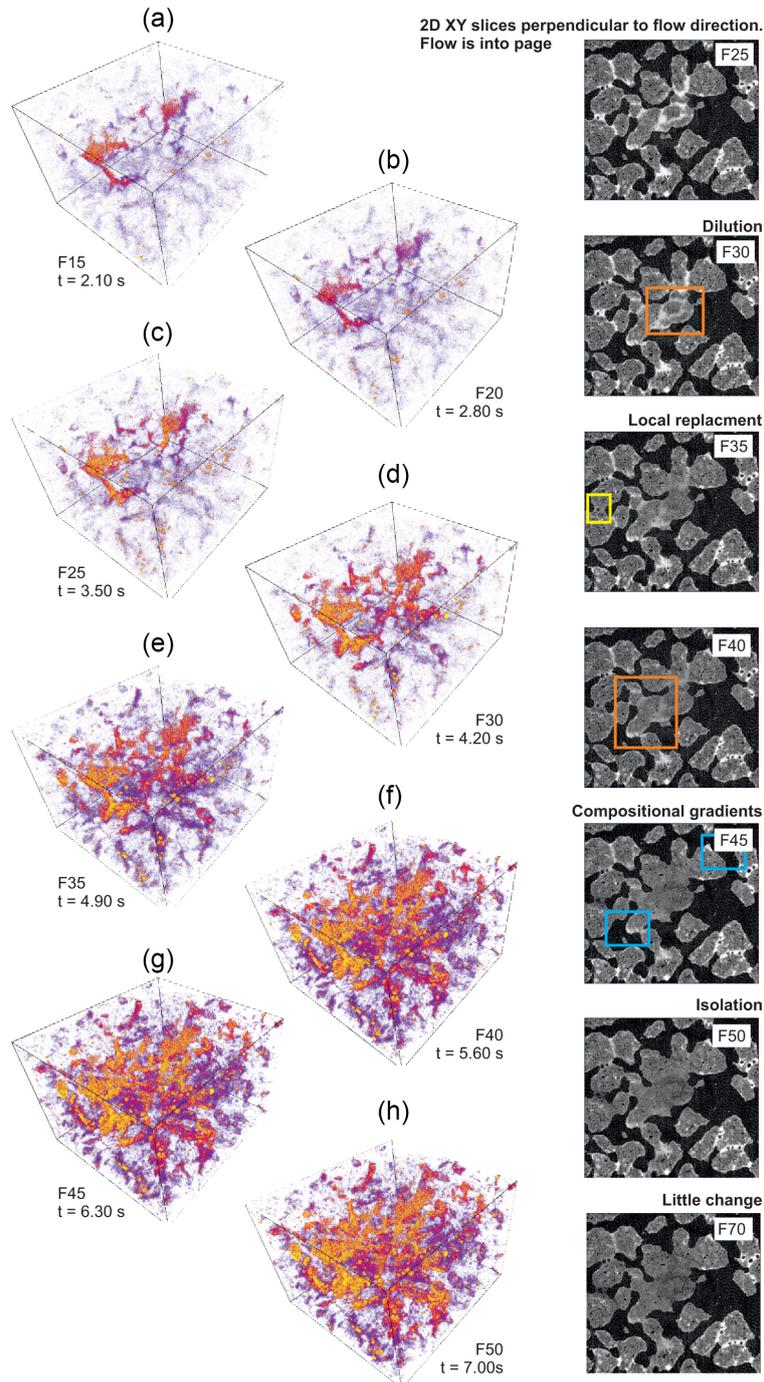


Figure 8. In situ observation of dilution. The 6M saline solution feed is replaced by H₂O at $t = 0$. (a–h) Three-dimensional renders of the differential density (greyscale) between the initial saline fluid volume and the saline + water mixture at the times shown. The frame acquisition rate was 0.07 s and acquisition frequency was 7 Hz (g). Data are from sample LD_3, scan 1. Yellow represents complete replacement with H₂O. Cooler colours represent volumes that are undergoing dilution. Volumes remaining at 6M KI are not shown, as they have undergone no change.

the overall experimental set-up can equally be used for solid core samples at the same spatial and temporal resolution. This experiment was performed at below maximum resolution, and imaging finer structures (grains or pores) at resolu-

tions of 1.2 μm is possible. For higher-resolution or higher-density samples slightly slower projection acquisition times may be needed, and in extreme circumstances this will reduce overall frame acquisition frequency. The ability of the

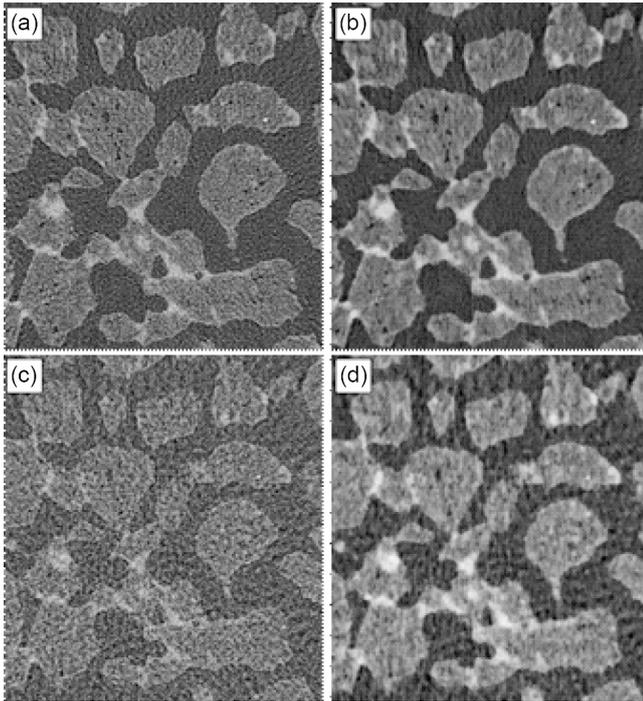


Figure 9. Two-dimensional reconstructed slices through a gravel sample using filtered back projection (**a**, **c**) and an iterative optimisation method with the total variation regularisation incorporated as prior information (**b**, **d**). Reconstructions are obtained using 180 projections (**a**, **b**) and 90 projections (**c**, **d**). The cropped area shown is approximately $3.5 \text{ mm} \times 3.5 \text{ mm}$. Note the increased definition of both the inter- and intra-grain porosity that can be achieved for a given number of projections. Data are from sample LD_3, first frame, scan 1. The optimised iterative methods took approximately 45 min (**b**, 180 projections) and 15 min (**d**, 90 projections) per 3-D frame, making them a viable element of the post-processing workflow for key datasets.

experimental set-up to acquire large numbers of projections at constant angular density means that even higher temporal resolution could be achieved. Reconstruction of volumes from datasets with small rotational offsets (typically about $10\text{--}20^\circ$, i.e. scan 1 reconstruction from 0 to 180° , scan 2 reconstruction from 10 to 190° , etc.) could increase temporal resolution by over an order of magnitude (Van Eyndhoven et al., 2015), especially when coupled to iterative methods that can reduce image noise and enhance phase boundary definition in the reconstructed data.

Complex multiphase, spatially heterogeneous micro- and pore-scale processes control many key macro-scale geological system responses. However, the method is not limited to the simple drip-fed experiments on geological materials presented here and can as be easily applied to high-speed processes operating in any material (metals, ceramics, plastics, biological materials, foodstuffs, etc.) suitable for synchrotron X-ray tomography at JEEP/i12. It can also be applied in other dynamic systems for samples and materials

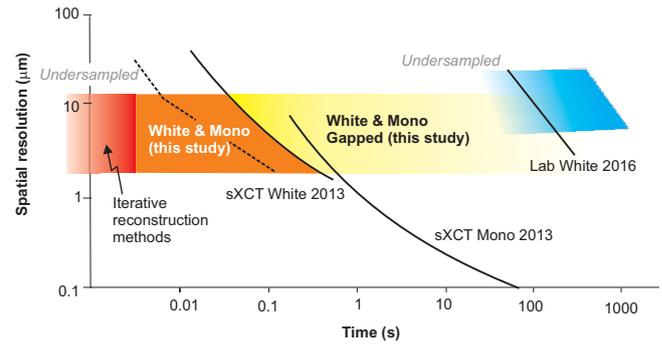


Figure 10. A schematic showing the current state-of-the-art imaging capabilities for synchrotron and laboratory XCT scanning (modified after Bultreys et al., 2016; Maire and Withers, 2014).

undergoing modification through treatment at extreme temperature (hot or cold) or experiencing deformation (tension, compression, cyclic fatigue, fracture, or shear). The development of the real-time 4-D imaging methods we have presented means that quantitative understanding of all of these processes is now possible, in porous media and beyond.

6 Data availability

The projection data and the reconstructed volumes of these experiments are available from the authors.

Appendix A: Post-processing and image analysis

Visualisation and quantification was performed using Avizo™ and MATLAB. The greyscale volumes were cropped to the sample boundaries to reduce all subsequent processing times. When processing continuous acquisition data the frames were first aligned to allow direct comparison and quantitative analysis. For the gapped scans, all data are acquired over the same angle and no additional alignment is required. The italicised steps in the following refer to built-in algorithms in Avizo™, and the reader is directed to the associated documentation for further details of the implementation.

The edge-preserving Bilateral Filter was applied to the greyscale-filtered back-projection reconstructions to reduce image noise. Water, air, and rock phases were segmented using the Interactive Thresholding tool, and the resultant binary data volumes were used in subsequent analysis. The inter- and intra-grain porosity was defined by creating two separate binary masks: one by segmentation of the wet and dry grains and one by segmentation of the fluid (inter- and intra-granular). Logical combination of these allows calculation of the inter- and intra-granular fluid volumes. Fluid volumes were calculated using in-built image measures (Volume3-D).

Additional quantitative analysis of the intra-grain porosity was performed on single grains isolated using the MATLAB command `imfill` followed by Watershed segmentation. The pore network within the grain was displayed using the Centerline Tree module, which applies the TEASAR algorithm (Sato et al., 2000) defining local Euclidian distance to the nearest object boundary. Fluid thickness was calculated by applying a Chamfer 3-D distance map to the fluid label volume to calculate the distance to the nearest surface. The exchange of KI for H₂O was tracked by calculating a differential image between the initial scan and each subsequent frame.

The Supplement related to this article is available online at doi:10.5194/se-7-1059-2016-supplement.

Author contributions. The study was designed by Katherine J. Dobson and Sam A. McDonald; reconstruction and visualisation was led by Sophia B. Coban and Katherine J. Dobson with assistance from Samuel A. McDonald. Robert C. Atwood led camera integration at Diamond Light Source. Katherine J. Dobson, Sophia B. Coban, Sam A. McDonald, Joanna N. Walsh, and Robert C. Atwood performed the experiments at i12/JEEP. The lead author was Katherine J. Dobson, with all other authors contributing to the manuscript.

Acknowledgements. Katherine J. Dobson was supported in this work by EVOKES ERC247076 and NERC NE/M018678/1; Sophia B. Coban is supported by EPSRC EP/J010456/1, the School of Mathematics, and the BP International Centre for Advanced Materials. Philip J. Withers wishes to acknowledge the BP International Centre for Advanced Materials for funding under ICAM 03. Samuel A. McDonald is grateful to Zeiss for funding his fellowship. We thank Diamond Light Source for access to beamline I12 (beam time award EE10500-1).

Edited by: M. Halisch

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