Viscosities of melts in the Na₂O-FeO-Fe₂O₃-SiO₂ system and factors controlling relative viscosities of fully polymerized silicate melts

DONALD B. DINGWELL* and DAVID VIRGO

Geophysical Laboratory, Carnegie Institution of Washington, 2801 Upton Street NW. Washington, D.C., 20008, U.S.A.

(Received May 12, 1987; accepted in revised form November 13, 1987)

Abstract—The viscosity-temperature relationships of nine melts in the Na₂O-Fe₂O₃-SiO₂ system (in equilibrium with air) have been measured in the temperature range of 1450-800°C, using the concentric cylinder method. ⁵⁷Fe Mössbauer spectra were obtained on quenched samples and show that all melts with ≥ 20 mole% ferrite component contain $\geq 95\%$ Fe as tetrahedrally-coordinated ferric ions.

The compositions investigated lie along the SiO_2 -NaFeO₂ and Na₂Si₄O₉-Na₆Fe₄O₉ joins. The viscosities of melts decrease strongly with decreasing silicate content along these joins. In contrast, the temperature dependence of viscosity does not vary significantly along these joins.

Comparison with equivalent melt compositions (related by the trivalent cation exchange operators $AlFe_{-1}$, BFe_{-1} and $GaFe_{-1}$) in the Na₂O-Al₂O₃-SiO₂, Na₂O-B₂O₃-SiO₂ and Na₂O-Ga₂O₃-SiO₂ systems, shows that viscosities decrease in the order aluminosilicate > ferrosilicate, (galliosilicate ?) > borosilicate.

The electronegativities of the trivalent cations are inversely correlated with the relative viscosities of melts in these systems. Similarly, the electronegativities of network-stabilizing cations are inversely correlated with melt viscosity for alkali and alkaline-earth aluminosilicate melt systems. The variation in the viscosity of tectosilicate melts is correlated with estimated average T-O-T bond angles, and exothermic heats of solution of quench glasses. Structural controls of viscosity discussed are tetrahedral ordering and relative bond strengths.

The acmite component in natural, peralkaline, silicic volcanics will not contribute directly to high melt viscosities for these lavas.

INTRODUCTION

VISCOSITY IS PERHAPS the single most important physical property of magmas (and lavas) (WILLIAMS and MCBIRNEY, 1979; FISHER and SCHMINKE, 1984). For relatively dry, crystal-poor, peralkaline lavas such as those of the Kenya rift zone (BAILEY and MACDONALD, 1987) the determination of lava viscosity during eruptions is largely a question of estimating '1) the temperature of eruption, and 2) the temperature-viscosity relationship of the erupted silicate melts whose compositions are probably very similar to the compositions of the glassy peralkaline rocks.

One of the most diagnostic chemical expressions of peralkalinity in the Kenyan comendites, trachytes and pantellerites is the presence of CIPW-normative acmite. Estimation of the eruptive viscosities of these lavas is, therefore, hindered by the lack of data for the viscosity-temperature relationships in the system Na₂O-FeO-Fe₂O₃-SiO₂, particularly for the acmite component. In fact, as noted by BOTTINGA and WEILL (1972), there is a paucity of viscosity data for highly oxidized iron-bearing silicate melts.

A comparison of the viscosity-temperature relationships of ferric iron- and aluminum-bearing melts may provide further insight into the relationship between melt viscosity and melt structure because (like aluminum in the Na₂O-Al₂O₃-SiO₂ system) ferric iron appears to be in tetrahedral coordination in melts in the Na₂O-FeO-Fe₂O₃-SiO₂ system that are equilibrated with air (FOX *et al.*, 1982; HENDERSON *et al.*, 1984; MYSEN *et al.*, 1980; VIRGO *et al.*, 1983) but (unlike aluminum in the Na₂-Al₂O₃-SiO₂ system) may not exhibit a random distribution of ferrite and silicate tetrahedra (MYSEN *et al.*, 1985a; VIRGO and MYSEN, 1985). Thus a comparison of melt viscosities in the Na₂O-Fe₂O₃-SiO₂ and Na₂O-Al₂O₃-SiO₂ systems might contribute to a better understanding of the relationship between viscosity and ordering of tetrahedrally-coordinated cations in these melts.

The system Na₂O-FeO-Fe₂O₃-SiO₂ was selected for an investigation of the viscosity-temperature relationships of ironbearing silicate melts for the following additional reasons. Firstly, phase equilibria had been determined in air for the compositional region of this study (BOWEN et al., 1930) ensuring that the superliquidus viscosities of the compositions were experimentally accessible. Secondly, a direct comparison of the viscosities of aluminum- and ferric iron-bearing silicate melts was possible by using existing data for the viscositytemperature relationships of the equivalent aluminum-bearing melts in the Na₂O-Al₂O₃-SiO₂ system (RIEBLING, 1966; DINGWELL, unpublished). In addition, comparisons can be made with limited melt viscosity data for the Na₂O-B₂O₃-SiO₂ and Na₂O-Ga₂O₃-SiO₂ systems. Thirdly, previous studies of the redox state of iron in melts in the Na₂O-SiO₂-Fe-O system (VIRGO et al., 1983) indicated that these melts, when equilibrated with air, would be extremely oxidized and thus the effect of ferrous iron could be minimized.

EXPERIMENTAL METHOD

Starting compositions

^{*} Present address: Bayerishes Geoinstitut, Universität Bayreuth, Postfach 10 12 51, 8580 Bayreuth, FRG.

The compositions of melts investigated in this study are plotted in the ternary system Na₂O-Fe₂O₃-SiO₂ along the SiO₂-NaFeO₂ and Na₂Si₄O₉-Na₆Fe₄O₉ joins (Fig. 1). The viscosity-temperature relationships were determined for melts with 12.5, 25 and 50 mole percent NaFeO₂ (denoted SFN6, Fe-albite and acmite, respectively) on the



FIG. 1. The compositions of melts investigated in this study plotted in the Na₂O-Fe₂O₃-SiO₂ system. Diagram in wt.%. Asterisks refer to compositions whose oxidation state-viscosity relationships have been investigated (see text for discussion of oxidation states).

SiO₂-NaFeO₂ join, and for melts containing 5, 10, 20, 30, 40 and 50 mole percent Na₆Fe₄O₉ (denoted NS4F5, NS4F10, NS4F20, NS4F30, NS4F40 and NS4F50) on the Na₂Si₄O₉-Na₆Fe₄O₉ join.

The starting materials for viscosity determinations were glasses prepared from mixtures of reagent-grade Na₂CO₃ and Fe₂O₃, and purified quartz sand. Batches of carbonate and oxides equivalent to be decarbonated weight of 70 g were ground under alcohol in an agate mortar for 2 hours and then melted into a 25 cc thin-walled platinum crucible at 1400°C for 2 hours. These batches were poured and broken from the flexible platinum crucible and the quenched glasses were melted into a thick-walled viscometry crucible at 1450°C and stirred with a viscometry spindle for 30 min. The viscometry crucible and spindle are described below.

Viscometry

Viscosities were measured in air, in the temperature range of 1450-800°C, using the concentric cylinder method. The design and operation of the viscometer used in this study have been described by DINGWELL (1986). The samples were contained in a cylindrical Pt₈₀Rh₂₀ crucible (5.1 cm in height, 2.56 cm inner diameter, 0.1 cm wall thickness) in a SiC-heated vertical tube furnace. The temperatures were monitored during viscometry runs with a Pt-Pt₉₀Rh₁₀ thermocouple (shielded in Pt tubing) calibrated against the melting point of Au and were checked periodically with a second, shielded Pt-Pt₉₀Rh₁₀ thermocouple immersed in the melt sample. The accuracy of temperature determinations is taken as $\pm 2^{\circ}$ C.

The viscosities were determined with a Brookfield RVTD viscometer head. This viscometer head drives a spindle at a range of constant angular velocities (0.5 to 100 rpm) and digitally records the torque exerted on the spindle by the sample. Two Pt₈₀Rh₂₀ spindles were used in this study. Both spindles have the cross-section of a cylinder with 45° conical ends and 0.24 cm diameter stems to reduce endeffects. The low viscosity spindle has a 1.44 cm diameter and a 3.32 cm length and the high viscosity spindle has a 0.32 cm diameter and a 4.2 cm length. The viscometer head and spindles were calibrated for viscosity determinations using NBS SRM 711 standard lead-silica glass for which the temperature-viscosity relationship is accurately known. The precision of determinations involving separate occupations of the temperature set-point and separate immersions of the spindles is $\pm 3\%$ (2 σ). The accuracy of determinations is taken as the sum of uncertainties from standard and sample determinations at $\pm 6\%$ (2 σ). For each sample, viscosity determinations were initiated at the highest temperature and then made at successively lower temperatures. Thermal equilibrium was monitored with a chart recording of the sample viscosity during the decreasing temperature steps. One hour was sufficient for equilibration of the sample over each 50-

100°C temperature decrease. At the termination of each viscometry run a high temperature viscosity measurement was redetermined to check for instrumental drift during the run. No drift was observed for any of these samples.

Torque measurements were made over a range of angular velocities for each sample. In all cases, the viscosities obtained were independent of angular velocity. At each temperature step, a Pt wire loop was dipped into the melt and a small ($\simeq 100$ mg) sample of melt was drawn and guenched in water. Viscosity measurements were continued with decreasing temperature until crystallization occurred. Crystallization during the final cooling step for each sample resulted in erratic viscosity readings and was easily confirmed by inspection of loop samples obtained at those times.

After viscometry runs, the samples were poured from the crucible onto a steel plate. Chips of each sample were analysed by electron microprobe for composition and homogeneity and the results of these analyses are compared with the "weighed in" compositions in Table 1. Excellent agreement was found between the initial and final (analvsed) compositions.

57Fe Mössbauer spectra

The oxidation state and structural state of iron in guenched loop samples collected during the viscometry runs were determined using ⁵⁷Fe Mössbauer spectroscopy. The resonant absorption spectra were recorded at 298 K with a 25 mC 57Co/Pd source on powdered samples mixed with plastic transoptic powder and pressed into thin discs. The absorber thickness was constrained to be $\approx 5 \text{ mg Fe/cm}^2$. Mirror image spectra were recorded over 512 channels. The data were analysed separately and the results are the average values obtained from the left and right sides. The spectral data were deconvoluted with a least-squares routine using Lorentzian lines and with area and width constraints as discussed by VIRGO and MYSEN (1985). It was not possible to fit the low-velocity component of the ferrous doublet of highly oxidized glasses (Table 2), despite the use of area constraints.

The calculated $Fe^{3+}/\Sigma Fe$ value and the hyperfine parameters, quadrupole splittings and isomer shifts for both ferrous and ferric iron are given in Table 2.

Quench glasses with compositions along the tectosilicate join are highly oxidized and Fe³⁺/2Fe is invariant as a function of Si/Fe³⁺ within experimental error. For the tetrasilicate quench glasses, $Fe^{3+}/$ ΣFe systematically decreases with increasing Fe³⁺/Si. A similar effect is observed with increasing Al/Si in iron-bearing, alkali and alkalineearth aluminosilicate glasses (MYSEN, 1988).

The values of ferric isomer shift for the alkali-silicate glasses are in the range of 0.23-0.29 mm/sec and are similar to the values found for oxidized Na2O-SiO2 glasses containing 5 wt.% Fe2O3 (VIRGO et

TABLE 1. Analyzed melt compositions.

Electron microprobe analysis ^{1,2} Stoichiometric							
Sample	Na ₂ 0	Fe2033	Si0 ₂	total	Na ₂ 0	Fe ₂ 03	\$10 ₂
Na ₂ Si40g-Na	6 ^{Fe} 4 ⁰ 9			/			
NS4-F5 NS4-F10 NS4-F20 NS4-F30 NS4-F40 NS4-F50	21.21 24.03 25.93 27.59 29.05 31.61	5.31 10.29 18.93 27.66 34.27 40.18	73.48 65.68 55.14 44.76 36.68 28.22	(98.08) (98.26) (99.67) (100.23) (101.51) (101.75)	21.82 23.05 25.30 27.30 29.10 30.70	5,11 9,90 18,63 26,38 33,31 39,55	73.07 67.05 56.07 46.32 37.60 29.76
SiO ₂ -NaFeO ₂							
SFN6 Fe-albite Acmite	5.96 10.89 14.10	15.54 27.78 35.78	78.50 61.34 50.12	(97,45) (100,70) (100,51)	5.83 10.65 13.41	15.02 27.43 34.56	79.14 61.93 52.02

 1 Wavelength dispersive analyses using a JEOL JSM-35 instrument and Krisel control system at: 15 kV accelerating voltage, 60 nA beam current on carbon, 30 sec maximum count times and continuously moving the sample under a 10 x 10 micron rastered beam. Synthetic glass standards were used. Maximum relative uncertainties at 3 s.d. are: 2.2.4% (Na), 2.0% (Fe) and 1.7% (Si). Oxygen by stoichiometry. 3 Analyses normalized to 100% for comparison with stoichiometry. All iron recalculated as Fe_0.3.

 ^{57}Fe Mössbauer parameters at 298K for glasses,along the $Si0_2\text{-NaFe}0_2$ and $Na_2Si_40_9\text{-Na}_6Fe_40_9$ joins . Fe²⁺ Fe³⁺ ferric/total iron Sample 1.S.² Q.S.³ Q.S. I,S. NS4F5 NS4F10 0.82 0.89 0.95 0.28 0.29 0.25 0.24 0.89 1.96 2.21 2.04 0.83 NS4F20 NS4F30 0.86 1.02 (4) (5) (4) (4) 0.86 1.0 0.98 0.22 0.86 (5) (4) NS4F40 NS4F50 1.0 1.83 SFN6 0.92 0.23 0.90 0.83

0.24

0.86

(5) 0,87

(5) 1.83

all melts equilibrated with air.

0.94

2 dif melts equilibre with the 3 gisomer shift (mm/sec). 4 guadrupole splitting (mm/sec). 5 line fit, see text.

53 line fit, see text.

TABLE 2.

Fe-albite

Acmite

al., 1983). The isomer shift values for Fe^{2+} are >0.83 mm/sec and are also similar to those found for oxidized alkali and alkaline-earth aluminosilicates (VIRGO et al., 1983; VIRGO and MYSEN, 1985). The isomer shift values for ferric iron are consistent with tetrahedrallycoordinated iron (see VIRGO and MYSEN, 1985, for a detailed discussion and review).

RESULTS

The Na₂O-FeO-Fe₂O₃-SiO₂ system

The viscosity data are presented in Table 3 and plotted as function of reciprocal absolute temperature in Fig. 2 (Na₂Si₄O₉-Na₆Fe₄O₉) and Fig. 3 (SiO₂-NaFeO₂). The viscosities of melts on the Na2Si4O9-Na6Fe4O9 join exhibit a trend

		TABLE 3.	Viscometry d	ata.	
Na ₂ Si ₄ 0 ₉ -N	^{a6^{Fe}4⁰9}				
NS4F5	NS4F10	NS4F20	NS4F30	NS4F40	NS4F50
T log n	⊺ log n	T log n	T log n	T log n	T log n
1401 1.94 1346 2.11 1287 2.31 1231 2.50 1179 2.70 1127 2.94 1073 3.20 1018 3.52 961 3.83 Si0 ₂ -NaFeO.	1417 1.73 1358 1.80 1301 1.98 1242 2.19 1188 2.37 1135 2.58 1081 2.83 1027 3.19 975 3.37 922 3.71 860 4.14 809 4.59	1402 1.20 1351 1.33 1292 1.47 1233 1.68 1178 1.87 1128 2.06 1074 2.30 1024 2.52 970 2.80 914 3.12 879 3.34	1392 0.74 1353 0.84 1294 1.02 1236 1.22 1182 1.40 1126 1.62 1089 1.76 1019 2.12 962 2.36 906 2.67 844 3.04	1408 0.33 1352 0.47 1293 0.67 1293 0.81 1186 0.98 1124 1.21 1068 1.41 1011 1.71 957 1.97 905 2.27 851 2.61	1384 0.07 1332 0.20 1278 0.37 1226 0.52 1172 0.71 1119 0.90 1067 1.11 1012 1.30 956 1.65
SFN6	Fe-albite	Acmite			
1426 3.51 1372 3.75 1319 4.01 1267 4.33	1447 2.12 1403 2.28 1353 2.46 1301 2.69 1250 2.92 1202 3.16 1149 3.44 1101 3.77 1053 4.14	1408 1.41 1362 1.54 1312 1.73 1260 1.94 			

Temperatures are \pm 2°C, viscosities are \pm 6% or 0.05 poise, whichever is larger.



FIG. 2. Viscosity temperature relationships on the Na₂Si₄O₉-Na₆Fe₄O₉ join.

of decreasing viscosity with decreasing Na₂Si₄O₉ content (Fig. 2). Similarly, the viscosities of melts on the SiO₂-NaFeO₂ join decrease with decreasing silicate content.

The temperature-dependence of the viscosities of some melts investigated in the Na₂O-FeO-Fe₂O₃-SiO₂ system may



FIG. 3. Viscosity temperature relationships on the SiO₂-NaFeO₂ join.

be described within error as linear functions of reciprocal temperature and thus may be fitted to an equation of the form:

$$\log_{10} n = \log_{10} n_0 + Ea/2.303RT \tag{1}$$

where n_0 and Ea are constants, termed the pre-exponential factor and the activation energy, respectively, R is the gas constant and T is the absolute temperature. The "Arrhenius" equations derived from these linear reciprocal temperature fits are presented in Table 4.

More generally, however, the observed viscosity-temperature relationships of Figs. 2 and 3 illustrate a non-linear (non-Arrhenian) dependence of viscosity on reciprocal absolute temperature. Several investigators have observed this behavior for silicate melts (e.g., RICHET, 1984). In this study, data were obtainable over the largest temperature ranges for Fe-albite and NS4F10 and these melts best illustrate the nonlinearity for the SiO₂-NaFeO₂ and Na₂Si₄O₉-Na₆Fe₄O₉ joins, respectively.

The non-linearity of the viscosity-temperature relationship may be incorporated into Eqn. (1) by the addition of a third adjustable parameter, denoted T_0 , and termed the fictive temperature:

$$\log_{10} n = \log_{10} n_0 + Et/2.303R(T - T_0).$$
(2)

The results of fits to this equation, termed the TVF (Tammann-Vogel-Fulcher) equation, are also presented in Table 4. The value of the slope parameters (Ea and Et) obtained from these two equations indicate that the viscosity-temperature relationships of melts on both the tetrasilicate and tectosilicate joins are not strong functions of composition.

Existing data for the Na₂O-Al₂O₃-SiO₂, Na₂O-B₂O₃-SiO₂ and Na₂O-Ga₂O₃-SiO₂ systems

The viscosity-temperature relationships of melts on the Na₂Si₄O₉-Na₆Fe₄O₉ and SiO₂-NaFeO₂ joins (Figs. 2, 3) are compared with the viscosity-temperature relationships for aluminum-bearing melts of equivalent stoichiometry (i,e., compositionally related by the Fe³⁺Al-1 exchange operator) in the system Na₂O-Al₂O₃-SiO₂ in Figs. 4 and 5 (inset) (Na₂Si₄O₉-Na₆Al₄O₉ data from DINGWELL, unpublished; SiO₂-NaAlO₂ data from RIEBLING, 1966). The viscosities of melts on the Na₂Si₄O₉-Na₆Al₄O₉ and SiO₂-NaAlO₂ joins are considerably higher than the viscosities of melts on the

TABLE 4. Arrhenius and TVF coefficients.

Sample	Arrhenius		Tammann-Vogel-Fulcher		
	Ea ¹	⁻¹⁰⁹ 10 ⁿ 0	Et ¹	-lag ₁₀ n0	т _о
NS4F5	40.8	3.41	23.9	1.95	333
NS4F10	39.8	3.53	20.4	1./1	3/3
NS4F202	35.3	3,38	22.1	2.32	
NS4F40	35.2	4.27	21.2	3.02	301
NS4F50	33.7	4.39	21.4	3.34	287
20002					
SEN6	61.9	4.41	~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~	<u> </u>	-
Fe-albite	52.7	4.61	27.0	2.46	427

Ea and Et in units of kcal/mole Arrhenius behavior within error



FIG. 4. Comparison of viscosities of melts on the Na₂Si₄O₉-Na₆Fe₄O₉ and Na₂Si₄O₉-Na₆Al₄O₉ joins at 1150°C. Boron-bearing melt data from TAIT et al. (1984), Aluminum-bearing melt data from DINGWELL (unpublished).

Na₂Si₄O₉-Na₆Fe₄O₉ and SiO₂-NaFeO₂ joins, respectively. The viscosity-composition relationships are, however, qualitatively similar for the SiO₂-NaFeO₂ and SiO₂-NaAlO₂ joins



FIG. 5. Comparison of the viscosity temperature relationships of melts of NaAlSi₃O₈, NaFeSi₃O₈, NaBSi₃O₈ and NaGaSi₃O₈. (inset: comparison of the viscosities of melts on the SiO₂-NaAlO₂, SiO₂-NaFeO2 and SiO2-NaBO2 joins at 1400°C.) Data from DAY and RINDONE (1962) and RIEBLING (1966), this work, TAIT et al. (1984) and LAPP and SHELBY (1984).

but quite different for the $Na_2Si_4O_9$ - $Na_6Fe_4O_9$ and $Na_2Si_4O_9$ - $Na_6Al_4O_9$ joins.

TAIT et al. (1984) have determined the viscosity-temperature relationships of sixteen melts in the Na₂O-B₂O₃-SiO₂ system in air, over the temperature range of 950-1500°C. The composition range of their investigation includes portions of the Na₂Si₄O₉-Na₆B₄O₄ (mole fraction Na₆B₄O₉ = 0.13 to (0.38) and SiO₂-NaBO₂ (mole fraction NaBO₂ = (0.20 to 0.55)) compositional joins. TAIT et al. (1984) fitted their isothermal viscosity-composition data to a series of cubic polynomials at four temperatures (1050, 1100, 1150 and 1200°C). The resulting residuals are <0.03 log₁₀ units of viscosity. The viscosity-temperature relationships of the borosilicate melts represented in Figs. 4 and 5 (inset) are derived from Eqn. 5 of TAIT et al. (1984). The viscosity-temperature relationships of sodium borosilicate melts in Figs. 4 and 5 (inset) are restricted to the ranges of composition and temperature of the TAIT et al. (1984) study, and are thus interpolations of the data.

The principal observation from Figs. 4 and 5 (inset) is that viscosities for Na₂O-Fe₂O₃-SiO₂ and Na₂O-B₂O₃-SiO₂ melts are similar. The viscosity of melts along the Na₂Si₄O₉-Na₆B₄O₉ and SiO₂-NaBO₂ joins decrease smoothly with decreasing silicate content. The viscosities of the boron-bearing melts are systematically lower than the equivalent iron-bearing melts over the entire range of comparison. The similarity of the bulk viscosities and the viscosity-temperature relationships of melts in the Na₂O-Fe₂O₃-SiO₂ and Na₂O-B₂O₃-SiO₂ systems contrasts with the much higher bulk viscosities and larger temperature dependencies of melts in the Na₂O-Al₂O₃-SiO₂ system.

SHELBY and coworkers (PIGUET and SHELBY, 1985; PI-GUET et al., 1985; LAPP and SHELBY, 1986) have investigated the (glass) transformation-range viscosity of several compositions in alkali galliosilicate systems. LAPP and SHELBY (1986) report viscosity data for several melts/glasses in the $Na_2O-Ga_2O_3-SiO_2$ system. The data were obtained using a beam-bending viscometer in the temperature range of 400 to 800°C. The data, although not tabulated, are presented in the form of isokom (iso-viscosity) curves of temperature versus composition at constant SiO₂ or Na₂O content. The 10¹² poise (10¹¹ Pa-s) isokom passes through NaGaSi₃O₈ (Ga-albite) composition at 750°C (as read from Fig. 9; LAPP and SHELBY, 1986); this data has been plotted in Fig. 5. The viscosity-temperature relationship is not reported for NaGaSi₃O₈ but is reported (from 670-730°C) for another composition on the SiO₂-NaGaO₂ join (0.6SiO₂-0.4NaGaO₂, Fig. 8, LAPP and SHELBY, 1986). This viscosity-temperature relationship is shown in Fig. 5 and has been used as a first approximation to the slope of the viscosity-temperature relationship of NaGaSi₃O₈ melt (dashed line, Fig. 5). The justification for this approximation is that the SiO₂-NaFeO₂, SiO₂-NaAlO₂ and SiO₂-NaBO₂ joins show a weak composition-dependence of the slope of the viscosity-temperature relationship in the range of mole fraction $SiO_2 = 0.5$ to 0.75.

Figure 5 provides a comparison of the viscosity-temperature relationships of NaAlSi₃O₈, NaFeSi₃O₈, NaBSi₃O₈ and NaGaSi₃O₈. The viscosity-temperature of NaAlSi₃O₈ has been extended to lower temperatures in Fig. 5 using the low temperature data of TAYLOR and RINDONE (1970). The important point to be noted from Fig. 5 is that the viscosity of NaGaSi₃O₈ is significantly lower than that of NaAlSi₃O₈ (approximately 3.75 log₁₀ units). Additionally, the measured viscosity-temperature relationship of NaAlSi₃O₈ is subparallel to that which was estimated above for NaGaSi₃O₈. We do not advise a quantitative extrapolation of the estimated NaGaSi₃O₈ viscosity-temperature relationship to higher temperature, but we consider it likely, based on the data in Fig. 5, that the high temperature viscosity of NaGaSi₃O₈ melts.

Thus, it is clear from the available data that the viscosities of melts in the Na₂O-Fe₂O₃-SiO₂, Na₂O-B₂O₃-SiO₂ and Na₂O-Ga₂O₃-SiO₂ systems are significantly lower than the viscosities of equivalent melts in the Na₂O-Al₂O₃-SiO₂ system. Moreover, all the available data on the viscosity-temperature relationships of melts in these systems indicate that the *relative* bulk viscosities of melts in these four systems are not sensitive functions of temperature. On the basis of the premise that melt structure determines melt viscosity, and on the evidence (reviewed below) for the coordination of Al³⁺, Fe³⁺, B³⁺ and Ga³⁺ in these melts; we now propose to rationalize the relative bulk viscosities of melts in the systems Na₂O-Al₂O₃-SiO₂, Na₂O-Fe₂O₃-SiO₂, Na₂O-B₂O₃-SiO₂ and Na₂O-Ga₂O₃-SiO₂ in terms of the structural effects of substituting each of these trivalent cations for Si in these melts.

DISCUSSION

Discussion of the significance of the temperature- and composition-dependence of viscosities of any iron-bearing melts is made difficult by changes in redox state as a function of temperature and bulk composition. The relationship between the redox state and viscosities of melts in the Na₂O-FeO-Fe₂O₃-SiO₂ system has recently been investigated in a companion paper (DINGWELL and VIRGO, 1987). Figures 2 and 3 of DINGWELL and VIRGO (1987) illustrate that viscosities of melts in this study containing $\approx 95\%$ ferric iron are lower than the viscosities one would estimate for the fully oxidized (100% ferric iron) equivalent melts by a log₁₀ factor of ≈ 0.05 . This difference between viscosities of melts equilibrated in air and their fully oxidized equivalents is insignificant in the discussion and comparisons to follow where the iron-bearing melts of this study are approximated to represent fully oxidized melts.

Similarly, the temperature-dependence of these melts is complicated by the temperature-dependence of redox equilibria. These melts will become more oxidized with decreasing temperature with the result that the temperature-dependence of their viscosities will be larger than that for an isochemical (constant ferric/ferrous ratio) melt. No further quantitative discussion of the temperature-dependence of viscosities of these iron-bearing melts is attempted at this time.

Coordination of Al^{3+} , Fe^{3+} , B^{3+} and Ga^{3+}

The short-range structure of melts in the Na₂O-Al₂O₃-SiO₂ system has been studied extensively. The structural role of aluminum in subaluminous (mol% Na = mol% Al) and peralkaline (mol% Na > mol% Al) melts in the Na₂O-Al₂O₃-SiO₂ system is generally agreed to be that of a tetrahedrallycoordinated trivalent cation which may be classified as a network-former. The charge deficiency of a trivalent cation coordinating four bridging oxygens is compensated by a sodium cation near the aluminate tetrahedron which stabilizes the tetrahedral coordination of aluminum and the sharing of bridging oxygens between aluminate and silicate tetrahedra. The sharing of oxygens between tetrahedral cations is the essential requirement for polymerization or network formation in molten silicates so the sodium cation may be termed a network stabilizer.

Evidence for this structure model of melts in the portion of the Na₂O-Al₂O₃-SiO₂ system of concern to us here (that portion bounded by the tectosilicate and tetrasilicate joins) comes from a variety of structural probes, including Raman spectroscopy (SHARMA et al., 1978; MYSEN et al., 1981, 1985b; MCMILLAN et al., 1982; SEIFERT et al., 1982; MATSON and SHARMA, 1985; MATSON et al., 1986), X-ray fluorescence spectra (DAY and RINDONE, 1962), X-ray radial distribution functions (XRDF) (TAYLOR and BROWN, 1979a,b) and luminescence and X-ray photoelectron spectroscopy (ONOR-ATO et al., 1986). Additionally, indirect evidence for fully polymerized melts and glasses along the SiO₂-NaAlO₂ join (and therefore complete tetrahedral coordination of Al) is provided by discontinuities in composition-dependence of several physical properties at Na/Al = 1 (e.g., viscosity, RIE-BLING, 1966; density and index of refraction, DAY and RIN-DONE, 1962; He diffusivity, SHELBY and EAGAN, 1976 and the redox ratio Fe²⁺/Fe³⁺, DICKENSON and HESS, 1981,1986a).

The coordination of ferric iron in silicate melts has also received considerable attention in recent years. Compositions in the Na₂O-FeO-Fe₂O₃-SiO₂ system have been investigated by several methods, including ⁵⁷Fe Mössbauer spectroscopy (DANCKWERTH and VIRGO, 1982; FLEET *et al.*, 1984; DING-WELL and VIRGO, 1987), EXAFS (BROWN *et al.*, 1979), XRDF (HENDERSON *et al.*, 1984), optical absorption and luminescence spectroscopy (Fox *et al.*, 1982), and ESCA (GOLDMAN, 1986). The consensus of the above workers is that ferric iron is tetrahedrally coordinated in the relatively oxidized melts in the Na₂O-FeO-Fe₂O₃-SiO₂ system that result from equilibration with air.

Glasses in the Na₂O-B₂O₃-SiO₂ system have been investigated using B¹¹ NMR spectroscopy (YUN and BRAY, 1978). Several B-bearing structural species, involving both trigonal and tetrahedral coordination of B³⁺, have been identified. YUN and BRAY suggest that the B¹¹ NMR spectra of glasses on the SiO₂-NaBO₂ join are consistent with tetrahedral coordination of B^{3+} (as $BSi_4O_{10}^-$ units) in the silicate network, from pure SiO₂ to SiO₂/B₂O₃ = 8 (or mole% NaBO₂ = 20%) on the SiO₂-NaBO₂ join). Beyond this NaBO₂ content, the structure of the glasses contains an increasing fraction of B^{3+} which is excluded from the silicate network to form borate structural units involving tetrahedral and/or trigonal coordination of B³⁺ (e.g., ring-type metaborate, diborate; YUN and BRAY, 1978). OESTRIKE et al. (1986) have provided additional high resolution NMR evidence for the both tetrahedral and trigonal coordination of B³⁺ in glasses on the reedmergnerite-albite join (NaBSi₃O₈-NaAlSi₃O₈). The proportion of trigonally coordinated B3+ decreases with increasing reedmergnerite content, and the proportion of B^{3+} in trigonal coordination in reedmergnerite glass (25 mole%) NaBO₂) is given by OESTRIKE et al. (1986) as 29%.

The coordination of gallium in glasses in the Na₂O-Ga₂O₃-SiO₂ system has been investigated by Raman spectroscopy (VIRGO *et al.*, 1979; HENDERSON *et al.*, 1985; MATSON and SHARMA, 1985) and X-ray absorption spectroscopy (EX-ANES and EXAFS; FLEET *et al.*, 1984). These structural studies are in agreement in concluding that Ga³⁺ is tetrahedrally-coordinated in melts in the Na₂O-Ga₂O₃-SiO₂ system. This structural inference is supported by the binary nature of Ga-Al interdiffusion reported by KUSHIRO (1983) and (similar to the case for aluminum) by discontinuities in the temperature-dependence of physical properties such as viscosity (LAPP and SHELBY, 1986), He diffusivity (SHELBY, 1981) and the Fe²⁺/Fe³⁺ ratio (DICKENSON and HESS, 1986b) at the Na/Ga = 1 compositional join which suggest that these (Na/Ga = 1) melts are fully polymerized.

Comparison of structure and properties of "tectosilicate" melts

If melts along the SiO₂-Na(X)O₂ joins (X = Al, Fe, B, Ga) exhibited tetrahedral coordination of all trivalent cations, then each of these "tectosilicate" compositions would yield structures that are equivalent in the sense that all are fully polymerized tectosilicate frameworks. As reviewed above, completely tetrahedral coordination is evidenced by spectroscopic studies of aluminosilicate, ferrosilicate and galliosilicate glasses, whereas borosilicate glasses are depolymerized by the existence of trigonally-coordinated B3+. Non-tetrahedral coordination of some B³⁺ is also supported by the observation that melts along the SiO₂-NaBO₂ join do not represent viscosity maxima along constant SiO₂ joins (TAIT et al., 1984) as is the case in the aluminosilicate and galliosilicate systems (RIEBLING, 1966; LAPP and SHELBY, 1986). Thus, the lower bulk viscosities of the borosilicate melts may be readily attributed to the trigonal coordination of some B^{3+} and the resulting depolymerized state of these melts. The contrast between the bulk viscosities of the aluminosilicate melts versus the ferrosilicate and galliosilicate melts requires a different structural explanation.

Cationic electronegativities

The relative viscosities of melts on the SiO2-NaFeO2, SiO2-NaAlO2, SiO2-NaBO2 (and possibly SiO2-NaGaO2) joins are inversely related to the electronegativities of the trivalent cations. The electronegativity of a cation may be viewed as the electron affinity of an ion during bond formation and the difference in cationic and anionic electronegativities is a measure of the ionicity (or percent ionic character) of the chemical bond. Although quantitative application of published cationic electronegativity scales (derived for the gaseous state) to actual cation-oxygen bond strengths in condensed silicates is not recommended, some qualitative observations may be made. Typical cationic electronegativities of Ga, Fe and B are 1.6, 1.9 and 2.0, respectively (PAULING, 1960). All are higher than Al (1.45) and thus Fe, B and Ga can be expected to compete with Si for the bonding electrons of bridging oxygens more effectively than Al. Such competition may result in increased tetrahedral bond angle distortion and the resulting creation of more than one structural unit of intermediate range. The importance of electronegativities in influencing the Raman spectra of sodium alumino- and galliosilicate melts has been noted previously by MATSON and SHARMA (1985).

Viscosity data from URBAIN *et al.* (1982) for alkali and alkaline earth aluminosilicate melts, presented in Fig. 6 reveal further systematics involving cationic electronegativities and melt viscosities. Figure 6 illustrates that melts along the tectosilicate joins, SiO_2 -XAl₂O₄ (where X = Ba, Sr, Ca, Mg) exhibit an inverse correlation between viscosity and the electronegativity of the (in this case) network-stabilizing, alkaline earth cation [Pauling electronegativities are: Ba (0.9), Sr (1.0), Ca (1.0), Mg (1.2)]. The data available for these melts (Fig. 6) indicate that the order of decreasing viscosity does not depend on either temperature or SiO₂ content.

Figure 6 illustrates the same general behavior for alkalibearing melts. The sparse data for melts on the SiO₂-XAlO₂ joins (where X = K, Na, Li; LiAlSi₃O₈ data from SHELBY, 1978) are consistent with an inverse relationship between melt viscosity and cation electronegativity [Pauling electronegativities are: K (0.8), Na (0.9), Li (1.0)].

One further point of comparison between tectosilicate melt viscosities is that the viscosities of the alkaline earth aluminosilicate melts are lower than those of the alkali aluminosilicate melts. The lower viscosities of the alkaline earth melts would then indicate that our electronegativity-viscosity relationship cannot be applied directly across chemical groups. We will return to the comparison of alkali and alkaline earth aluminosilicate melt viscosities in a later section on calorimetric data.

T-O-T bond angle variation

Several studies have proposed that the NaAlSi₋₁ substitution results in a decrease in average T-O-T bond angle of melts along the SiO₂-NaAlO₂ join and the estimates of relative change in T-O-T bond angle from SiO₂ to NaAlSiO₄ (*i.e.* 50% completion of the NaAlSi₋₁ substitution) from Raman spectra (SEIFERT *et al.*, 1982) and from XRD data (KONNERT and KARLE, 1973; TAYLOR and BROWN, 1979b; DEJONG and BROWN, 1980) are (comparable at) 4.6% and 5.8%, respectively. Estimation of the average T-O-T bond angle in vitreous SiO₂ and NaFeSi₃O₈ by HENDERSON *et al.* (1984) yields a change in T-O-T bond angle, with only 25% completion of the NaFeSi₋₁ exchange (from silica to Fe-albite), of 7.3%. Thus, comparison of the decreases of the average T-O-T bond angles for the SiO₂-NaAlO₂ and SiO₂-NaFeO₂

joins is consistent with the suggestion that the effect of the NaFeSi₋₁ substitution is much larger (7.3% at 25% substitution) than the effect of the NaAlSi₋₁ substitution (\approx 4-6% at 50% substitution).

VIRGO et al. (1979) and HENDERSON et al. (1985) have discussed their Raman spectral data of glasses along the SiO₂-NaGaO₂ join in terms of potential intermediate range structures. It is evident from Fig. 33 of VIRGO et al. (1979) that the frequency shift for the Si-O (bridging) stretching mode (at approx. 1070 cm⁻¹ in pure SiO₂) as a function of composition is larger for glasses on the SiO₂-NaFeO₂ and SiO₂-NaGaO₂ joins. This observation is consistent with the suggestion that the NaFeSi₋₁ and NaGaSi₋₁ substitutions have a larger effect than the NaAlSi₋₁ substitution on the tectosilicate structure.

Thus, the evidence for the structural effects of both the NaGaSi₋₁ and the NaFeSi₋₁ substitutions may be interpreted similarly in that incorporation of Fe-tetrahedra or Ga-tetrahedra into SiO₂ melt results in perturbation of the network structure as evidenced by the large decrease observed in the average T-O-T bond angle.

Related information for a hypothetical tectosilicate structure of melts on the SiO₂-NaBO₂ join comes from the ab initio molecular orbital calculations of GEISINGER et al. (1985). NAVROTSKY et al. (1985) present the variation of calculated potential energy as a function of T-O-T bond angle for B-O-B, B-O-Si, Al-O-Al, Al-O-Si and Si-O-Si derived from calculations for groups of the general formula $H_6X_2O_7^{n-1}$ (where X = Si, Al, B and n = 2-4). The potential energy curves indicate that for the aluminosilicate group, the Si-O-Al, Al-O-Al and Si-O-Si bonds have relatively broad minima as a function of T-O-T bond angle and the minima are in the T-O-T bond angle range of 135 to 150°. In contrast, the borosilicate groups exhibit relatively narrow potential energy minima for the B-O-B and B-O-Si bonds that occur at lower values of the T-O-T bond angle (125 to 133°). NAVROTSKY et al. (1985) suggest that the lower value of T-O-T bond angle and the relative "inflexibility" of this angle for bonds involving B^{3+} rather than Al^{3+} poses a potential energy barrier to the substitution of borate for silicate tetrahedra in borosilicate melts.

Enthalpy of solution

The heat of reaction obtained from solution of glasses $(LN_2--quenched from liquids at 1893 K)$ of the general formula $X_{1/n}^{n+}AlO_2$ -SiO₂ (where M = K, Na, Ba, Sr, Ca, Mg) into 2PbO \cdot B₂O₃ at 973 K are provided by ROY and NAVROTSKY (1984, Fig. 2). The relative values of heats of solution (H_{sol}) provide information on the relative strengths of bonds broken during the dissolution process. The values of H_{sol} become more exothermic with increasing cation electronegativity (from K to Mg) and, as noted by NAVROTSKY *et al.* (1985), the decrease in average bond strength, inferred from more exothermic values of H_{sol} , correlates with decreasing viscosities for these compositions.

Figure 7 is a comparison of viscosity (1400°C) and H_{sol} (2PbO \cdot B₂O₃, 973 K) for aluminosilicate melts of the general formula $X_{1/n}^{n+}$ AlSi₃O₈ (where X = K, Na, Ba, Sr, Ca, Mg). A positive correlation exists between viscosity and H_{sol} for these melts containing 75 mol% SiO₂. In order to minimize the

FIG. 6. Comparison of the viscosity-temperature relationships of both alkaline earth and alkali aluminosilicate melts. Data from SHELBY (1978); URBAIN *et al.* (1982).





FIG. 7. Viscosities of melts of the general composition $X_{1/n}^{n+}AlSi_3O_8$ (where X = Mg, Ca, Sr, Ba, Na, K) at 1400°C plotted versus the enthalpy of solution of equivalent glasses obtained in molten 2PbO \cdot B₂O₃ at 973 K. Viscosity data (in poise) from same sources as Fig. 6; enthalpy data (in kJ/mole based on 2 oxygens) from ROY and NAVROTSKY (1984). Discrepant albite (Na) data are from URBAIN *et al.* (1982) (higher value) and SCARFE and CRONIN (1986) (lower value); see text for discussion.

imprecision of viscosity data arising from comparison of different studies, all viscosity data were chosen from the study of URBAIN *et al.* (1982). The H_{sol} data are all from the study by ROY and NAVROTSKY (1984). A positive correlation between viscosity and H_{sol} is apparent. The data are insufficient to indicate whether the correlation is best described as one curve including both alkali- and alkaline earth-bearing melts or two subparallel curves, one each for alkali- and alkaline earth-bearing melts. High temperature viscosity data for LiAlSi₃O₈ are needed.

If the correlation of Fig. 7 is a single curve then the origin of the differences in viscosity between aluminosilicates on the tectosilicate joins may be attributed to differences in average bond strength. If there are, instead, two curves in Fig. 7, then average bond strength cannot entirely explain the relative melt viscosities, and aluminate tetrahedral ordering becomes a likely explanation for the difference between the two curves.

CONCLUSIONS

The examination of the available data on the structure and viscosity-temperature relationships of tectosilicate melts in the Na₂O-Al₂O₃-SiO₂, Na₂O-Fe₂O₃-SiO₂, Na₂O-B₂O₃-SiO₂ and Na₂O-Ga₂O₃-SiO₂ systems and comparison with equivalent data for both the alkaline earth and other alkali aluminosilicate systems, in combination with calorimetric and spectroscopic studies, leads to the following conclusions.

The viscosity-temperature relationships of these melts yield isothermal bulk viscosities that decrease in the order of aluminosilicate > ferrosilicate, (galliosilicate?) > borosilicate. The relative viscosities do not appear to be sensitive to the temperature of comparison.

The relative viscosities of tectosilicate melts are inversely correlated with the electronegativity of the substituted cation both for monovalent and divalent network-stabilizer cations e.g., Li, Na, K and Mg, Ca, Sr, Ba) and for trivalent networkformer cations. A substitution of the monovalent or divalent network-stabilizer cation or the trivalent network-former cation for a cation of greater electronegativity, yields a lower viscosity that results from weaker T-O-T bonds and a more distorted T-O-T bond angle population. The latter of these two effects (T-O-T bond angle distortion) may result, in the case of the network-former substitution, in a change in the ordering of trivalent network-formers. Thus the viscosity reduction with substitution to cations of higher electronegativity is explicable in terms of (1) weakened T-O bonds, and (2) T-O-T bond angle distortion (with possible ordering of tetrahedral cations). The estimate of average bond strength from solution calorimetry of silicate glasses may provide, in the future, evidence for a distinction between these two effects.

Acknowledgements—The authors wish to thank P. Hess, M. Fleet and an anonymous reviewer for their comments.

Editorial handling: P. C. Hess

REFERENCES

- BAILEY D. K. and MACDONALD R. (1987) Dry peralkaline felsic liquids and carbon dioxide flux through the Kenya rift zone. In Magmatic Processes: Physicochemical Principles (ed. B. O. MY-SEN). Geochem. Soc., Special Publication 1, 91-105.
- BOTTINGA Y. and WEILL D. F. (1972) The viscosity of magmatic silicate liquids: A model for calculation. *Amer. J. Sci.* 272, 438-475.
- BOWEN N. L., SCHAIRER J. F. and WILLEMS H. W. V. (1930) The ternary system: Na₂SiO₃-Fe₂O₃-SiO₂. Amer. J. Sci. 20, 405–455.
- BROWN G., KEEFER K. D. and FENN P. M. (1979) Extended X-ray absorption fine structure (EXAFS) study of iron-bearing silicate glasses: Iron coordination environment and oxidation state. (abstr.) Geol. Soc. Amer. Abstr. Prog. 11, 373.
- DANCKWERTH P. A. and VIRGO D. (1982) Structural state of iron in the system Na₂O-SiO₂-Fe-O. Carnegie Inst. Wash. Yearb. 81, 340-342.
- DAY D. E. and RINDONE G. E. (1962) Properties of soda aluminosilicate glasses. J. Amer. Ceram. Soc. 45, 489-504, 579-581.
- DEJONG B. H. W. S. and BROWN G. E. (1980) The polymerization of silicate and aluminate tetrahedra in glasses, melts and aqueous solutions—I. Electronic structure of H₆Si₂O₇, H₆AlSi₂O₇ and H₆Al₂O₇. Geochim. Cosmochim. Acta 44, 491-511.
- DICKENSON M. P. and HESS P. C. (1981) Redox equilibria and the structural role of iron in aluminosilicate melts. *Contrib. Mineral. Petrol.* 78, 352–357.
- DICKENSON M. P. and HESS P. C. (1986a) The structural role and homogeneous redox equilibria of iron in peraluminous, metaluminous and peralkaline silicate melts. *Contrib. Mineral. Petrol* 92, 207-217.
- DICKENSON M. P. and HESS P. C. (1986b) The structural role of Fe^{3+} , Ga^{3+} , Al^{3+} and homogeneous iron redox equilibria in K₂O-Al₂O₃-Ga₂O₃-SiO₂-Fe₂O₃-FeO melts. J. Non-Cryst. Solids **86**, 303-310.
- DINGWELL D. B. (1986) Viscosity-temperature relationships in the system Na₂Si₂O₅-Na₄Al₂O₅. Geochim Cosmochim. Acta 50, 1261– 1265.
- DINGWELL D. B. and VIRGO D. (1987) The effect of oxidation state on the viscosity of melts in the system Na_2O -FeO-Fe₂O₃-SiO₂. *Geochim. Cosmochim. Acta* **51**, 195-205.
- FISHER R. V. and SCHMINKE H.-U. (1984) Pyroclastic Rocks. Springer-Verlag, 472p.
- FLEET M. E., HERZBERG C. T., HENDERSON G.S., CROZIER E. D., OSBORNE M. D. and SCARFE C. M. (1984) Coordination of Fe. Ga and Ge in high pressure glasses by Mössbauer, Raman and Xray absorption spectroscopy, and geological implications. Geochim. Cosmochim. Acta 48, 1455-1466.
- FOX K. E., FURUKAWA T. and WHITE W. B. (1982) Transition metal ions in silicate melts. Part 2, Iron in sodium silicate glasses. *Phys. Chem. Glasses* 32, 169–178.
- GEISINGER K. L., GIBBS G. V. and NAVROTSKY A. (1985) A molecular orbital study of bond length and angle variations in framework structures. *Phys. Chem. Minerals* 11, 266-283.
- GOLDMAN D. S. (1986) Evaluation of the ratios of bridging to nonbridging oxygens in simple silicate glasses by electron spectroscopy for chemical analysis. *Phys. Chem. Glasses* 27, 128–133.

- HENDERSON G. S., FLEET M. S. and BANCROFT G. M. (1984) An X-ray scattering study of vitreous KFeSi₃O₈ and NaFeSi₃O₈ and reinvestigation of vitreous SiO₂ using quasi-crystalline modelling. J. Non-Cryst. Solids **68**, 333-349.
- HENDERSON G. S., BANCROFT G. M., FLEET M. E. and ROGERS D. J. (1985) Raman spectra of gallium and germanium substituted silicate glasses: Variations in intermediate range order. *Amer. Mineral.* **70**, 946–960.
- KONNERT J. H. and KARLE L. (1973) The computation of radial distribution functions for glassy materials. *Acta Crystal.* 29, 702-710.
- KUSHIRO I. (1983) Effect of pressure on the diffusivity of networkforming cations in melts of jadeitic composition. *Geochim. Cos*mochim. Acta 47, 1415-1422.
- LAPP J. C. and SHELBY J. E. (1986) Viscosity and thermal expansion of sodium and potassium galliositicate glasses. J. Amer. Ceram. Soc. 69, 126-131.
- MATSON D. W. and SHARMA S. K. (1985) Structures of sodium alumino- and gallosilicate glasses and their germanium analogs. *Geochim. Cosmochim. Acta* 49, 1913-1924.
- MATSON D. W. and SHARMA S. K. and PHILPOTTS J. A. (1986) Raman spectra of some tectosilicates and of glasses along the orthoclase-anorthite and nepheline-anorthite joins. *Amer. Mineral.* 71, 694-704.
- MCMILLAN P. F., PIRIOU B. and NAVROTSKY A. (1982) A Raman spectroscopic study of glasses along the joins silica-calcium aluminate, silica-sodium aluminate and silica potassium aluminate. *Geochim. Cosmochim. Acta* 46, 2021–2037.
- MYSEN B. O. (1988) Relationships between structure, redox equilibria of iron and properties of magmatic liquids. In *Advances in Physical Chemistry* (eds. L. PERCHUK and I. KUSHIRO). In press.
- MYSEN B. O., SEIFERT F. A. and VIRGO D. (1980) Structure and redox equilibria of iron-bearing silicate melts. *Amer. Mineral.* 65, 867–884.
- MYSEN B. O., VIRGO D. and KUSHIRO I. (1981) The structural role of aluminum in silicate melts—a Raman spectroscopic study at 1 atmosphere. *Amer. Mineral.* 66, 678–701.
- MYSEN B. O., VIRGO D., NEUMANN E.-R. and SEIFERT F. A. (1985a) Redox equilibria and the structural states of ferric and ferrous iron in melts in the system CaO-MgO-Al₂O₃-SiO₂-Fe-O: Relationships between redox equilibria, melt structure and liquidus phase equilibria. *Amer. Mineral.* **70**, 317-331.
- MYSEN B. O., VIRGO D. and SEIFERT F. A. (1985b) Relationships between properties and structure of aluminosilicate melts. *Amer. Mineral.* **70**, 834–847.
- NAVROTSKY A., GEISINGER K. L., MCMILLAN P. and GIBBS G. V. (1985) The tetrahedral framework in glasses and melts—inferences from molecular orbital calculations and implications for structure, thermodynamics, and physical properties. *Phys. Chem. Mineral.* 11, 284–298.
- OESTRIKE R., GEISINGER K., NAVROTSKY A. TURNER G. L. and KIRKPATRICK R. J. (1986) Structure and thermochemistry of glasses along the join NaAlSi₃O₈-NaBSi₃O₈: The effect of boron. *Geol. Soc. Amer. Abstr. Prog.* 18, 709.
- ONORATO P. I. K., ALEXANDER M. N., STRUCK C. W., TASKER G. W. and UHLMANN D.R. (1986) Bridging and nonbridging oxygen atoms in alkali aluminosilicate glasses. J. Amer. Ceram. Soc. 68, C-148-150.
- PAULING L. (1960) The Nature of the Chemical Bond. Cornell Univ. Press, Ithaca.
- PIGUET J. L. and SHELBY J. E. (1985) Transformation-range behavior

of Li₂O-(Al,Ga)₂O₃-SiO₂ glasses. J. Amer. Ceram. Soc. 68, C-232-233.

- PIGUET J. L., LAPP J. C. and SHELBY J. E. (1985) Transformationrange behavior of lithium galliosilicate glasses. J. Amer. Ceram. Soc. 68, 326-329.
- RICHET P. (1984) Viscosity and configurational entropy of silicate melts. Geochim. Cosmochim. Acta 48, 471-483.
- RIEBLING E. F. (1966) Structure of sodium aluminosilicate melts containing at least 50 mole% SiO₂ at 1500°C. J. Chem. Phys. 44, 2857-2865.
- ROY B. N. and NAVROTSKY A. (1984) Thermochemistry of chargecoupled substitutions in silicate glasses: The systems M^{*}_{1/n}AlO₂-SiO₂ (M = Li, Na, K, Rb, Cs, Mg, Ca, Sr, Ba, Ph). J. Amer. Ceram. Soc. 67, 606-610.
- SCARFE C. M. and CRONIN D. J. (1986) Viscosity-temperature relationships of melts at 1-atm in the system diopside-albite. Amer. Mineral. 71, 767-771.
- SEIFERT F. A., MYSEN B. O. and VIRGO D. (1982) Three dimensional network structure of quenched melts (glass) in the systems SiO₂-NaAlO₂, SiO₂-CaAl₂O₄ and SiO₂-MgAl₂O₄. Amer. Mineral. 67, 696-717.
- SHARMA S. K., VIRGO D. and MYSEN B. O. (1978) Structure of melts along the join SiO₂-NaAlSiO₄ by Raman spectroscopy. Carnegie Inst. Wash. Yearb. 77, 652–658.
- SHELBY J. E. (1978) Viscosity and thermal expansion of lithium aluminosilicate glasses. J. Appl. Phys. 49, 5885-5891.
- SHELBY J. E. (1981) Helium migration in alkali galliosilicate glasses. J. Non-Cryst. Solids, 45, 411-418.
- SHELBY J. E. and EAGAN R. J. (1976) Helium migration in sodium aluminosilicate glasses. J. Amer. Ceram. Soc. 59, 420-425.
- TAIT J. C., MANDOLESI D. L. and RUMMENS H. E. C. (1984) Viscosity of melts in the sodium borosilicate system. *Phys. Chem. Glasses* 25, 100-104.
- TAYLOR M. and BROWN G. E. JR. (1979a) Structure of mineral glasses I. The feldspar glasses NaAlSi₃O₈, KAlSi₃O₈, CaAl₂Si₂O₈. *Geochim. Cosmochim. Acta* 43, 61–74.
- TAYLOR M. and BROWN G. E. JR. (1979b) Structure of mineral glasses II. The SiO₂-NaAlSiO₄ join. *Geochim. Cosmochim. Acta* 43, 1467-1473.
- TAYLOR I. D. and RINDONE G. E. (1970) Properties of aluminosilicate glasses: V. Low-temperature viscosities. J. Amer. Ceram. Soc. 53, 692-695.
- URBAIN G., BOTTINGA Y. and RICHET P. (1982) Viscosity of liquid silica, silicates and alumino-silicates. *Geochim. Cosmochim. Acta* **46**, 1061–1072.
- VIRGO D. and MYSEN B. O. (1985) The structural state of iron in oxidized versus reduced glasses at 1 atm: A ⁵⁷Fe Mössbauer study. *Phys. Chem. Mineral.* 12, 65-76.
- VIRGO D., SEIFERT F. A. and MYSEN B. O. (1979) Three-dimensional network structure of glasses in the systems CaAl₂O₄-SiO₂, NaAlO₂-SiO₂, NaFeO₂-SiO₂ and NaGaO₂-SiO₂ at 1 atm. *Carnegie Inst.* Wash. Yearb. 78, 506-511.
- VIRGO D., MYSEN B. O. and DANCKWERTH P. A. (1983) Redox equilibria and the anionic structure of Na₂O·xSiO₂-Fe-O melts: Effect of oxygen fugacity. *Carnegie Inst. Wash. Yearb.* 82, 305– 309.
- WILLIAMS H. and MCBIRNEY A. R. (1979) Volcanology. Freeman, Cooper & Co., San Francisco, 397p.
- YUN Y. H. and BRAY P. J. (1978) Nuclear Magnetic Resonance studies of glasses in the system Na₂O-B₂O₃-SiO₂. J. Non-Cryst. Solids 27, 363-380.